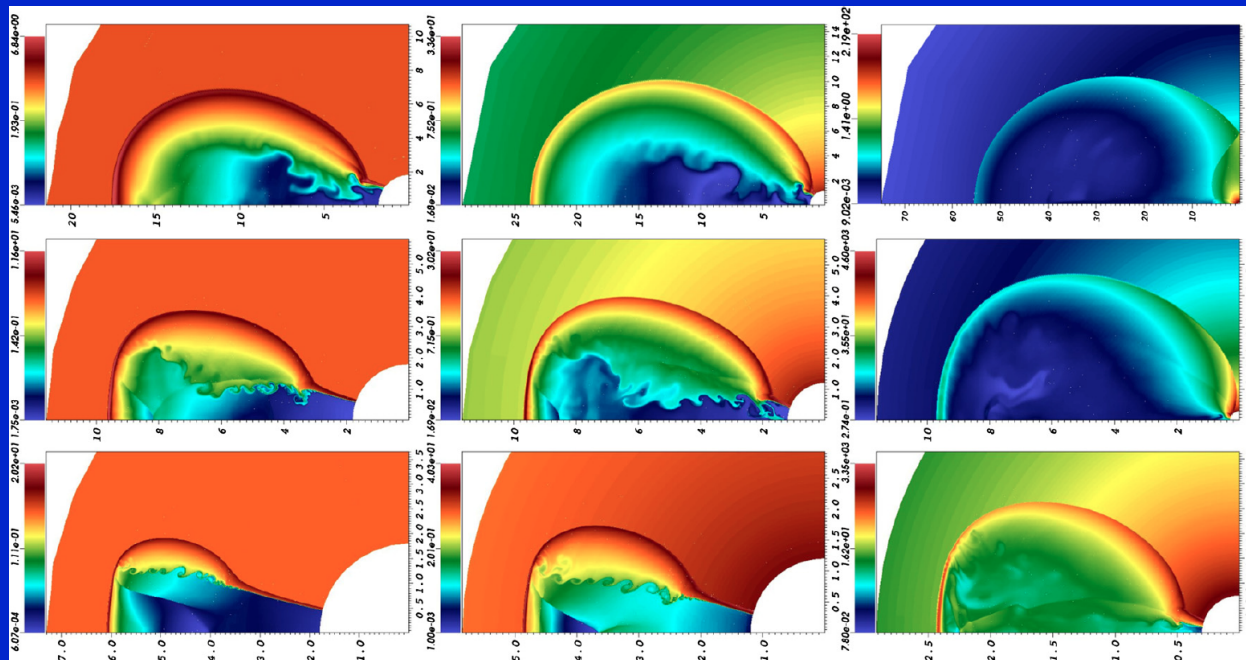


# GRB Jet Dynamics & Afterglow Lightcurves

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13<sup>th</sup> Marcel Grossmann meeting, Stockholm, July 6, 2012

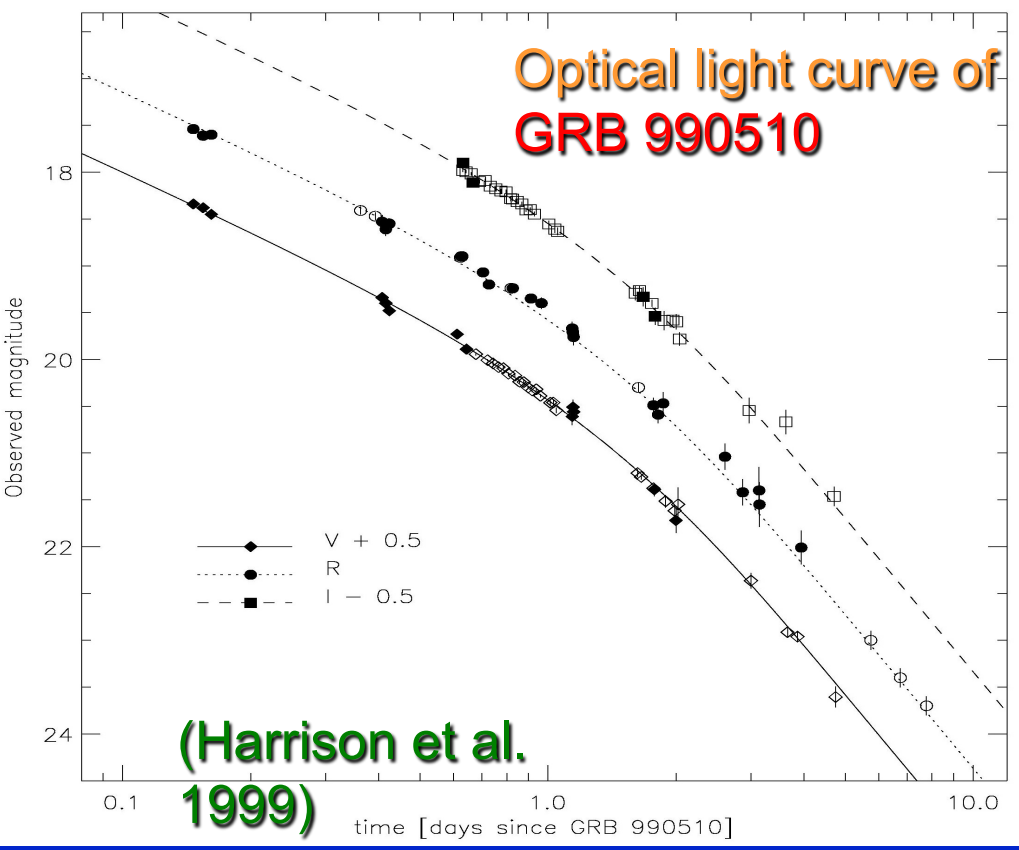
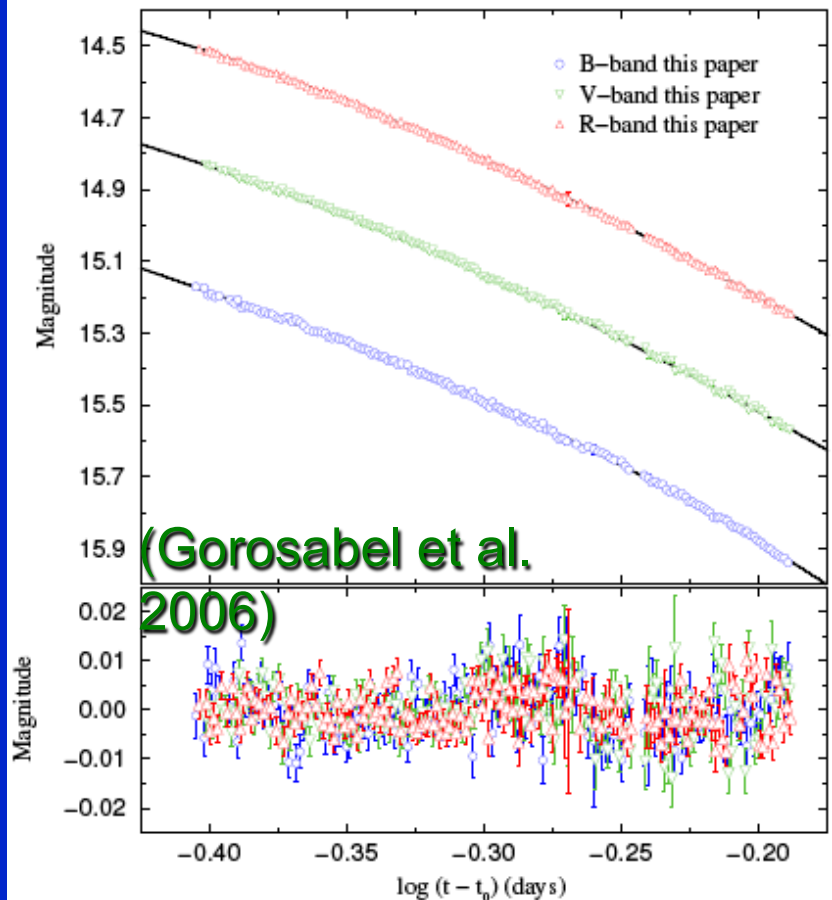
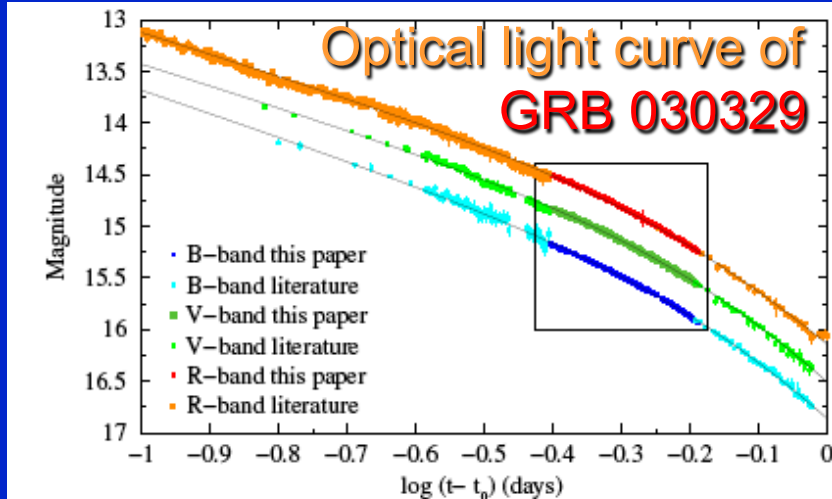
# Outline of the talk:

- Evidence for jets, angular structure, evolution stages
- Jet dynamics during the afterglow: an overview
- Analytic vs. numerical results: a problem?
- Recent numerical & analytic results: finally agree
- Simulations of an afterglow jet propagating into a stratified external medium:  $\rho_{\text{ext}} \propto R^{-k}$  for  $k = 0, 1, 2$
- Implications for GRBs: jet breaks, radio calorimetry

# Observational evidence for jets in GRBs:

- The energy output in  $\gamma$ -rays assuming isotropic emission ( $E_{\gamma,iso}$ ) approaches (and sometimes even exceeds)  $M_{\odot}c^2$ 
  - ◆ Difficult for a stellar mass progenitor
  - ◆ **True energy** is much smaller for a narrow jet
- At least some long GRBs occur together with a SN Ic
  - ◆ the outflow would contain  $>M_{\odot}$  if spherical
  - ◆ only a small part of this mass can reach  $\Gamma \gtrsim 100$   
& it would contain a small fraction of the energy
- Achromatic break or steepening of the afterglow light curves (“**jet break**”)

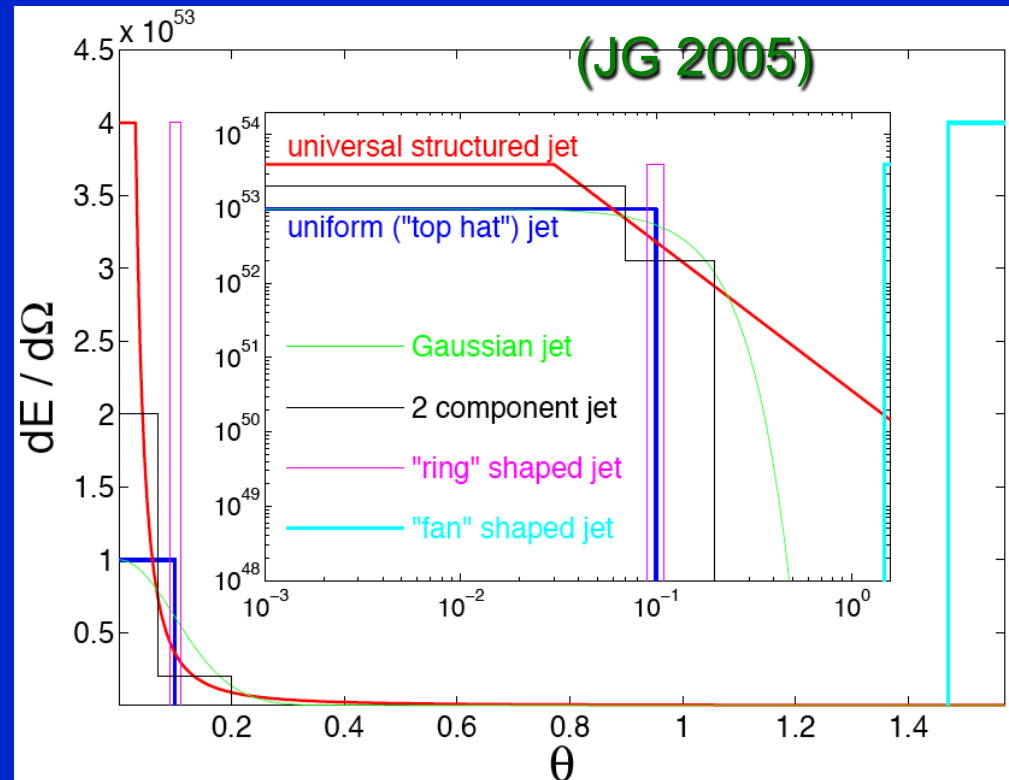
# Examples of Smooth & Achromatic Jet Breaks:





# The Angular Structure of GRB Jets:

- **Jet structure: unclear** (uniform, structured, hollow cone,...)
  - ◆ Affects  $E_{\gamma,iso} \rightarrow E_{\gamma}$  & observed GRB rate  $\rightarrow$  true rate
  - ◆ Viewing-angle effects (afterglow & prompt - XRF)
  - ◆ Can also affect late time radio calorimetry
- Here I consider mainly a uniform “top hat” jet

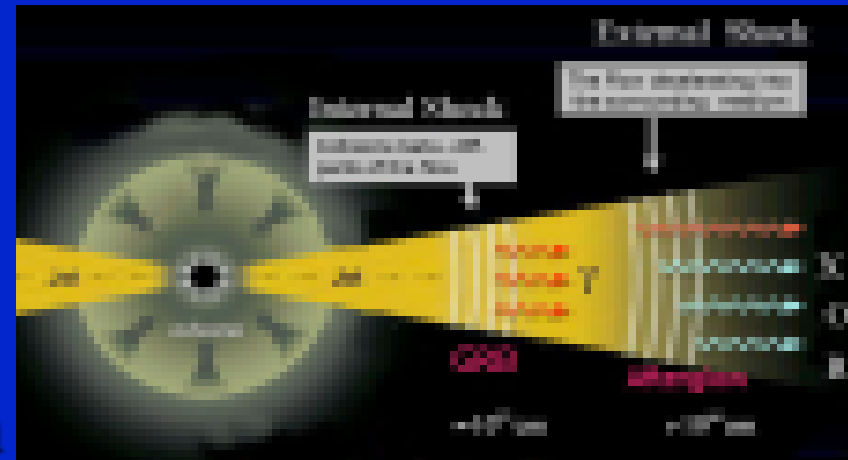


# Stages in the Dynamics of GRB Jets:

- **Launching** of the jet: magnetic (B-Z?) neutrino annihilation?
- **Acceleration**: magnetic or thermal?
- For long GRBs: propagation inside progenitor star
- **Collimation**: magnetic, stellar envelope, accretion disk wind
- **Coasting** phase that ends at the deceleration radius  $R_{\text{dec}}$
- At  $R > R_{\text{dec}}$  most of the energy is in the shocked external medium: the composition & radial profile are forgotten, but the angular profile persists (**locally: BM76 solution**)

■ Once  $\Gamma < 1/\theta_0$  at  $R > R_{\text{jet}}$  jet lateral expansion is possible

- Eventually the flow becomes spherical approaches the self-similar **Sedov-Taylor solution**



# Dynamics of GRB Jets: Lateral Expansion

- **Simple semi-analytic models** (Rhoads 97, 99; Sari, Piran & Halpern 99,...) make simplifying assumptions, such as:
  - ◆ The shock front is a part of a sphere within  $\theta < \theta_{\text{jet}}$
  - ◆ The velocity is in the radial direction (even at  $t > t_{\text{jet}}$ )
  - ◆ Lateral expansion at  $c_s \approx c/\sqrt{3}$  in the comoving frame
  - ◆ The jet dynamics are obtained by solving simple 1D equations for conservation of energy and momentum
- $\Rightarrow \Gamma \sim (c_s/c\theta_0)\exp(-R/R_{\text{jet}})$ ,  $\theta_{\text{jet}} \sim \theta_0(R_{\text{jet}}/R)\exp(R/R_{\text{jet}})$
- **Hydro-simulations**: these simplifying assumptions fail: shock front is aspherical, velocity is not radial,...
- **Very mild** lateral expansion while jet is relativistic
- Non-uniform shocked fluid: emission mainly from  $\theta < \theta_0$
- **Nevertheless**, despite the differences, there is a **sharp achromatic jet break** [for  $v > v_m(t_{\text{jet}})$ ] at a similar  $t_{\text{jet}}$

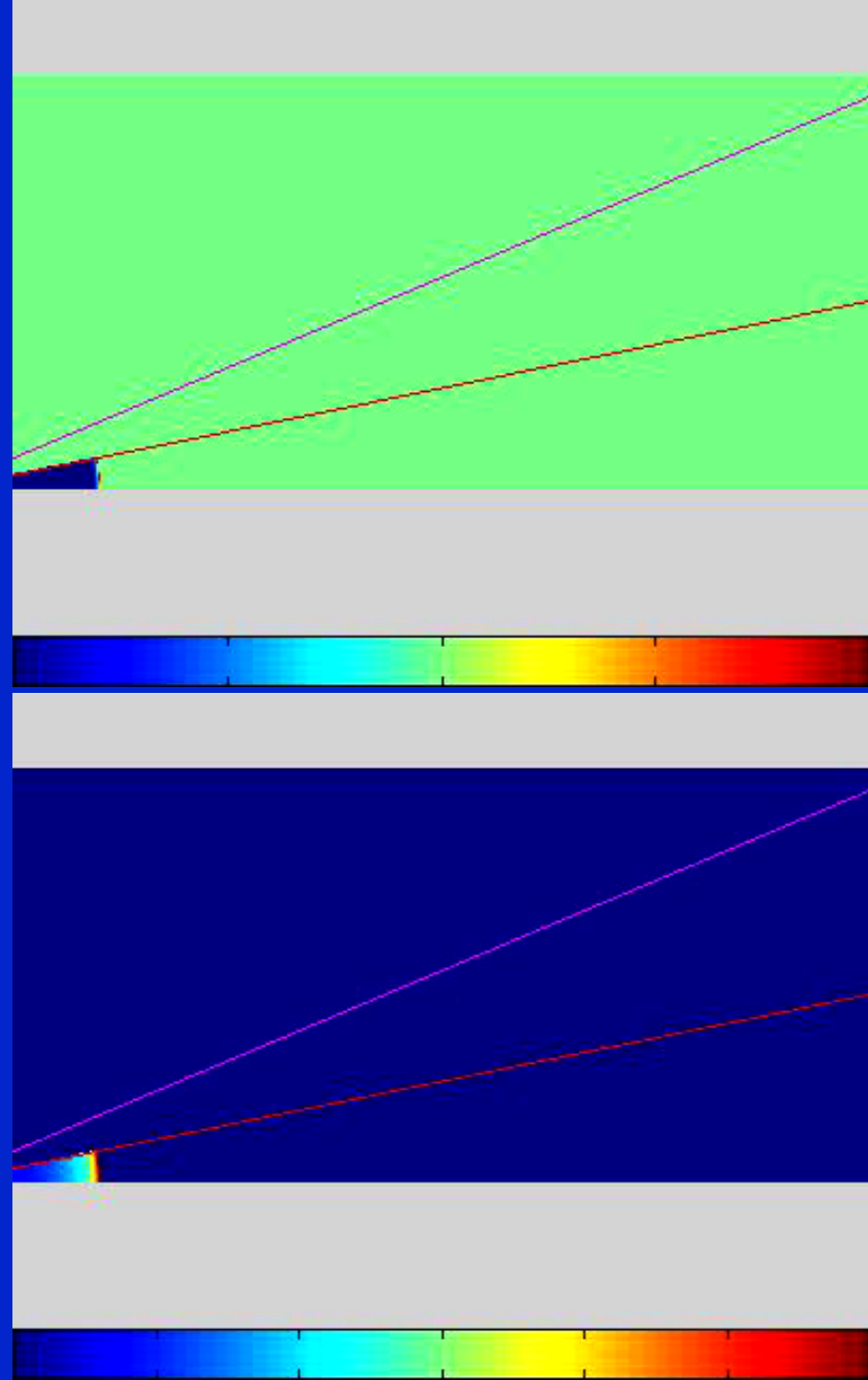
# 2D hydro-simulations

(JG et al. 2001)

Proper Density:  
(logarithmic color scale)

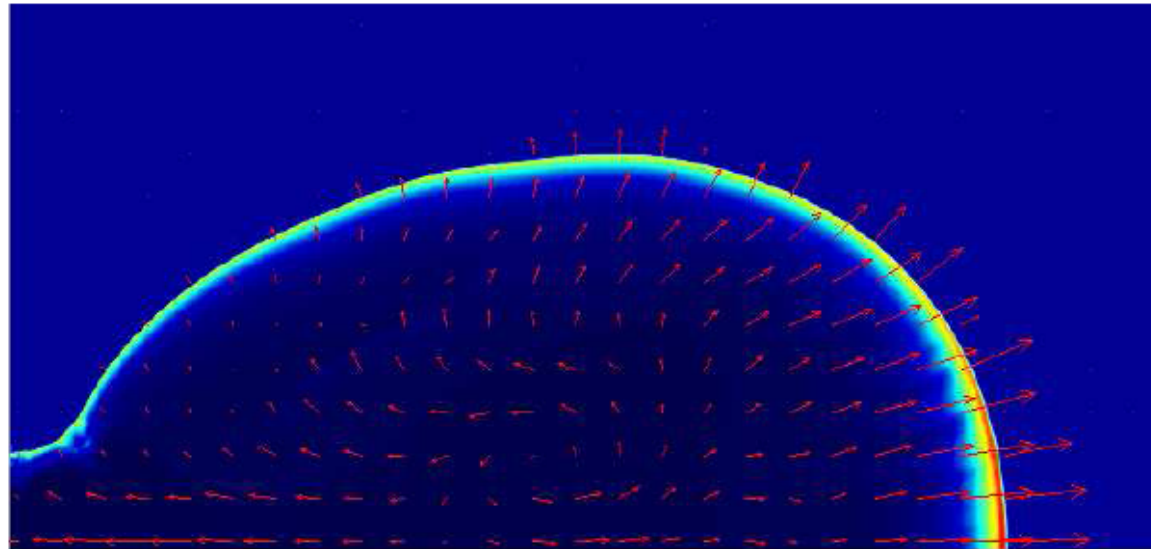
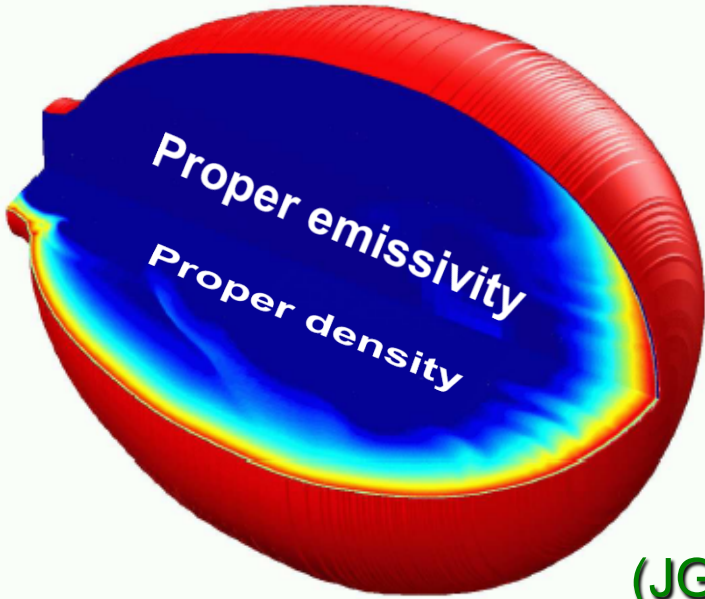
- Uniform external medium
- Initial conditions: a conical wedge from the BM solution

Bolometric  
Emissivity:  
(logarithmic color scale)





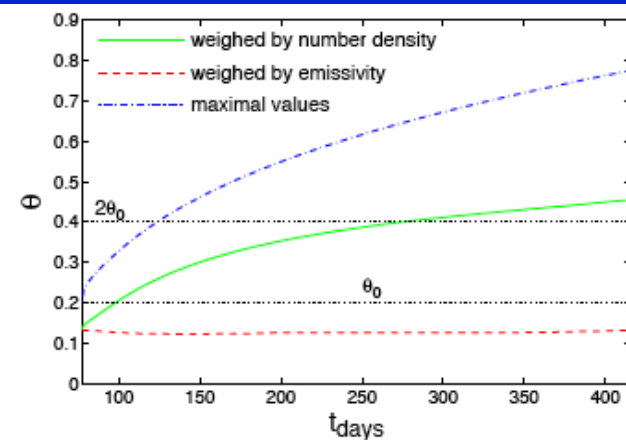
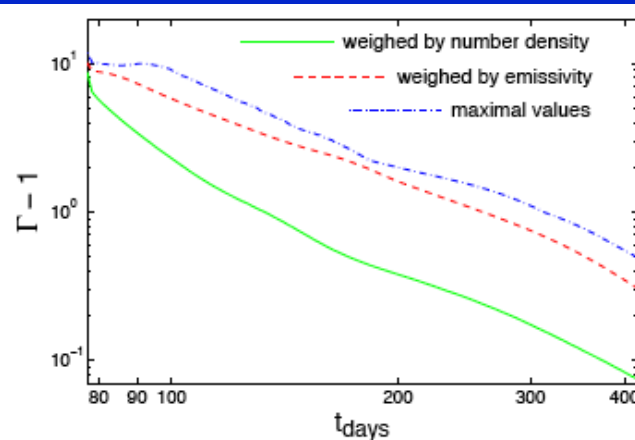
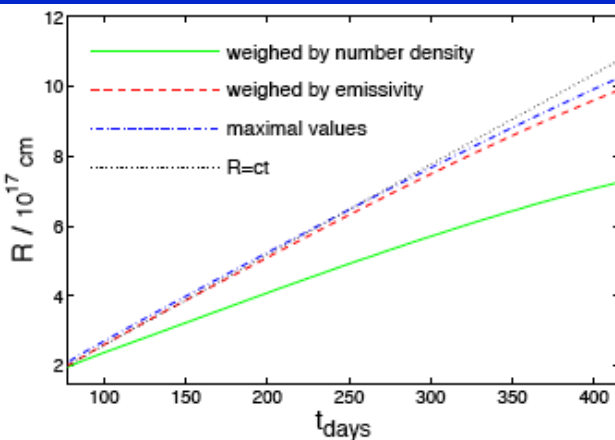
# The Jet Dynamics: very modest lateral expansion



(JG et al. 2001)

$0.911R^{-12}$   $0.163R^{-06}$

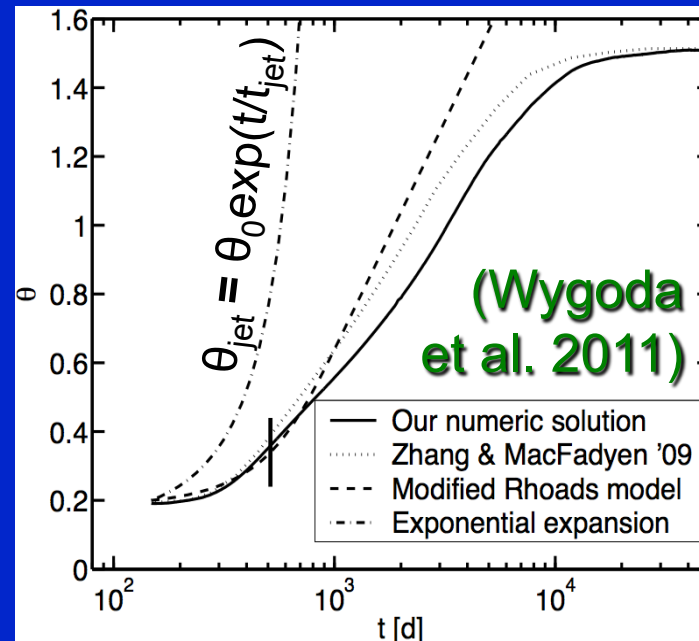
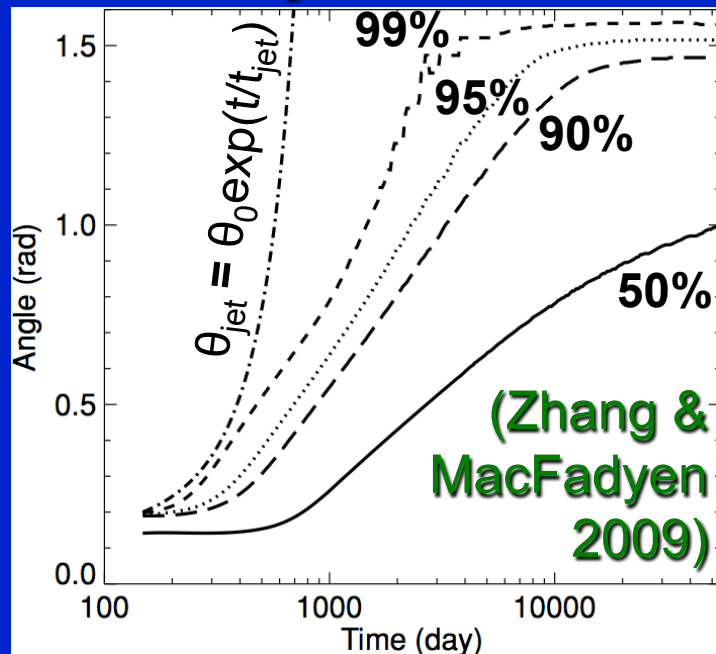
- There is slow material at the sides of the jet while most of the emission is from its front





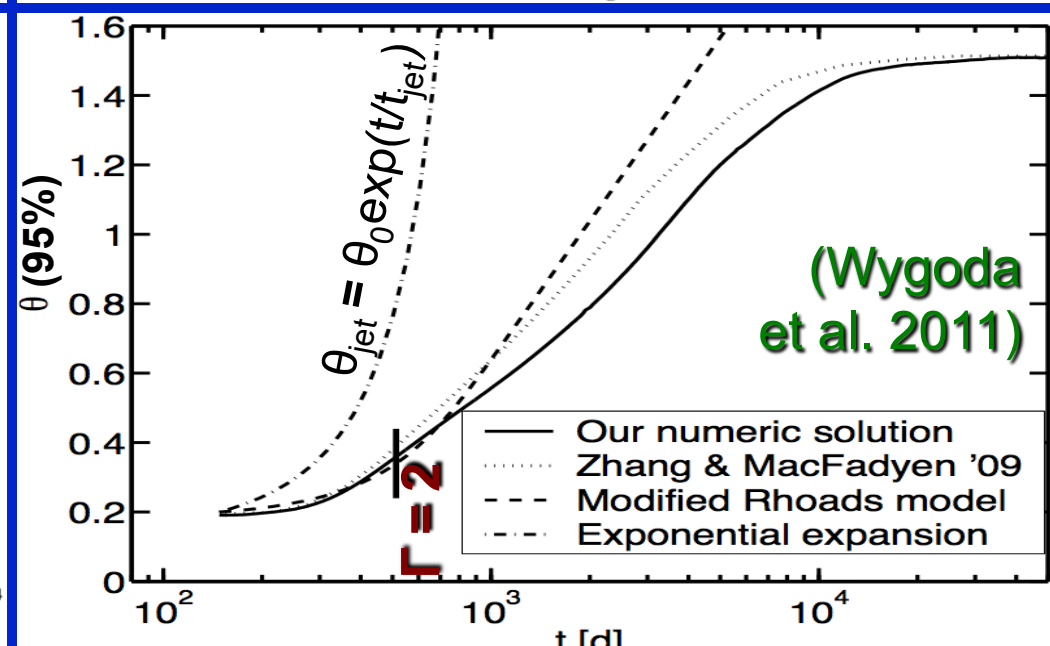
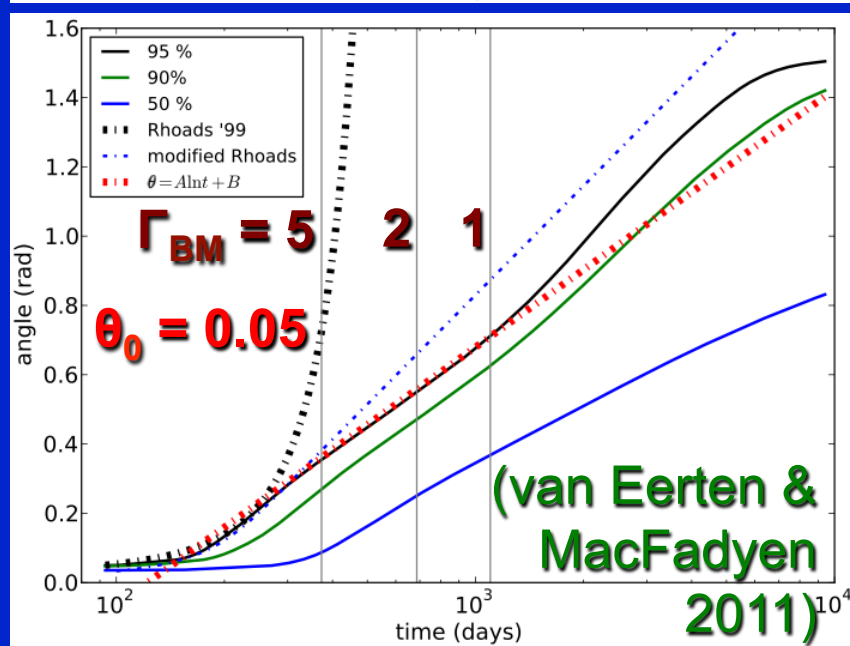
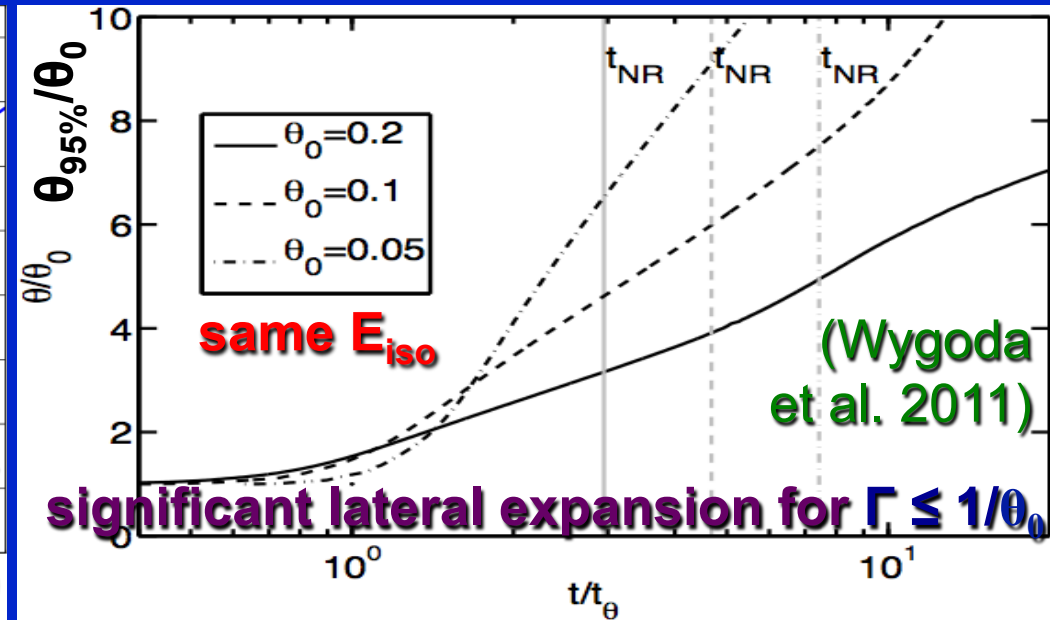
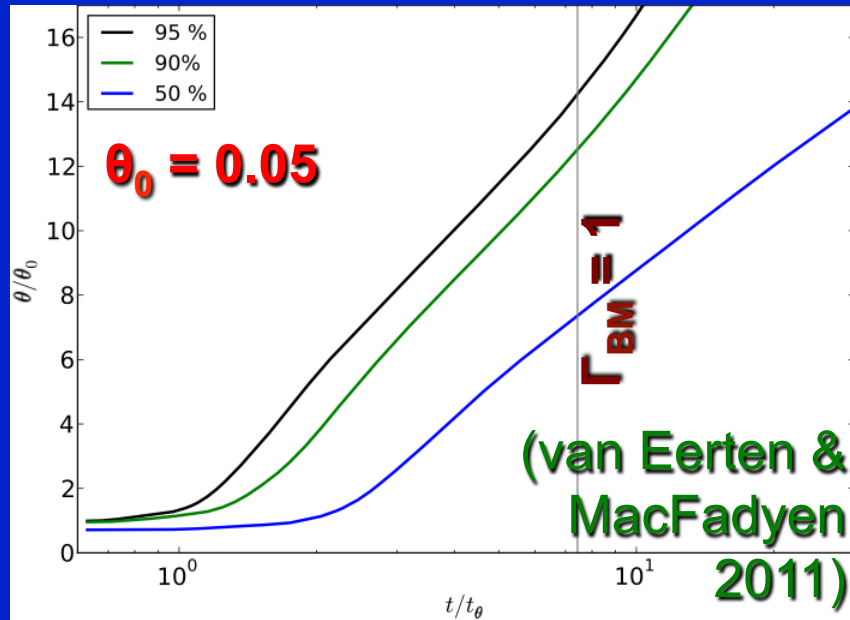
# Analytic vs. Numerical results: a problem?

- **Analytic results** (Rhoads 1997, 99; Sari, Piran & Halpern 99):  
**exponential lateral expansion** at  $R > R_{\text{jet}}$  e.g.  
 $\Gamma \sim (c_s/c\theta_0)\exp(-R/R_{\text{jet}})$ ,  $\theta_{\text{jet}} \sim \theta_0(R_{\text{jet}}/R)\exp(R/R_{\text{jet}})$ 
  - ◆ Supported by a self-similar solution (Gruvinov 2007)
- **Hydro-simulations:** very **mild** lateral expansion while jet is relativistic (also for simplified **2D  $\rightarrow$  1D**)



Modest  $\theta_0$   
 $\Rightarrow$  small region of validity

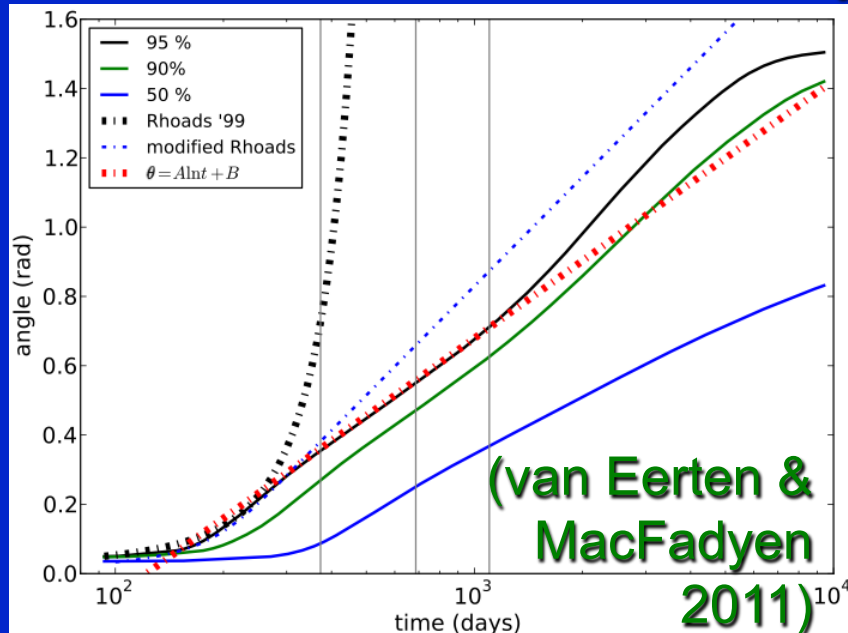
# Analytic vs. Numerical results: a problem?



# Analytic vs. Numerical results: a problem?

van Eerten & MacFadyen 11'

- No exponential lateral expansion even for  $\theta_0 = 0.05$
- Lateral expansion is instead only logarithmic:  $\theta_j \sim \theta_0 \ln(t/t_j)$
- Affects jet break shape +  $t_j$  & late time radio calorimetry



Lyutikov 2011

- Lateral expansion becomes significant only for  $\Gamma \leq \theta_0^{-1/2}$
- Based on thin shell approx.

$$\tan \alpha = -\frac{\partial \ln R}{\partial \theta} \quad (\text{Kumar \& JG 2003})$$

$$\Rightarrow \beta_\theta \sim \frac{1}{\Gamma^2 \Delta \theta} \sim \frac{1}{\Gamma^2 \theta_j}$$

$r = R(\theta) \rightarrow$  shock radius  
in spherical coordinates

$\alpha =$  angle between the shock normal  $\hat{n}$  and radial direction  $\hat{r}$

# Generalized Analytic model (JG & Piran 2012)

## ■ Lateral expansion:

1. new recipe:  $\beta_\theta/\beta_r \sim 1/(\Gamma^2\Delta\theta) \sim 1/(\Gamma^2\theta_j)$  (based on  $\hat{\beta} = \hat{n}$ )

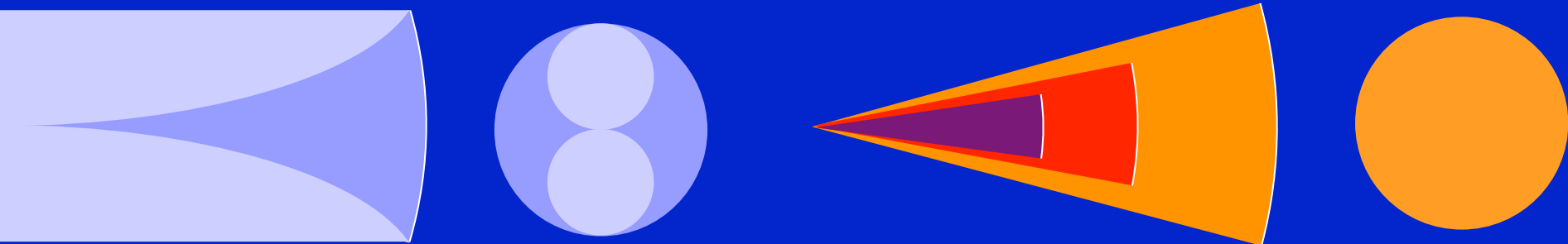
2. old recipe:  $\beta_\theta = u_\theta/\Gamma = u'_\theta/\Gamma \sim \beta_r/\Gamma$  (based on  $u'_\theta \sim 1$ )

Generalized recipe:  $\frac{d\theta_j}{d\ln R} = \frac{\beta_\theta}{\beta_r} \approx \frac{1}{\Gamma^{1+a}\theta_j^a}$ ,  $a = \begin{cases} 1 & (\hat{\beta} = \hat{n}) \\ 0 & (u'_\theta \sim 1) \end{cases}$

◆ New recipe: lower  $\beta_\theta$  for  $\Gamma > 1/\theta_0$  but higher  $\beta_\theta$  for  $\Gamma < 1/\theta_0$

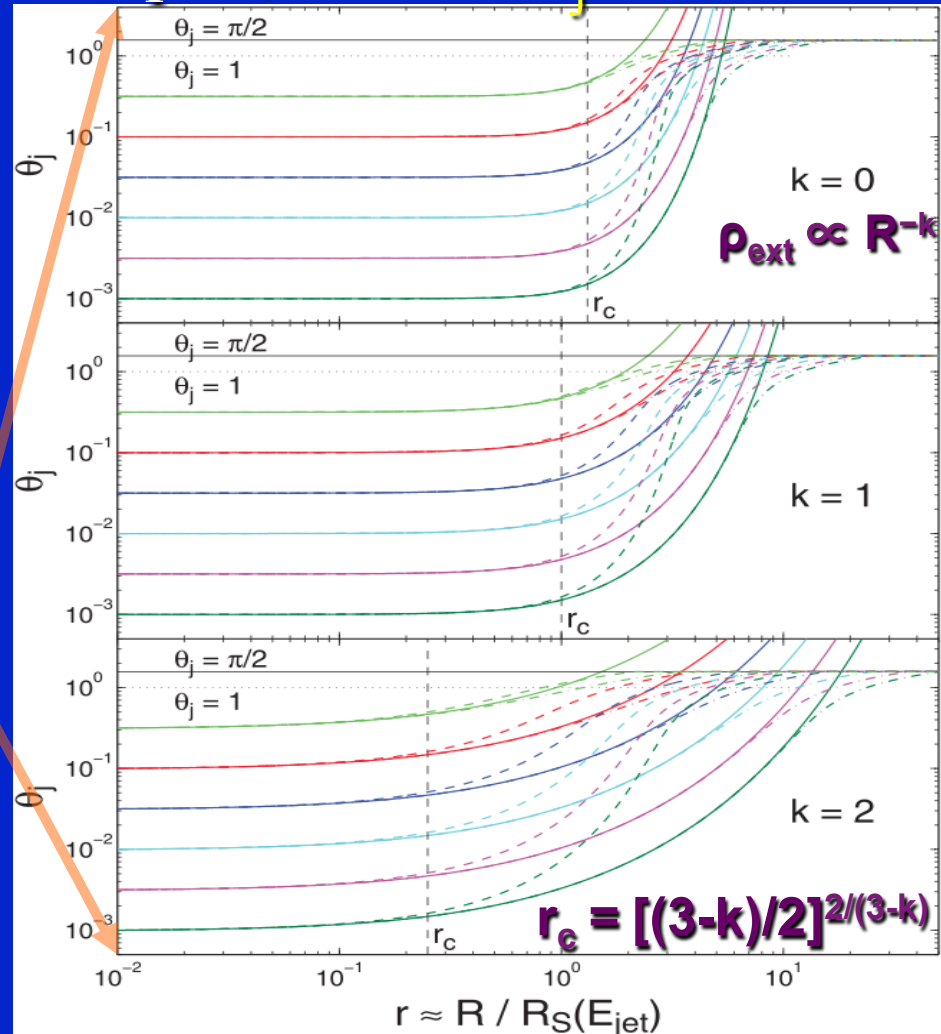
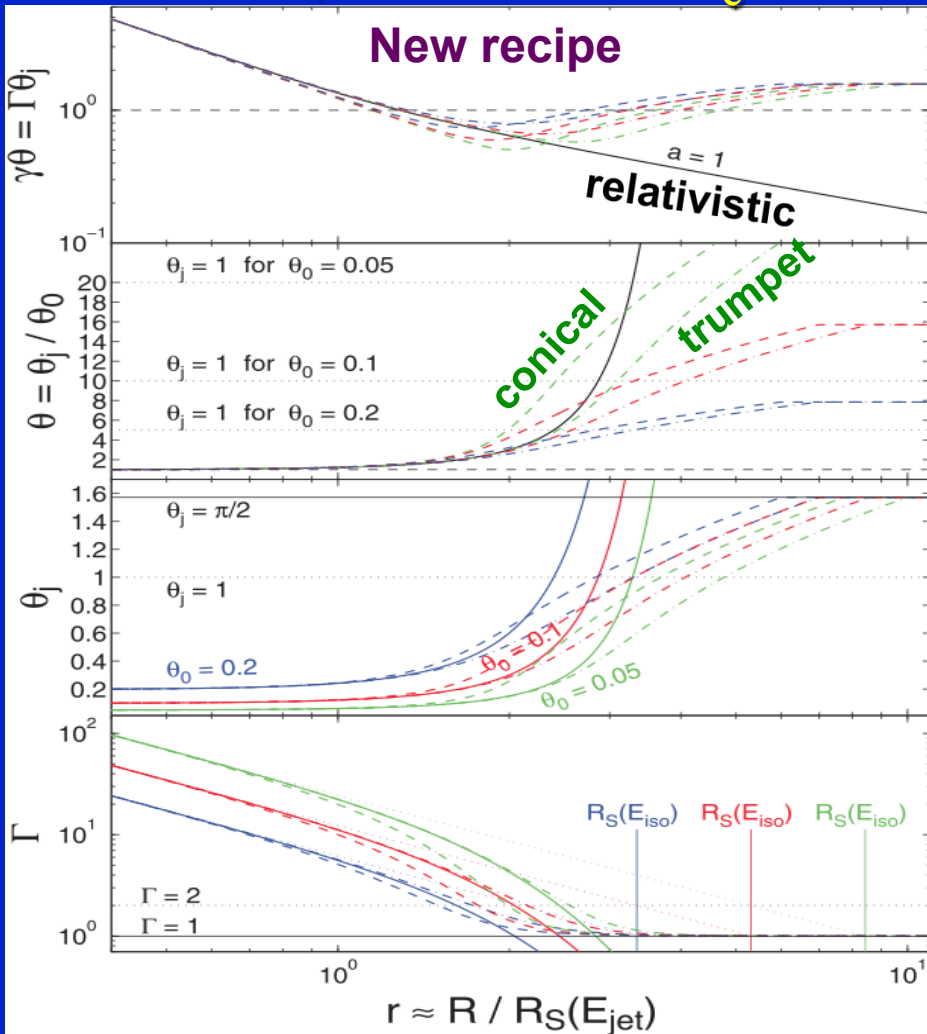
■ Does not assume  $\Gamma \gg 1$  or  $\theta_j \ll 1$  (& variable:  $\Gamma \rightarrow u = \Gamma\beta$ )

■ Sweeping-up external medium: trumpet vs. conical models



# Generalized Analytic model (JG & Piran 2012)

- Main effect of relaxing the  $\Gamma \gg 1, \theta_j \ll 1$  approximation: quasi-logarithmic (~~exponential~~) lateral expansion for  $\theta_0 \gtrsim 0.05$
- conical  $\neq$  rel. for  $r \gtrsim r_c$  while trumpet  $\neq$  rel. for  $\theta_j \gtrsim 0.2$



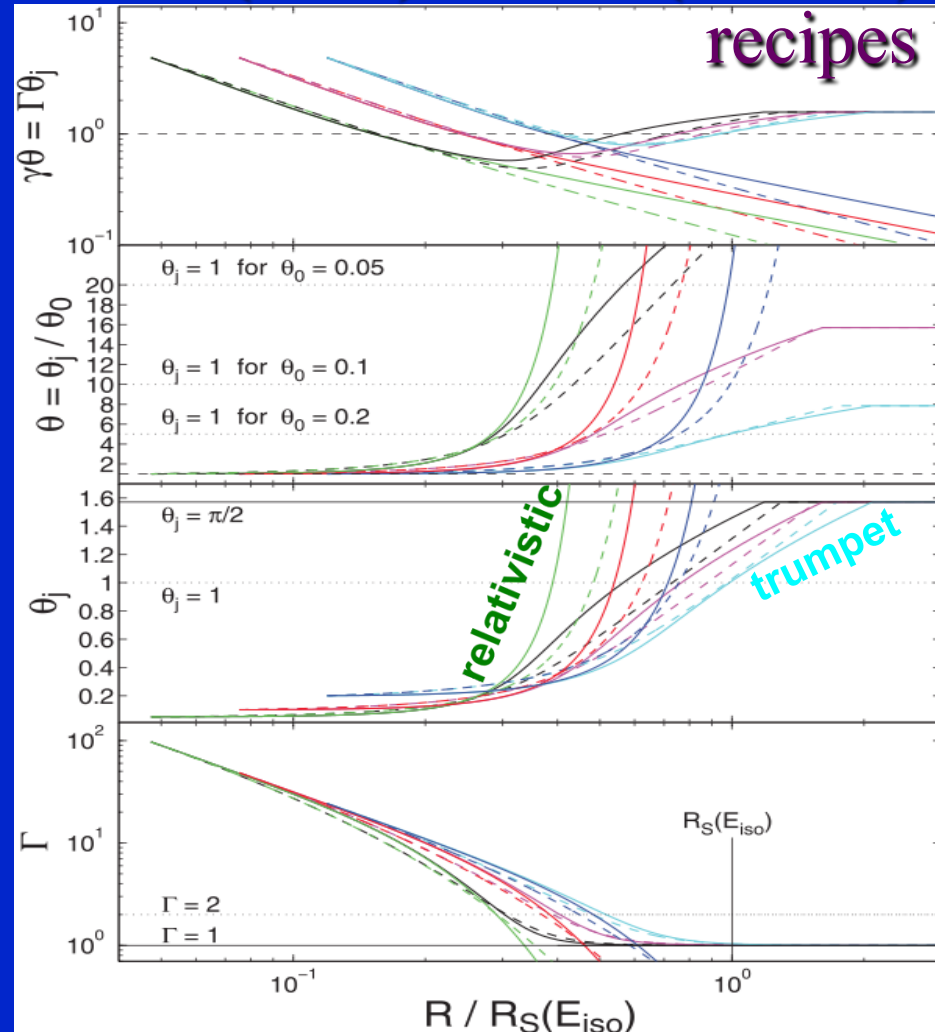
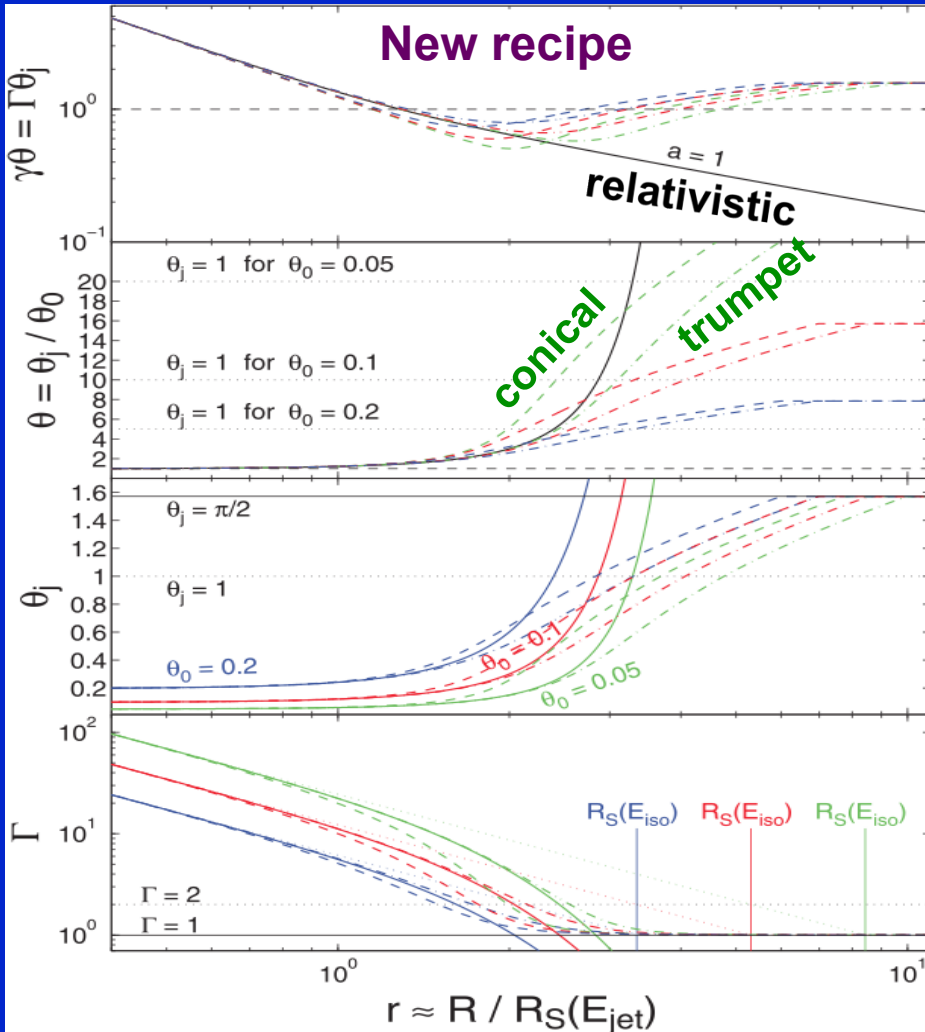


# Generalized Analytic model (JG & Piran 2012)

Conical: larger  $M(R)$  than trumpet  
 $\rightarrow$  lower  $\Gamma(R) \rightarrow$  larger  $\theta_j(R)$

New recipe: lower  $\beta_\theta$  for  $\Gamma > 1/\theta_0$   
 but higher  $\beta_\theta$  for  $\Gamma < 1/\theta_0$

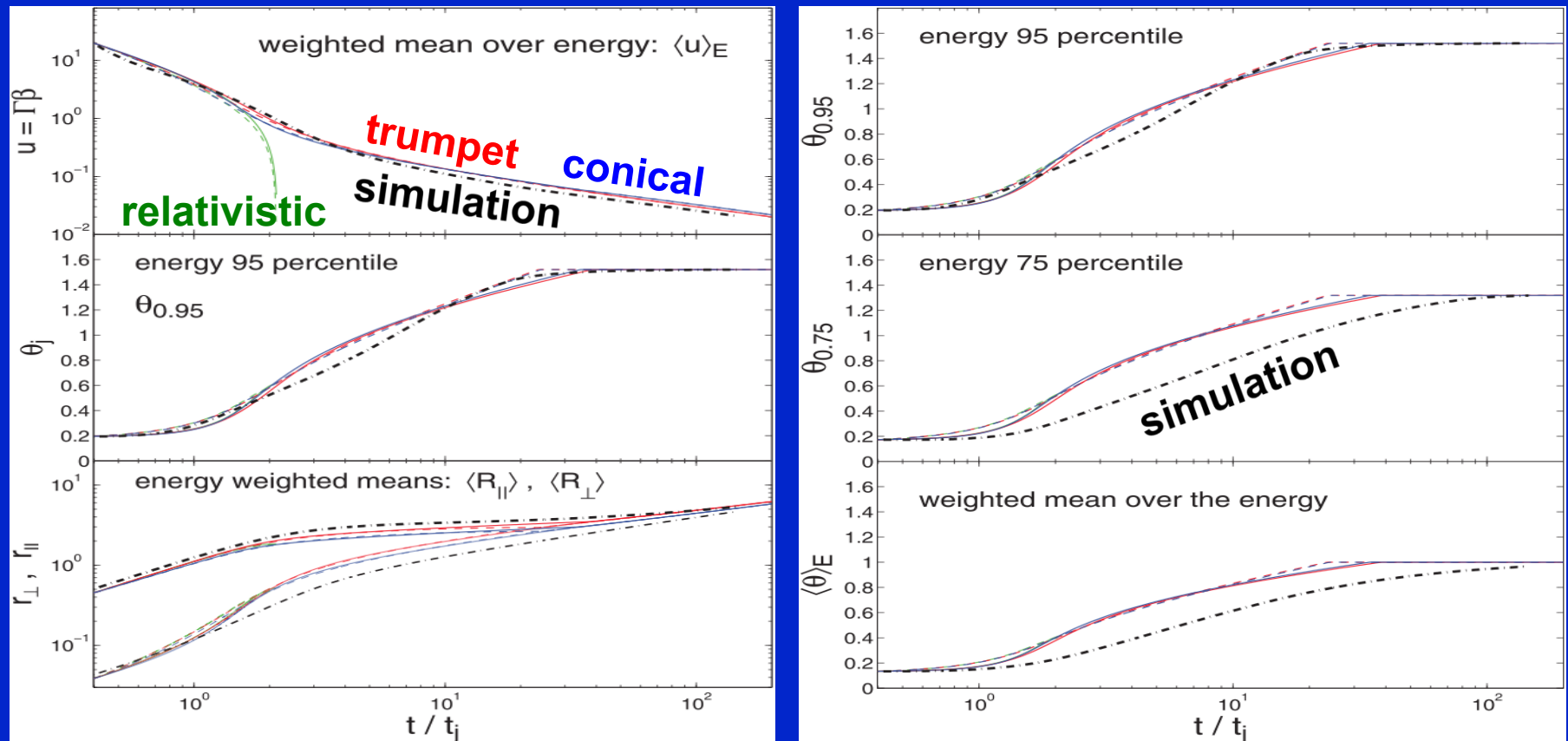
New (*solid*) vs. old (*dashed*)



# Comparison to Simulations (JG & Piran 2012)

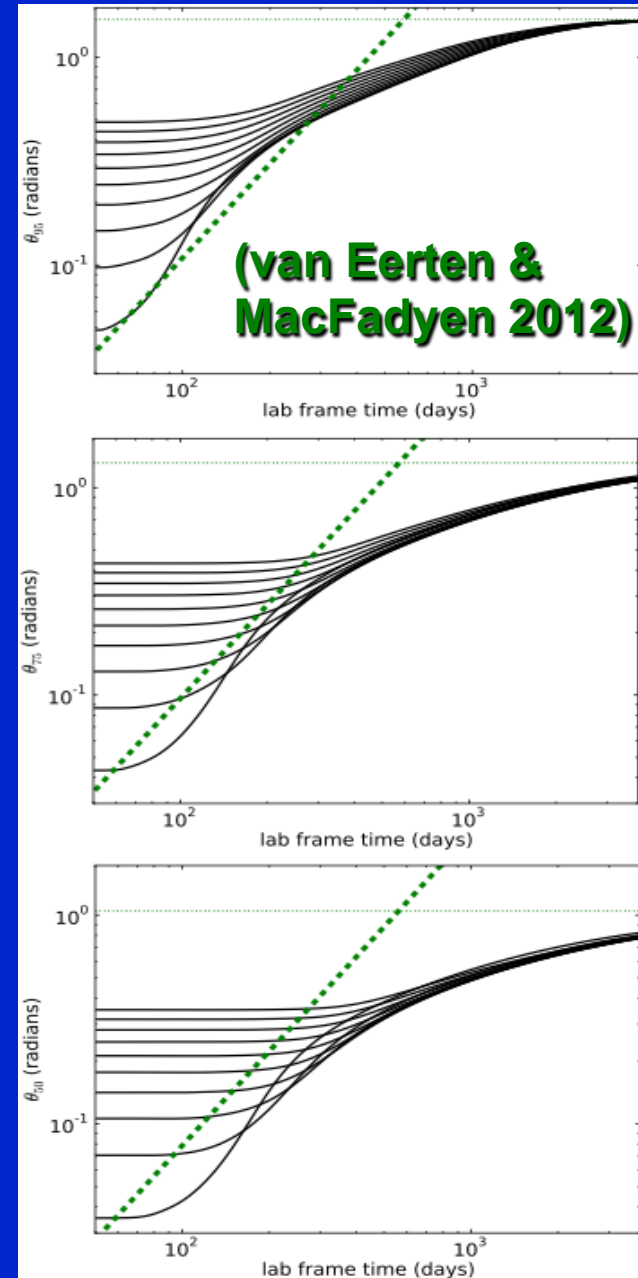
- There is a reasonable overall agreement between the analytic generalized models and the hydro-simulations
- Analytic models: over-simplified, but capture the essence

2D hydro-simulation by F. De Colle et al. 2012, with  $\theta_0 = 0.2$ ,  $k = 0$



# Jet Dynamics: Intermediate Conclusions

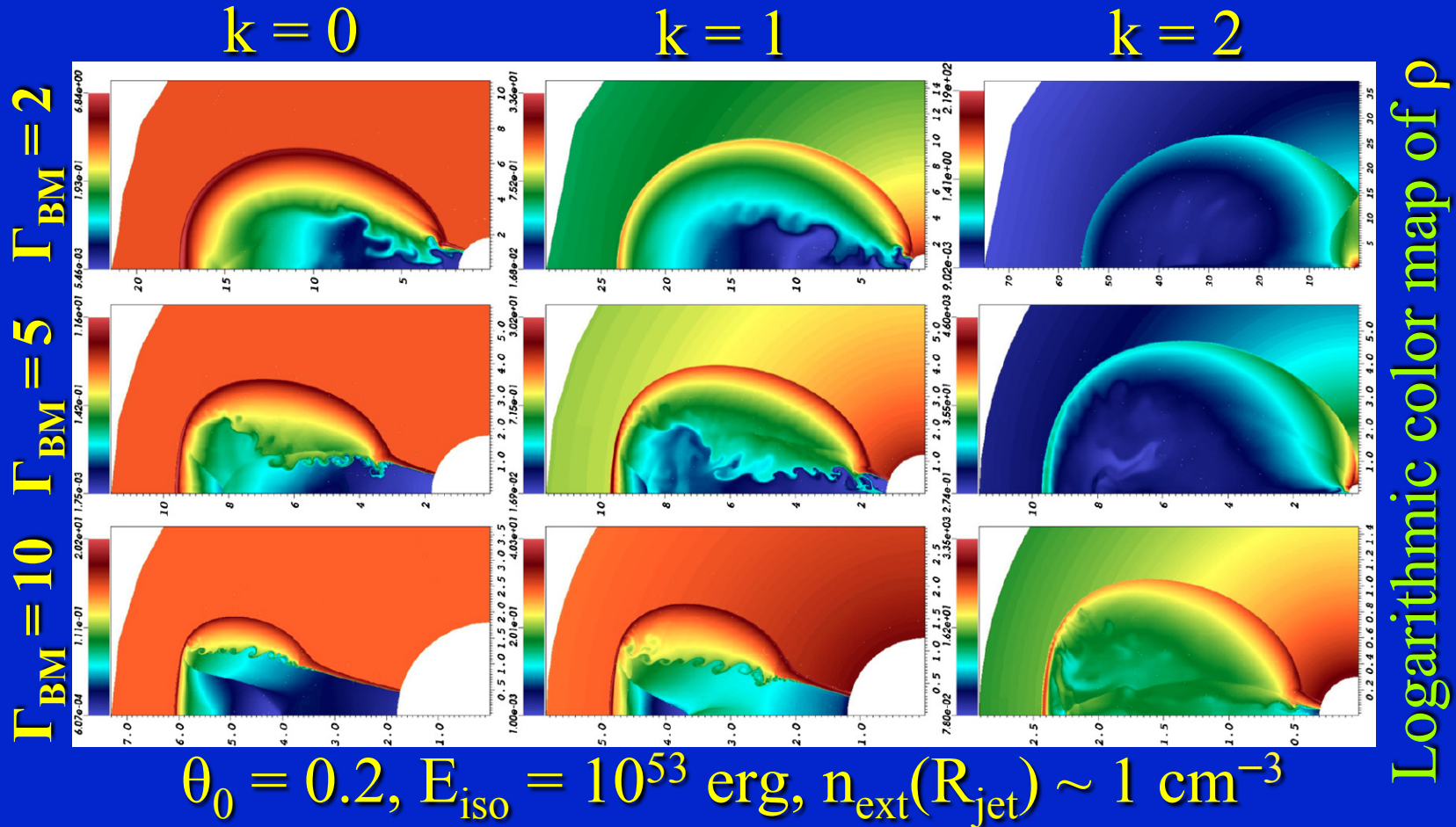
- For  $\theta_0 \gtrsim 0.05$  the lateral expansion is quasi-logarithmic (~~exponential~~), due to the small dynamical range  $1/\theta_0 > \Gamma \gg 1$
- For  $\theta_0 \ll 0.05$  there is an exponential lateral expansion phase (hinted also by van Eerten & MacFadyen's simulations) but such narrow GRB jets appear rare
- The jet first becomes sub-relativistic & only then gradually approaches spherical symmetry over a long time



# Afterglow jet in stratified external media

(De Colle, Ramirez-Ruiz, JG & Lopez-Camara 2012)

- Previous simulations were all for  $k = 0$  where  $\rho_{\text{ext}} \propto R^{-k}$
- Larger (e.g.  $k = 1, 2$ ) are motivated by the stellar wind of a massive star progenitor for long GRBs

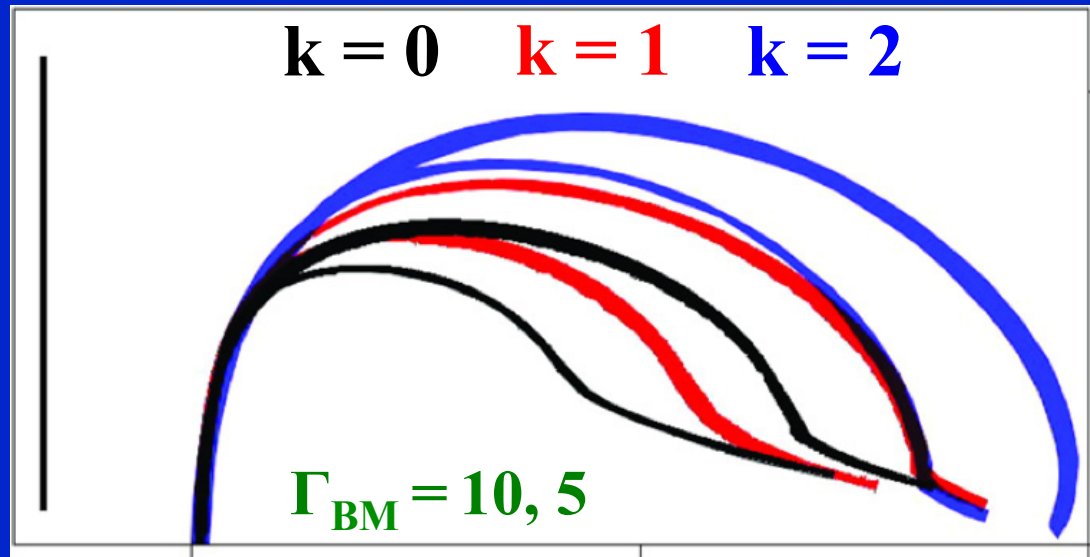




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- Previous simulations were all for  $k = 0$  where  $\rho_{\text{ext}} \propto R^{-k}$
- Larger (e.g.  $k = 1, 2$ ) are motivated by the stellar wind of a massive star progenitor for long GRBs



- At the same Lorentz factor larger  $k$  show larger sideways expansion since they sweep up mass and decelerate more slowly (e.g.  $M \propto R^{3-k}$ ,  $\Gamma \propto R^{(3-k)/2}$  in the spherical case) and spend more time at lower  $\Gamma$  (and  $\beta_{\theta}$  decreases with  $\Gamma$ )



# Afterglow jet in stratified external media

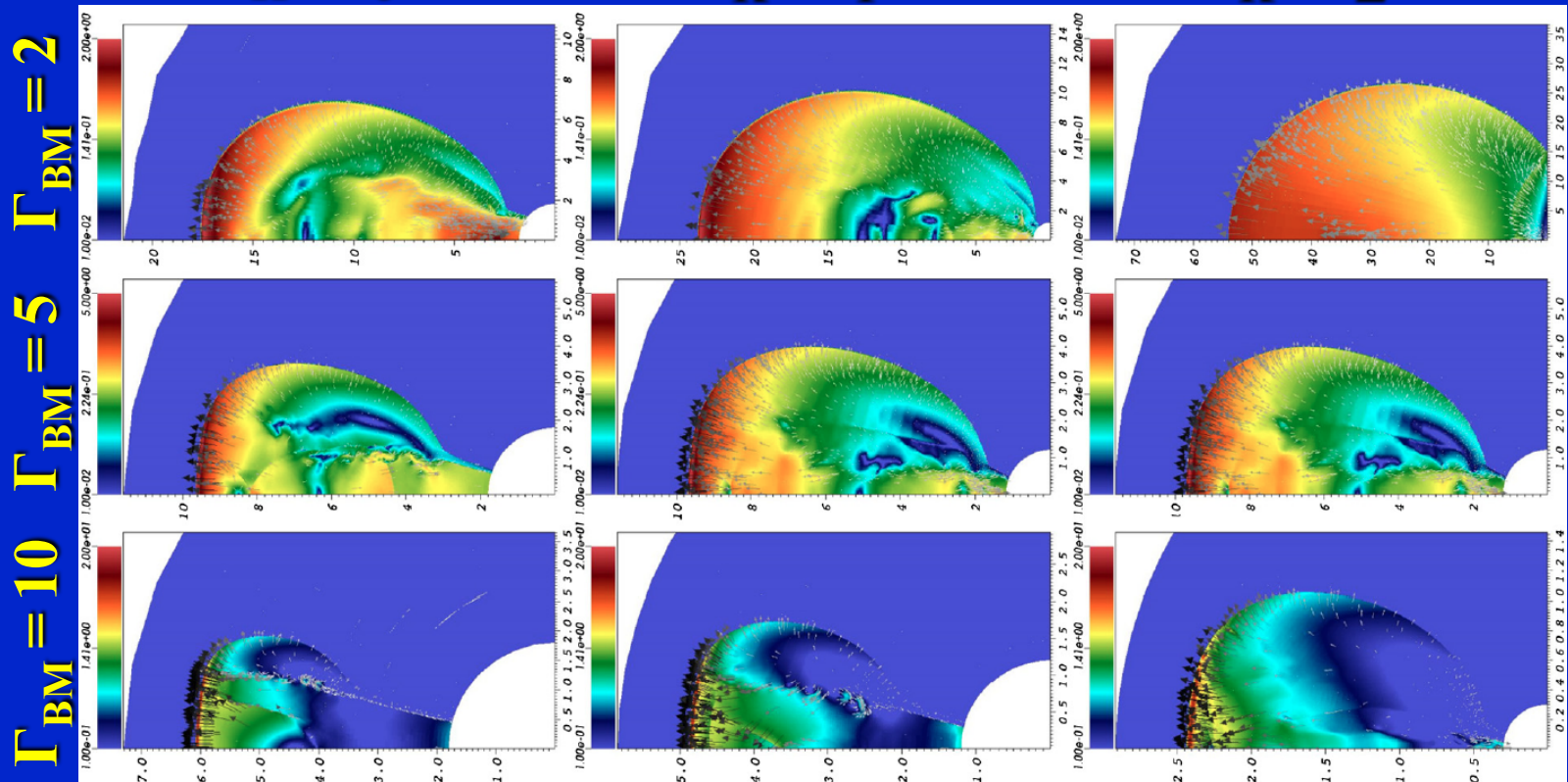
(De Colle, Ramirez-Ruiz, JG & Lopez-Camara 2012)

- The velocity just behind the shock is always normal to the shock front – radial near the head of the jet, while pointing sideways & non-relativistic at the sides of the jet

$k = 0$

$k = 1$

$k = 2$

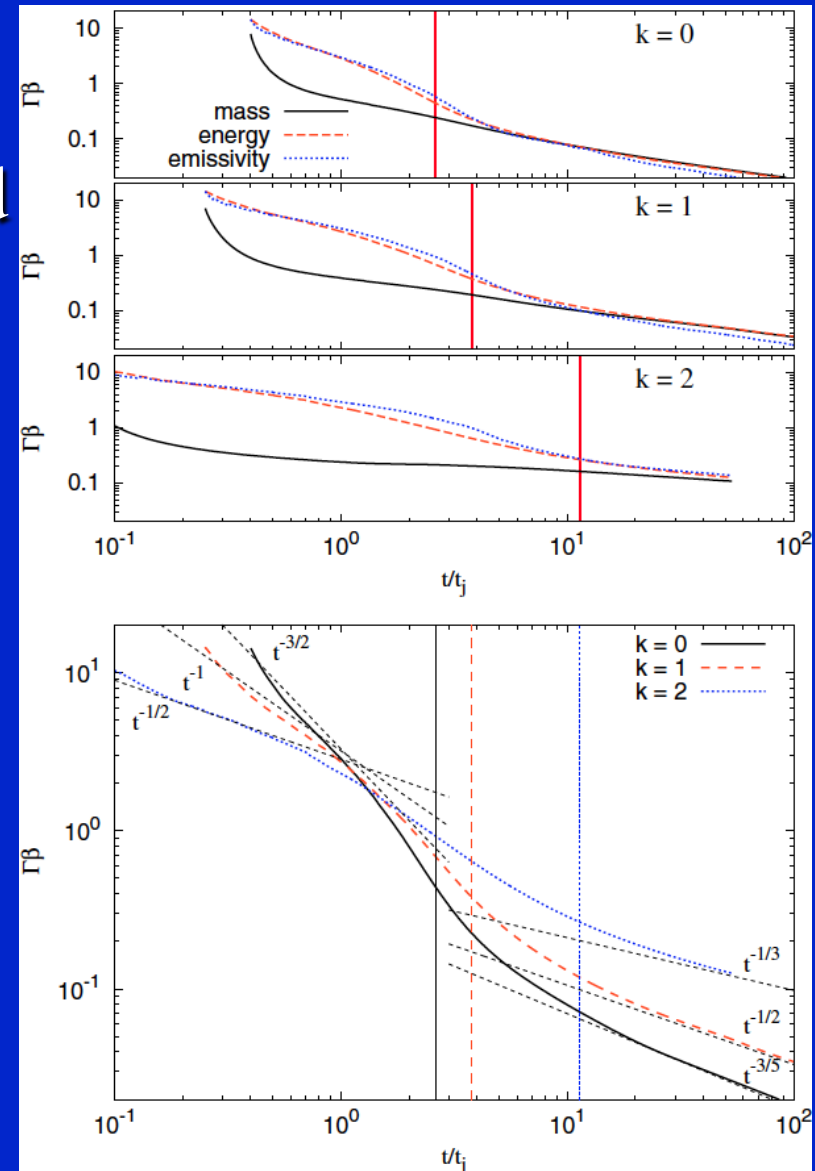
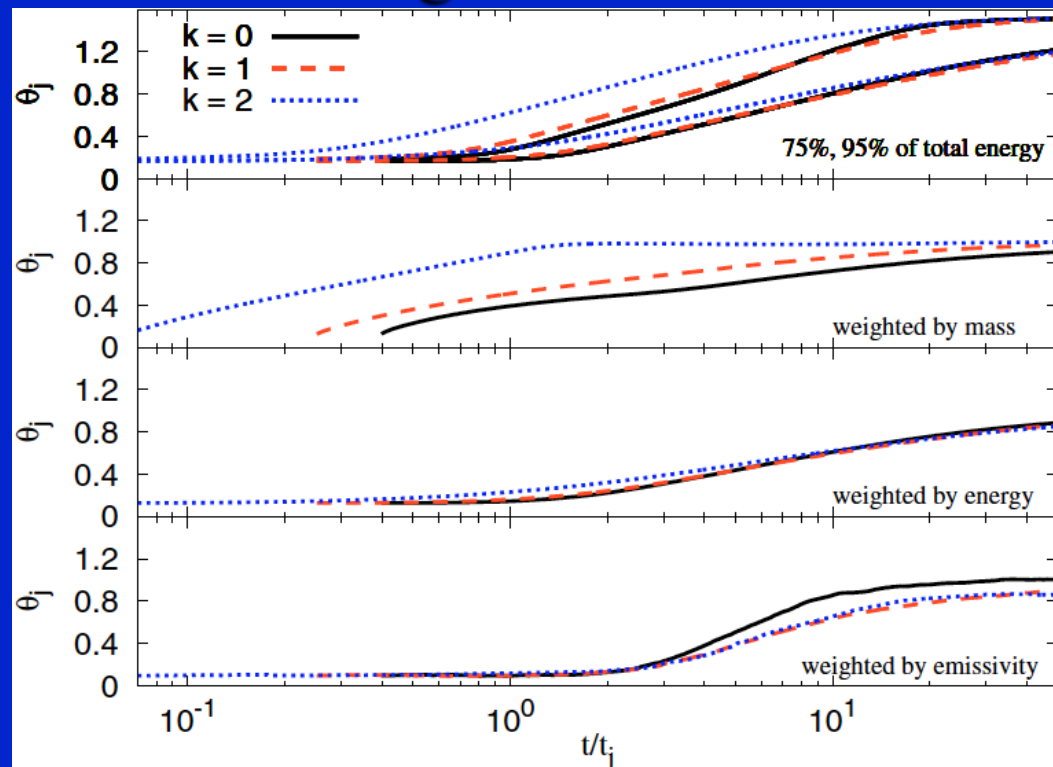


Velocity field ( $\beta$ ) on top of logarithmic color map of  $\rho$

# Afterglow jet in stratified external media

(De Colle, Ramirez-Ruiz, JG & Lopez-Camara 2012)

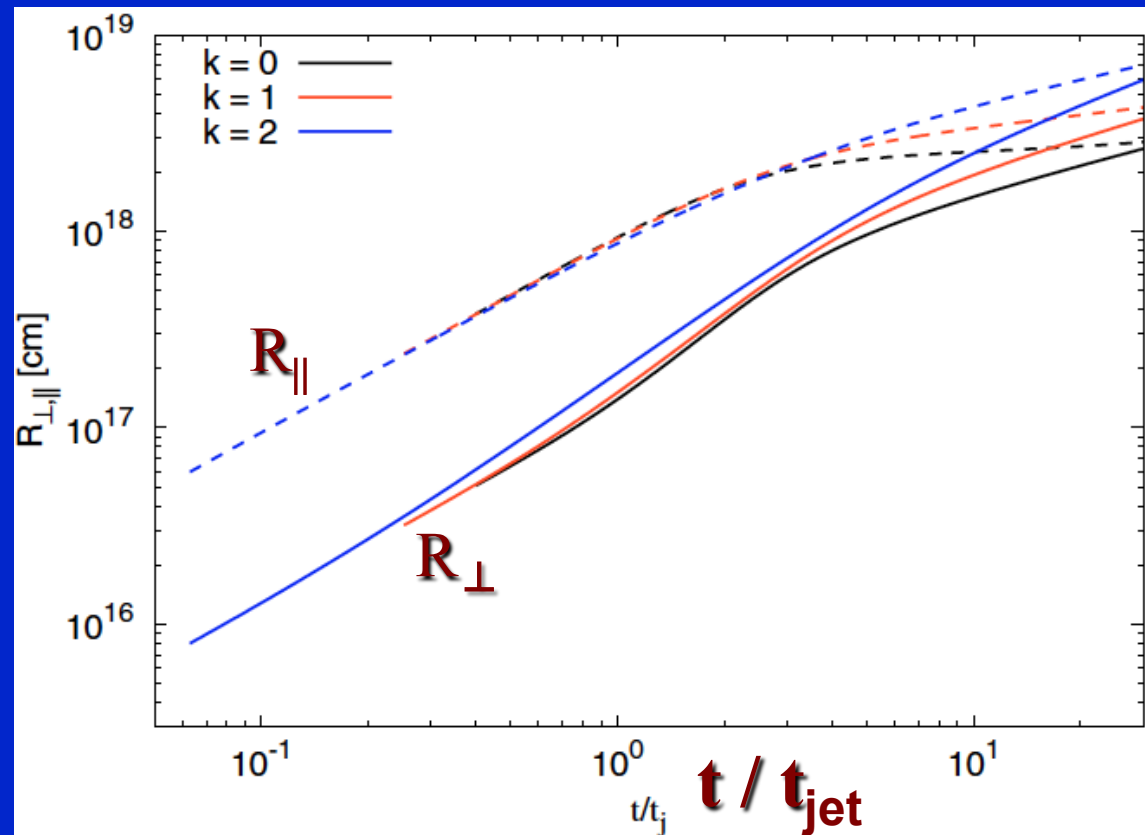
- Swept-up mass: a lot at the sides of the jet at large angles
- Energy, emissivity: near the head
- Spherical symmetry approached later for larger  $k$



# Afterglow jet in stratified external media

(De Colle, Ramirez-Ruiz, JG & Lopez-Camara 2012)

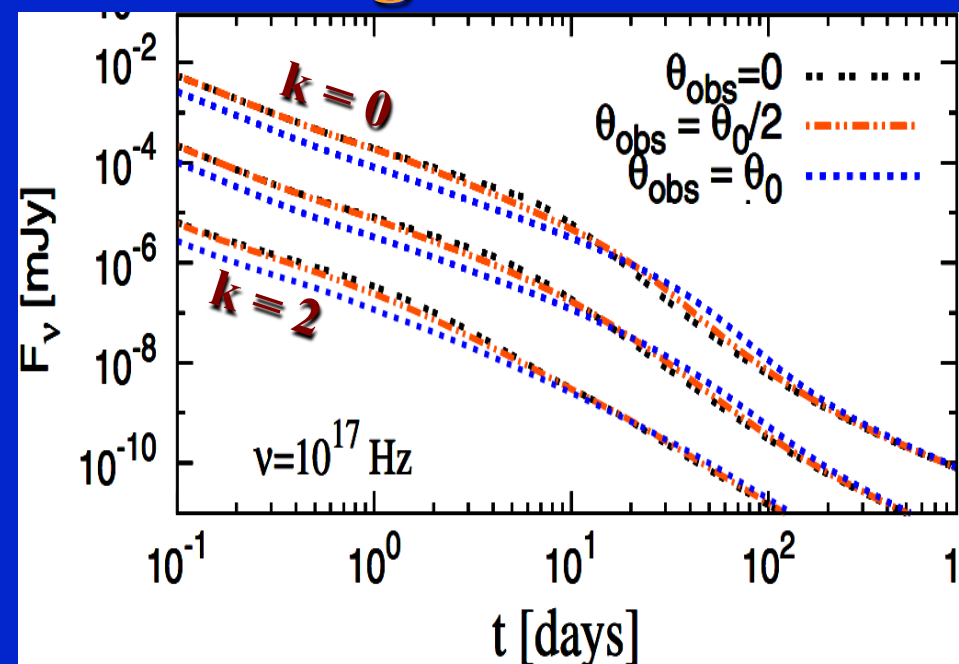
- For  $k = 0$  the growth of  $R_{\parallel}$  is stalled at  $t_{NR}(E_{iso})$  while  $R_{\perp}$  continues to grow  $\rightarrow$  helps approach spherical symmetry
- Less pronounced for larger  $k$  as the slower accumulation of mass enables  $R_{\parallel}$  to grow more  $\rightarrow$  become spherical more slowly



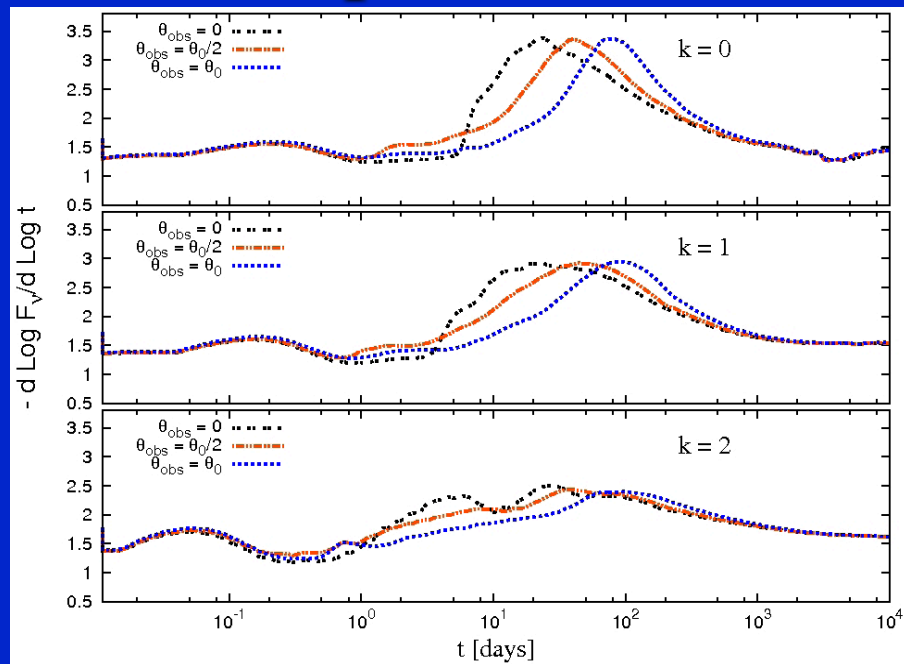
# The shape of the jet break

- Jet break becomes smoother with increasing  $k$  (as expected analytically; Kumar & Panaitescu 2000 – KP00)
- However, the jet break is significantly sharper than found by KP00 → better prospects for detection
- Varying  $\theta_{\text{obs}} < \theta_0$  dominates over varying  $k \lesssim 2$

## Lightcurves



## Temporal index

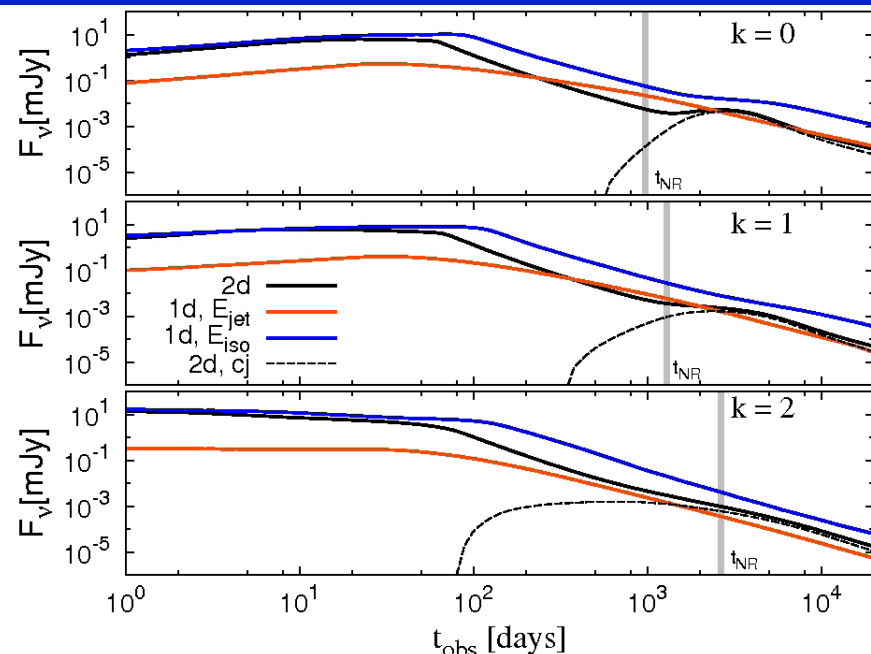




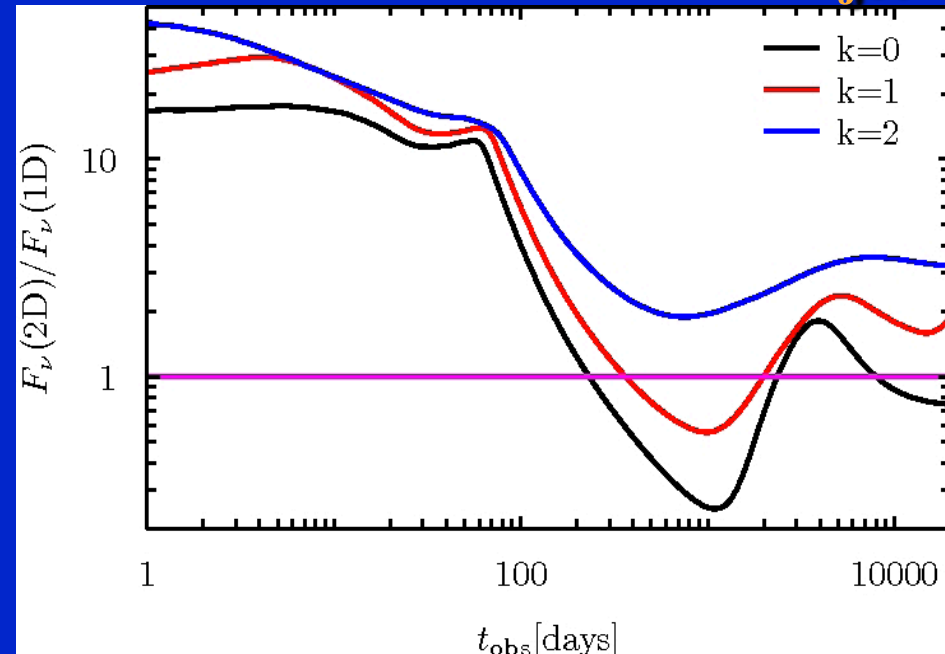
# Late time Radio emission & Calorimetry

- The bump in the lightcurve from the counter jet is much less pronounced for larger  $k$  (as the counter jet decelerates & becomes visible more slowly)  $\rightarrow$  hard to detect
- The error in the estimated energy assuming a spherical flow depends on the observation time  $t_{\text{obs}}$  & on  $k$

## Radio Lightcurves



## Flux Ratio: 2D/1D( $E_{\text{jet}}$ )





# Conclusions:

- **Jet lateral expansion:** analytic models & simulations agree
  - ◆ For  $\theta_0 \gtrsim 0.05$  the lateral expansion is quasi-logarithmic (~~exponential~~), due to small dynamic range  $1/\theta_0 > \Gamma \gg 1$
  - ◆ For  $\theta_0 \ll 0.05$  there is an exponential lateral expansion phase early on (but such narrow GRB jets appear rare)
  - ◆ The jet first becomes sub-relativistic & only then slowly approaches spherical symmetry over a long time
- **Jet in a stratified external medium:**  $\rho_{\text{ext}} \propto R^{-k}$  for  $k = 0, 1, 2$ 
  - ◆ larger  $k$  jets sweep-up mass & slow down more slowly
    - ➔ sideways expansion is faster at  $t < t_j$  & slower at  $t > t_j$
    - ➔ become spherical slower; harder to see counter jet
  - ◆ Jet break is smoother for larger  $k$  but possibly detectable
  - ◆ Jet break sharpness affected more by  $\theta_{\text{obs}} < \theta_0$  than  $k \lesssim 2$
  - ◆ Radio calorimetry accuracy affected both by  $t_{\text{obs}}$  &  $k$