

# The Synchrotron X-Ray Nebula around Swift J1834.9–0846

- Magnetic field in the X-ray emitting region is obtained from

$$B = \left( \frac{L_X \sigma_e}{AV} \frac{\Gamma - 2}{\Gamma - 1.5} \frac{\nu_1^{1.5-\Gamma} - \nu_2^{1.5-\Gamma}}{\nu_m^{2-\Gamma} - \nu_M^{2-\Gamma}} \right)^{2/7}$$

$$\simeq \begin{cases} 4.0 \xi \sigma_e^{2/7} d_4^{-2/7} \mu\text{G} & \text{(whole nebula) ,} \\ 5.0 \xi_{\text{in}} \sigma_e^{2/7} d_4^{-2/7} \mu\text{G} & \text{(inner nebula) .} \end{cases}$$

Assuming:

- Power-law electron distribution:  $\frac{dN_e}{d\gamma_e} \propto \gamma_e^{-p} \quad \gamma_1 < \gamma_e < \gamma_2$
- Electrons emitting in observed range:  $\gamma_1 < \gamma_m < \gamma_e < \gamma_M < \gamma_2$

- Magnetization:  $\sigma = \frac{B^2}{4\pi w} = \frac{3}{2} \frac{B^2}{8\pi e} = \frac{3}{2} \frac{E_B}{E_m} = \frac{3}{2} \sigma_e \epsilon_e$

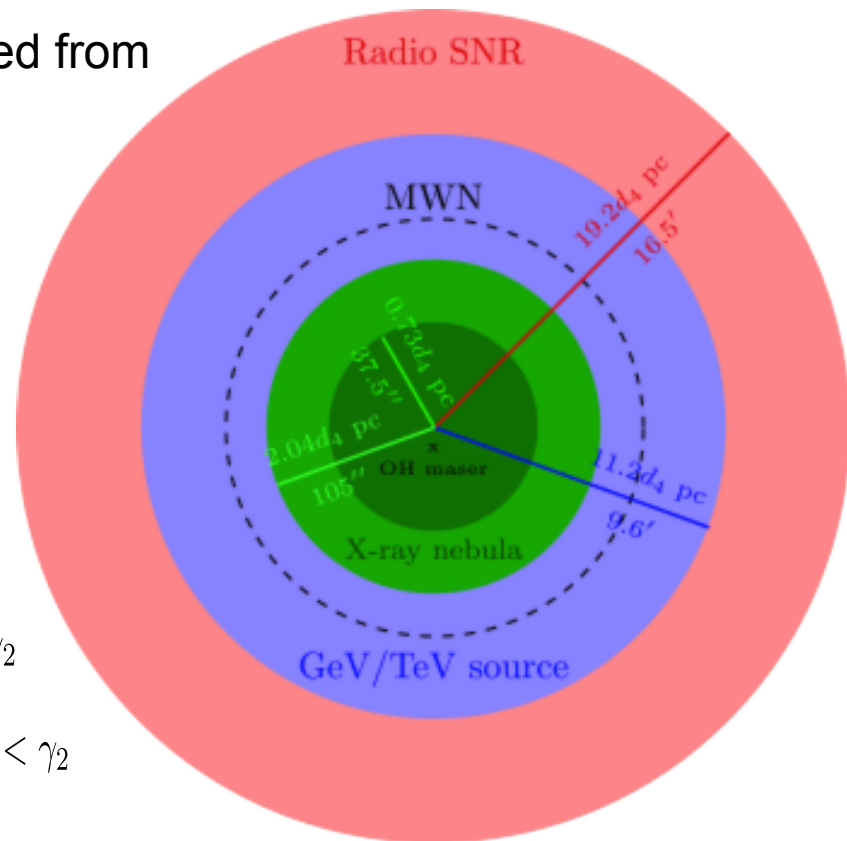
- Emission volume:  $V = \frac{4}{3} \pi R_x^3$

$$\sigma_e \equiv E_B / E_e$$

$$E_e = \epsilon_e E_m$$

( $E_m$  = total energy in matter in the emission region)

$$\xi^{7/2}(\nu_1, \nu_2) = \left( \frac{\nu_2^{1.5-\Gamma} - \nu_1^{1.5-\Gamma}}{\nu_M^{1.5-\Gamma} - \nu_m^{1.5-\Gamma}} \right) = \text{ratio of energy in all electrons to that in those radiating in the observed frequency range}$$



Total emitted power (0.5 – 30 keV)

$$L_{\text{x,tot}} = 2.74 \times 10^{33} d_4^2 \text{ erg s}^{-1}$$

The spin-down power

$$L_{\text{sd}} = 2.05 \times 10^{34} \text{ erg s}^{-1}$$

X-ray efficiency of MWN

$$\eta_X = \frac{L_{\text{X,tot}}}{L_{\text{sd}}} = 0.13 d_4^2$$

# Constraints From Maximum Electron Lorentz Factor

Maximum injected L.F.:  $\gamma_{\max} = \frac{eR_{\text{NS}}^3 \Omega^2 B_s}{m_e c^4} = \frac{e}{m_e c^2} \sqrt{\frac{L_{\text{sd}}}{fc}} = 4.9 \times 10^8 f^{-1/2}$   
 (corresponding to acc. in the potential difference across the open field lines)

How robust is this limit?

Max. obs. energy:

$$E_X = 30 E_{M,30} \text{ keV}$$

since  $\gamma_{\max} > \gamma_M$

$$B > B_{\min} \equiv \frac{m_e^3 c^6 f E_X}{\hbar e^3 L_{\text{sd}}} \simeq 11.0 f E_{M,30} \mu\text{G}$$

which further yields

$$\left. \begin{aligned} \sigma_e &> 0.043 d_4 \left( \frac{f E_{M,30}}{\xi_7} \right)^{7/2}, \\ \sigma_{e,\text{in}} &> 0.055 d_4 \left( \frac{f E_{M,30}}{\xi_{\text{in}}/5} \right)^{7/2}. \end{aligned} \right| \begin{aligned} \xi &> 16.1 f E_{M,30} \left( \frac{\epsilon_e d_4}{\sigma_{-2.5}} \right)^{2/7}, \\ \xi_{\text{in}} &> 12.7 f E_{M,30} \left( \frac{\epsilon_e d_4}{\sigma_{\text{in},-2.5}} \right)^{2/7}, \end{aligned}$$

Synchrotron **cooling time** of X-ray emitting electrons:  $t_{\text{syn}} = \frac{6\pi m_e c}{\sigma_T B^2 \gamma_e} \simeq 1.02 B_{15\mu\text{G}}^{-3/2} E_2^{-1/2} \text{ kyr}$ .  
 (at  $2E_2$  keV) is  $\ll$  system's age  $\rightarrow$  quasi-steady state

Electron energy balance:  $\langle \dot{E} \rangle = g L_{\text{sd}} = \frac{(1 + \sigma)}{\epsilon_e \epsilon_X} L_{X,\text{tot}}$

In terms of the observed X-ray efficiency:

$$g = \frac{L_{X,\text{tot}}}{L_{\text{sd}}} \frac{(1 + \sigma) \xi^{7/2}}{\epsilon_e} = \frac{\eta_X (1 + \sigma) \xi^{7/2}}{\epsilon_e}$$

Estimate of  $\xi_{\text{in}}$  from the inner nebula yields:  $\frac{g\sigma}{1 + \sigma} > 3.07 d_4^3 E_{M,30}^{7/2} f^{7/2}$

$(1 + \sigma)^{-1}$  = fraction of the total energy injected into the nebula going to particles (the rest goes into the B-field)

$\epsilon_e$  = fraction of that going into power-law energy distribution of electrons

$\epsilon_X = \xi^{-7/2}$  = fraction of that going into electrons radiating observed X-rays

# The g-sigma parameter space & Implications

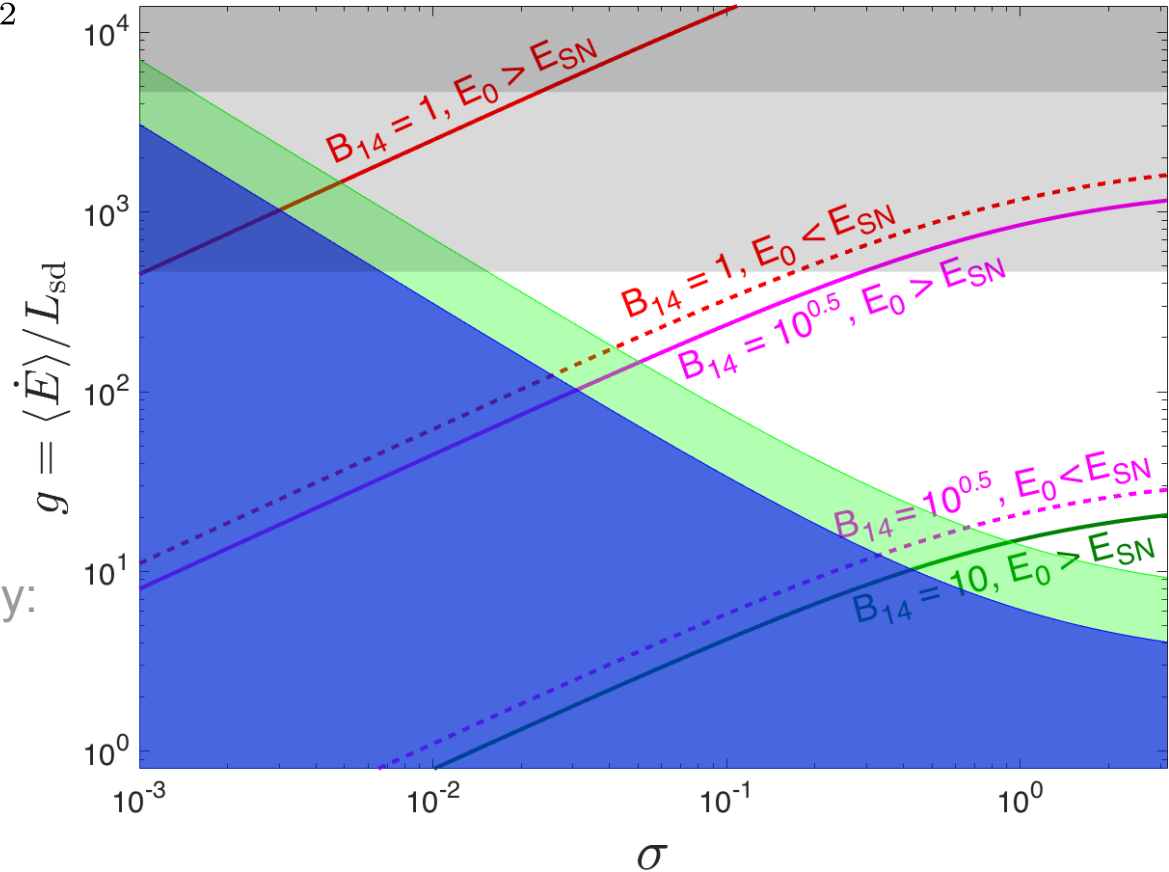
Electron energy balance:  $\frac{g\sigma}{1+\sigma} > 3.07 d_4^3 E_{M,30}^{7/2} f^{7/2}$

Equating  $B_{in}(L_x) = B_{in}(\text{dynamics})$  yields:

$$g = \begin{cases} \frac{1422 E_{\text{tot},52.3}^{\frac{7}{8}} \left(\frac{\kappa_1}{6.0f}\right)^{\frac{7}{4}}}{B_{14.5}^{\frac{7}{2}} d_4^4 M_3^{\frac{7}{8}}} \left(\frac{\sigma}{1+\sigma}\right)^{\frac{3}{4}}, & E_0 > E_{\text{SN}} \\ \frac{35.03 E_{\text{tot},51}^{\frac{7}{8}} \left(\frac{\kappa_2}{4.47f}\right)^{\frac{7}{4}}}{B_{14.5}^{\frac{7}{2}} E_{\text{SN},51}^{\frac{7}{20}} d_4^4 M_3^{\frac{7}{8}} M_3^{\frac{7}{8}} P_{0,-2}^{\frac{7}{10}}} \left(\frac{\sigma}{1+\sigma}\right)^{\frac{3}{4}}, & E_0 < E_{\text{SN}} \end{cases}$$

Limit on energy source: internal field decay:

$$g < \frac{2}{2+\alpha} \frac{R_{\text{NS}}^3 B_{\text{int,max}}^2}{6L_{\text{sd}}t} \simeq 4.65 \times 10^2 B_{M,16}^2 t_{4.5}^{-1} \\ \simeq 4.65 \times 10^3 B_{M,16.5}^2 t_{4.5}^{-1}$$



Dipole field decay isn't enough  $\Rightarrow$  another energy source is required – likely internal-field decay

Quiescent steady particle wind? What would drive it? Seems unlikely

Outflows associated to magnetar bursts – known to occur in **giant flares** (SGR 1806–20, 1900+14)

Burst distribution:  $dN/dE \propto E^{-5/3}$  (self-organized criticality)? **giant flares** dominate for rising  $E^2 dN/dE$

$$\Rightarrow \Delta t_{\text{GF}} = \frac{E}{\langle \dot{E} \rangle} = 100 \left(\frac{g}{50}\right)^{-1} \left(\frac{E}{10^{45.5} \text{ erg}}\right) \text{ yr} \quad \text{since} \quad \langle \dot{E} \rangle \approx \left(\frac{g}{50}\right) \times 10^{36} \text{ erg s}^{-1}$$