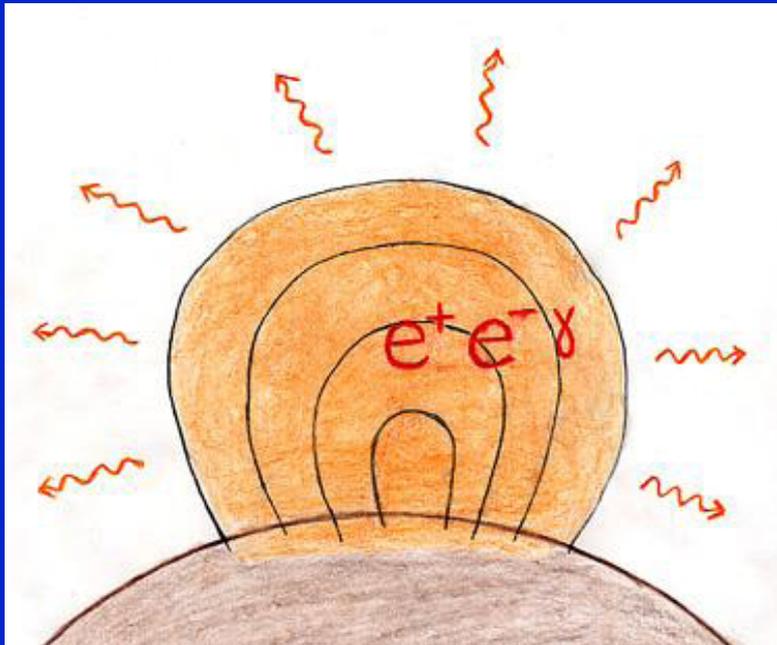


# Magnetic Field Decay in Magnetars & implications for evolutionary links

**Jonathan Granot**

Open University of Israel, Hebrew University

Collaborators: **Simone Dall'Osso, Tsvi Piran**



# Outline of the talk:

- Motivation & brief introduction
- Observational evidence for magnetar B-field decay
- Field decay: main properties (true vs. spin-down age) & observational constraints on dipole field decay
- $L_X > |\dot{E}_{B,dipole}| \rightarrow$  another energy source is required
- $\rightarrow$  Internal field decay? requires  $B_{internal,i} \gtrsim 10^{16} \text{ G}$
- Implications for evolutionary links
- Conclusion

# Motivation

X-ray emission of magnetars is powered by the decay of their super-strong magnetic field:

- a) eventually test this hypothesis;
- b) best objects to study B-field decay

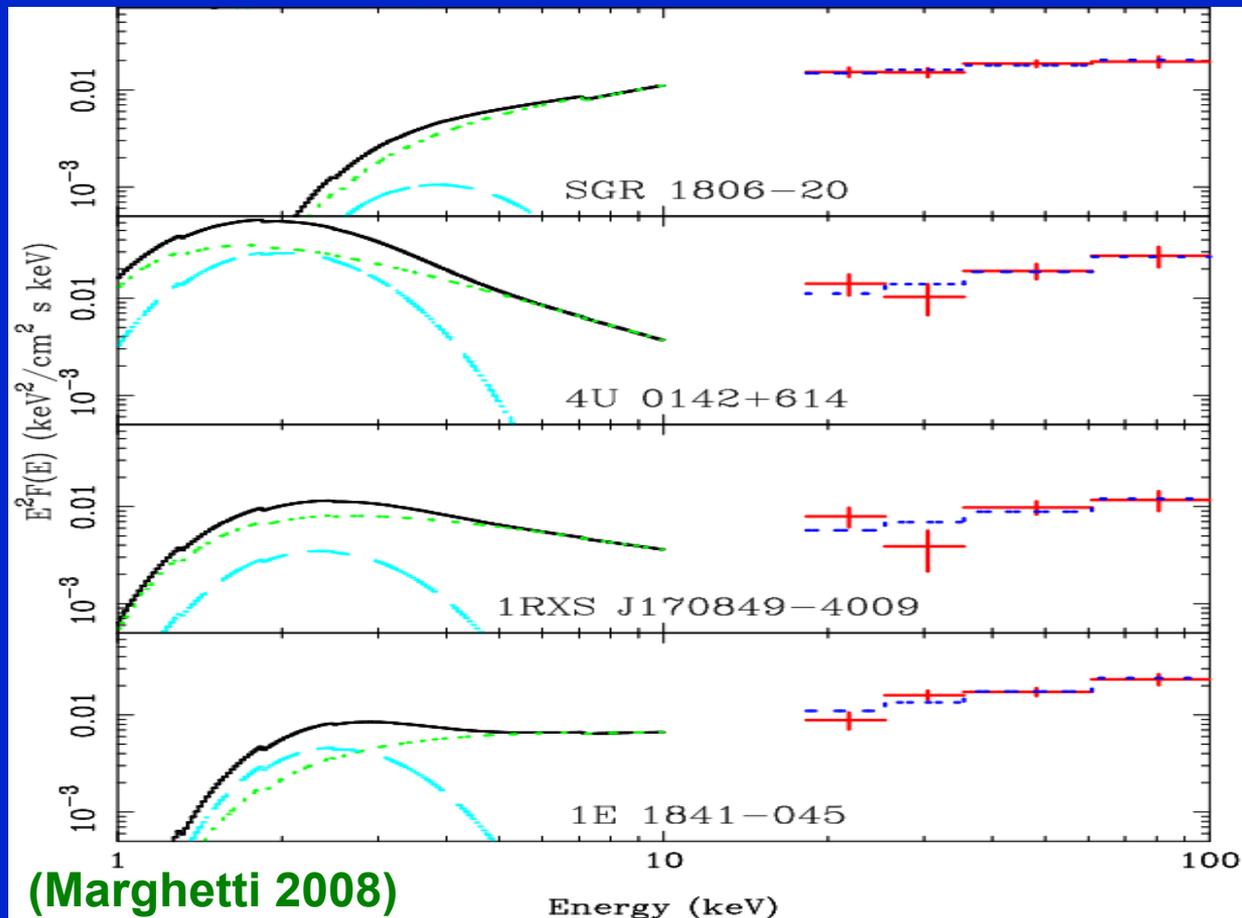
“Magnetar-like” emission from unsuspected magnetars:  
**Additional degree of freedom besides dipole field?**

Link between different classes of high-B NSs  
(SGRs, AXPs, “transient” AXP/SGRs, XDINs,...)

**Dall'Osso, JG & Piran 2012, MNRAS, 422, 2878**

# Source Classes

AXPs/SGRs: Persistent X-ray emission  $\gg dE_{\text{rot}}/dt$   
Thermal  $\rightarrow kT = (0.5-0.7) \text{ keV}$   
Hard-X spectral tails (up to **150 keV**)



# Source Classes

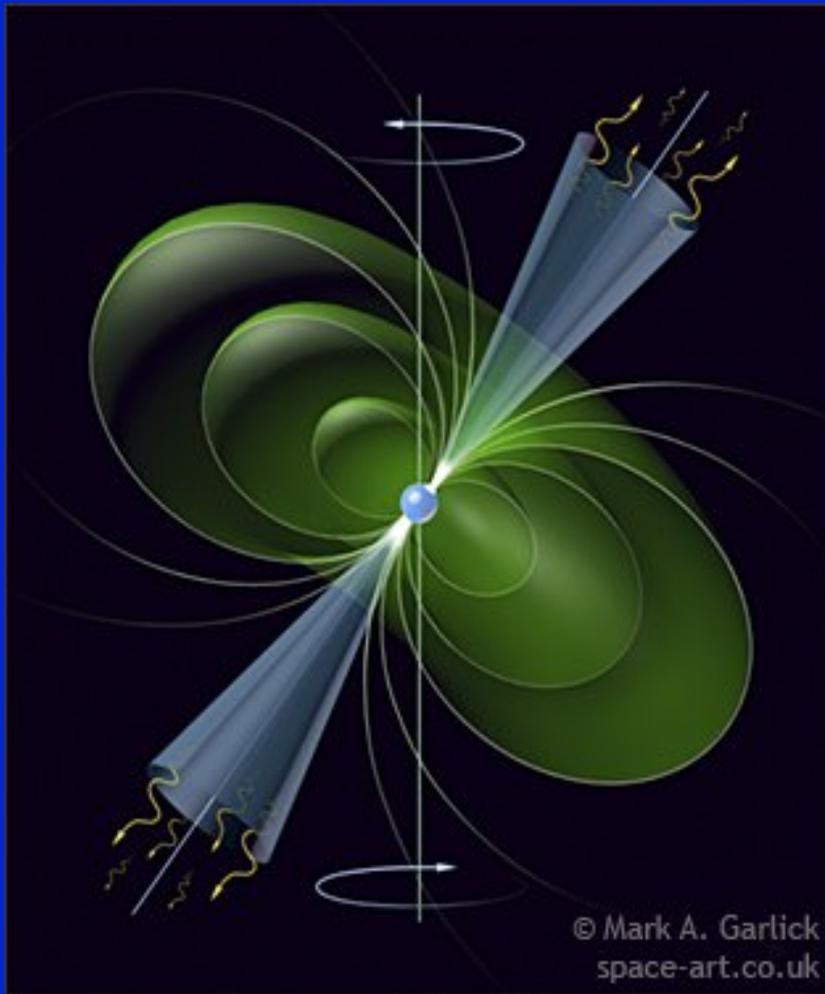
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Thermal  $\rightarrow kT = (0.5-0.7) \text{ keV}$   
Hard-X spectral tails (up to **150 keV**)

Transients: Quiescent X-ray emission  $\lesssim dE_{\text{rot}}/dt$   
Thermal  $\rightarrow kT = (0.4-0.5) \text{ keV} + A \ll A_{\text{NS}}$   
In outburst: X-ray emission  $\gg dE_{\text{rot}}/dt$   
& decays over years

X-ray Dim Isolated Neutron stars (XDINs):

Prototypical Isolated Neutron Stars,  
Nearly perfect thermal spectra and  
Stable X-rays  $kT = (0.04-0.1) \text{ keV}$

# PULSARS: BACK TO BASICS



Magnetic dipole spindown:

$$\dot{\Omega} = K B_d^2 \Omega^3$$

$$K_{\text{vacuum}} = 2R^6 \sin^2 \theta_B / 3Ic^3$$

$$K_{\text{plasma}} = R^6 (1 + \sin^2 \theta_B) / Ic^3$$

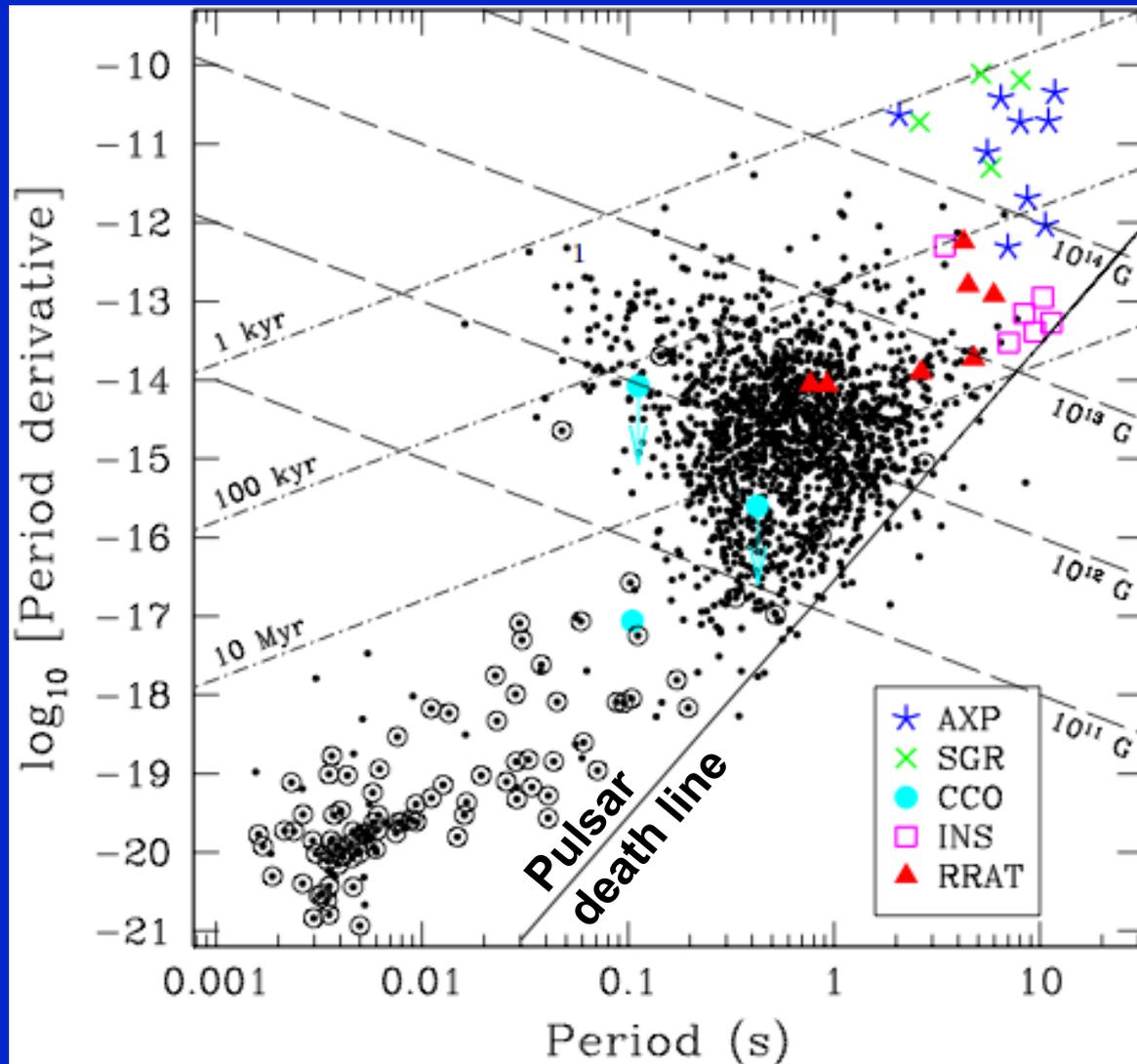
Characteristic (spindown) Age

$$\tau_c = -\Omega / 2\dot{\Omega} = P / 2\dot{P}$$

$$B_d = 3.2 \times 10^{19} (P[s] \dot{P})^{1/2} \text{ G}$$

$$= 3.2 \times 10^{19} P[s] (2\tau_c[s])^{-1/2} \text{ G}$$

# An observational perspective:



$P - \dot{P}$  diagram for 1704 objects, including 1674 RPPs (small black dots).

Magnetars have some of the lowest spins, despite being the most luminous. A lower- $B$  NS would be over the death line & thus no longer a RPP.

RPP = rotation powered pulsar  
CCO = central compact object  
INS = XDIN  
RRAT = rotating radio transient

(Kaspi 2010)

# An observational perspective:

**SGR0418+5729**: discovered by Fermi/GBM on 9 June 2009 through 2 weak bursts (**Van der Horst et al. 2010**)

+

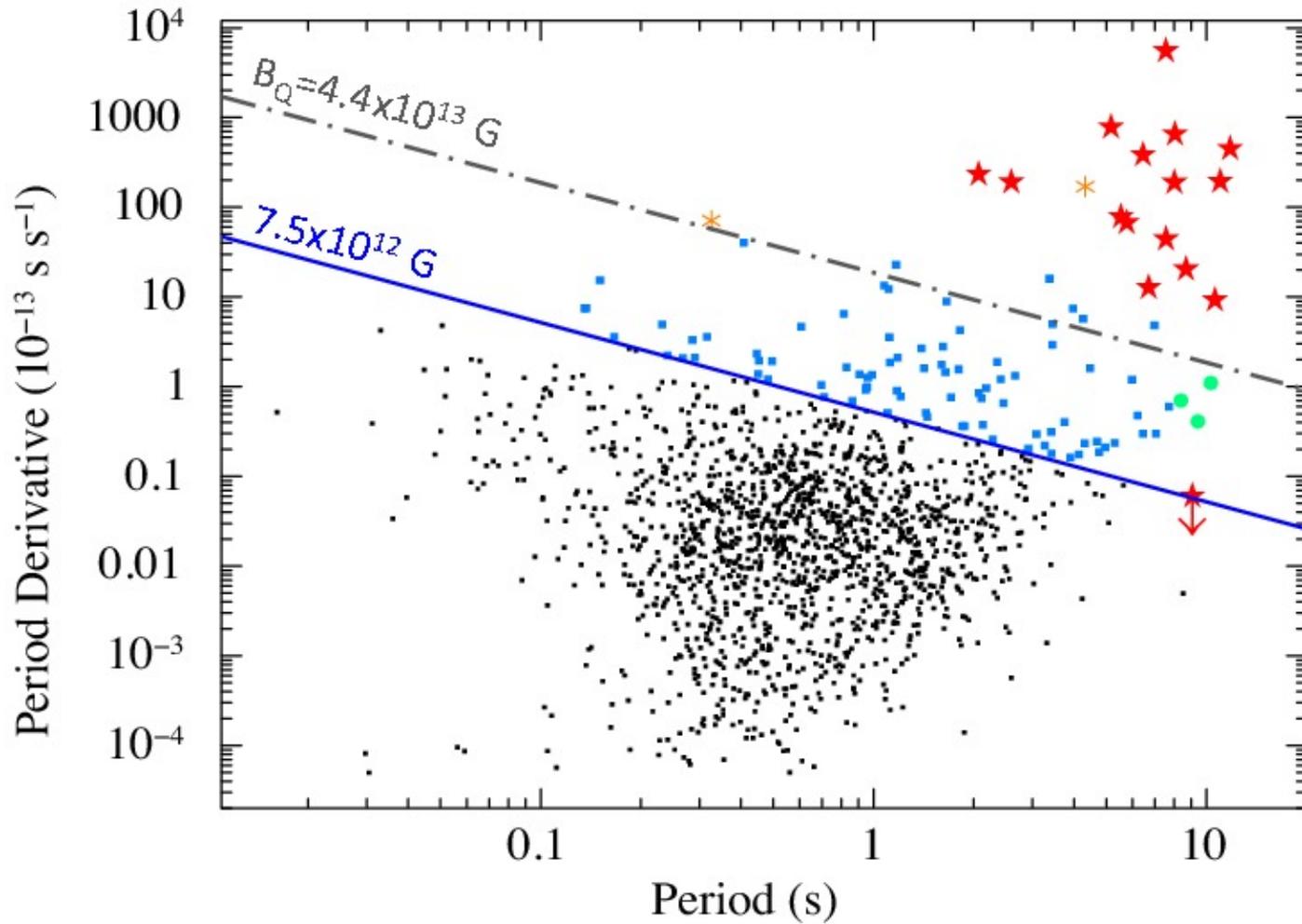
It displayed long-lasting enhanced X-ray emission & a spin period of  $P \approx 9.1$  s (**Esposito et al. 2010**)

By 23 Sep. 2010 its  $F_X$  decreased by  $\times 10^3$  to  $6 \times 10^{31}$  erg/s. Hot thermal emission ( $kT \approx 0.67$  keV) & very small emission region ( $R \approx 0.1$  km) (**Rea et al. 2010**)

No period derivative measured:  $\dot{P} < 6 \times 10^{-15}$  s/s

$\rightarrow B_{\text{dipole}} < 7.5 \times 10^{12}$  G &  $\tau_c > 2.4 \times 10^7$  yrs  
(**Rea et al. 2010**)

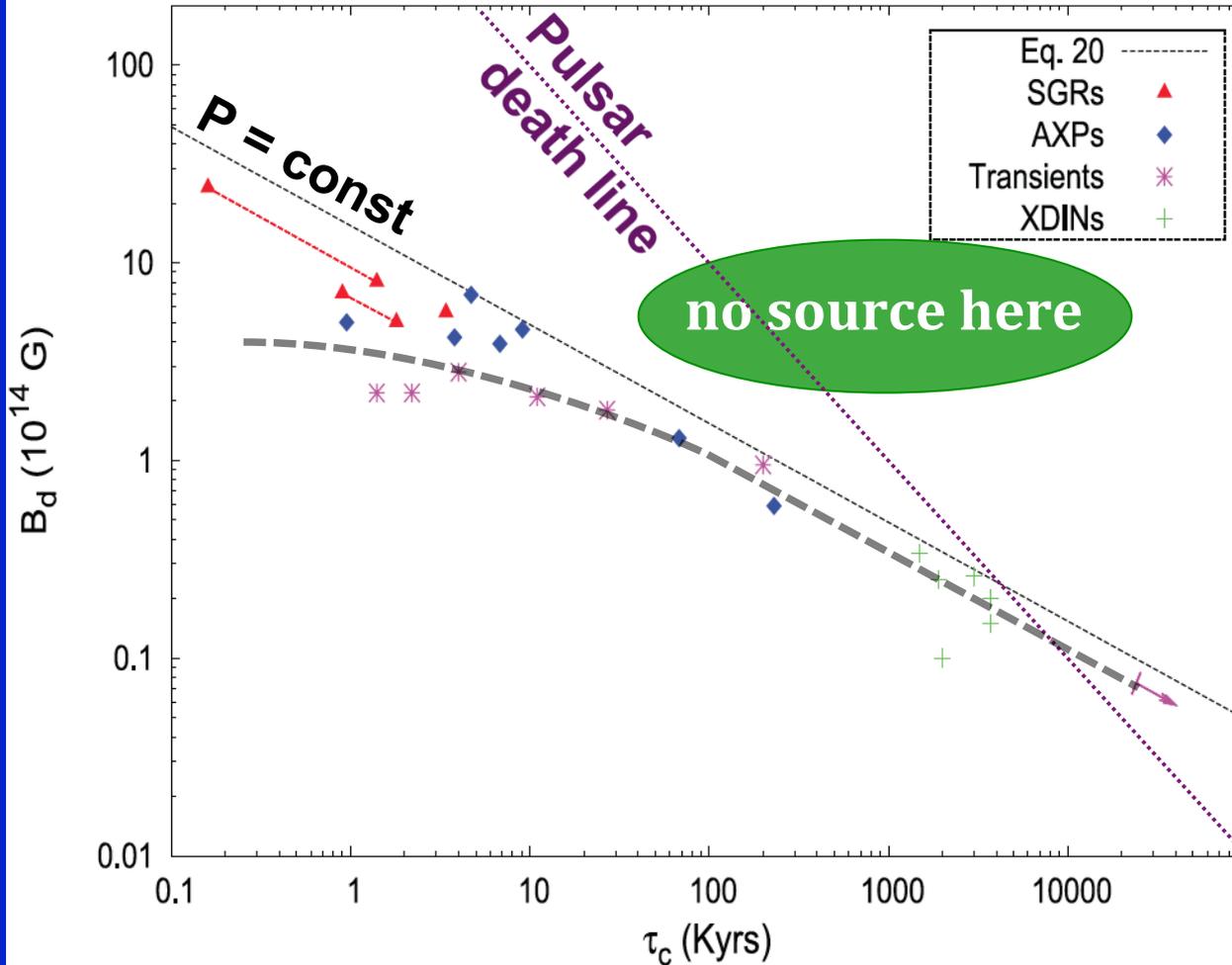
# An observational perspective:



(Rea et al. 2010)

# A physical perspective:

Observed Distribution of  $B_d$  vs.  $\tau_c$



No high- $B_d$  objects  
With an old  $\tau_c$

( $B_d$  = dipole B-field  
 $\tau_c$  = spin-down age)

There is a limiting  
spin period:  $P \lesssim 10$  s

It is expected if  $B_d$   
decays fast enough:

$$\dot{B}_d \propto B_d / \tau_d$$

$$\tau_d \propto B^{-\alpha} \text{ with } \alpha < 2$$

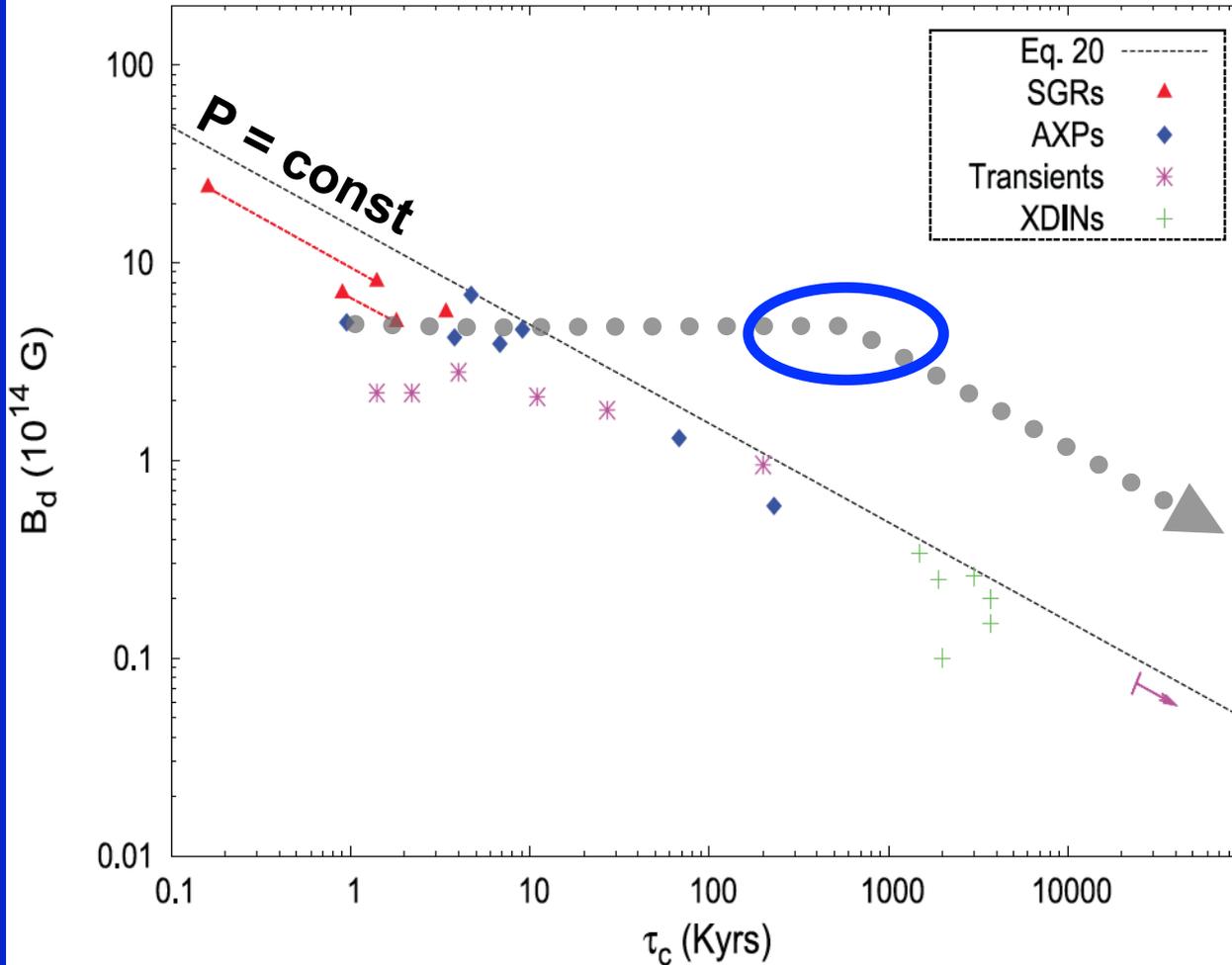
(Colpi et al. 2000)

(Dall'Osso, JG & Piran 2012)

(7 SGRs, 11 AXPs, 6/7 XDINs)

# Could this be a selection effect? **Unlikely!**

Observed Distribution of  $B_d$  vs.  $\tau_c$



**SGRs** are detected through their **bursts**:

Before  $B_d$  decays a lot,  $\tau_c \sim t \rightarrow$  more sources at larger  $\tau_c$

If bursts follow  $B_d \rightarrow$  SGRs should be detected mostly with  $P \gg 10$  s

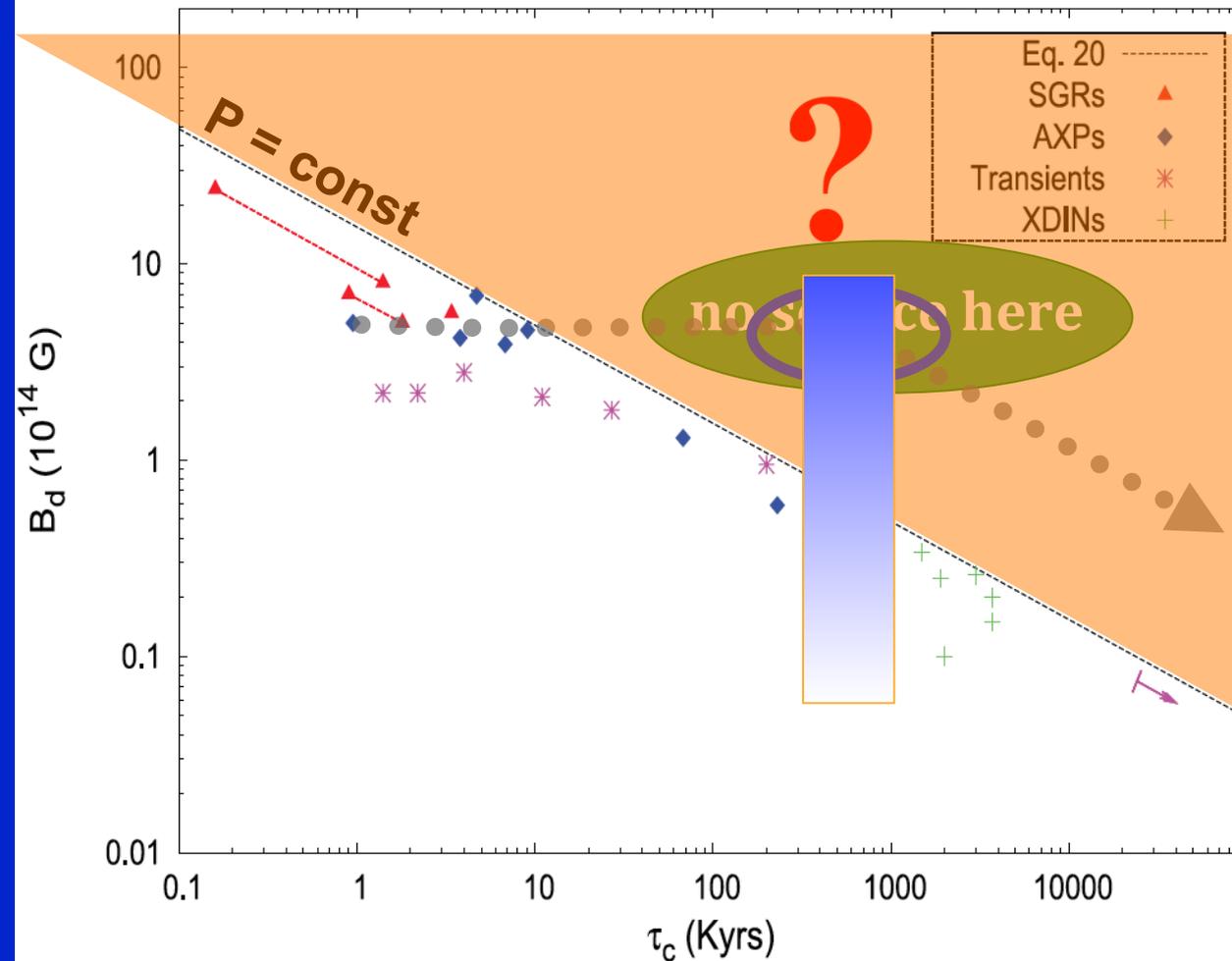
If bursts follow  $B_d$  how do they “know” to stop at  $P > 10$  s ?

(Dall’Osso, JG & Piran 2012)

■ For  $\tau_c = \text{const}$ , larger  $B_d$  is easier to detect

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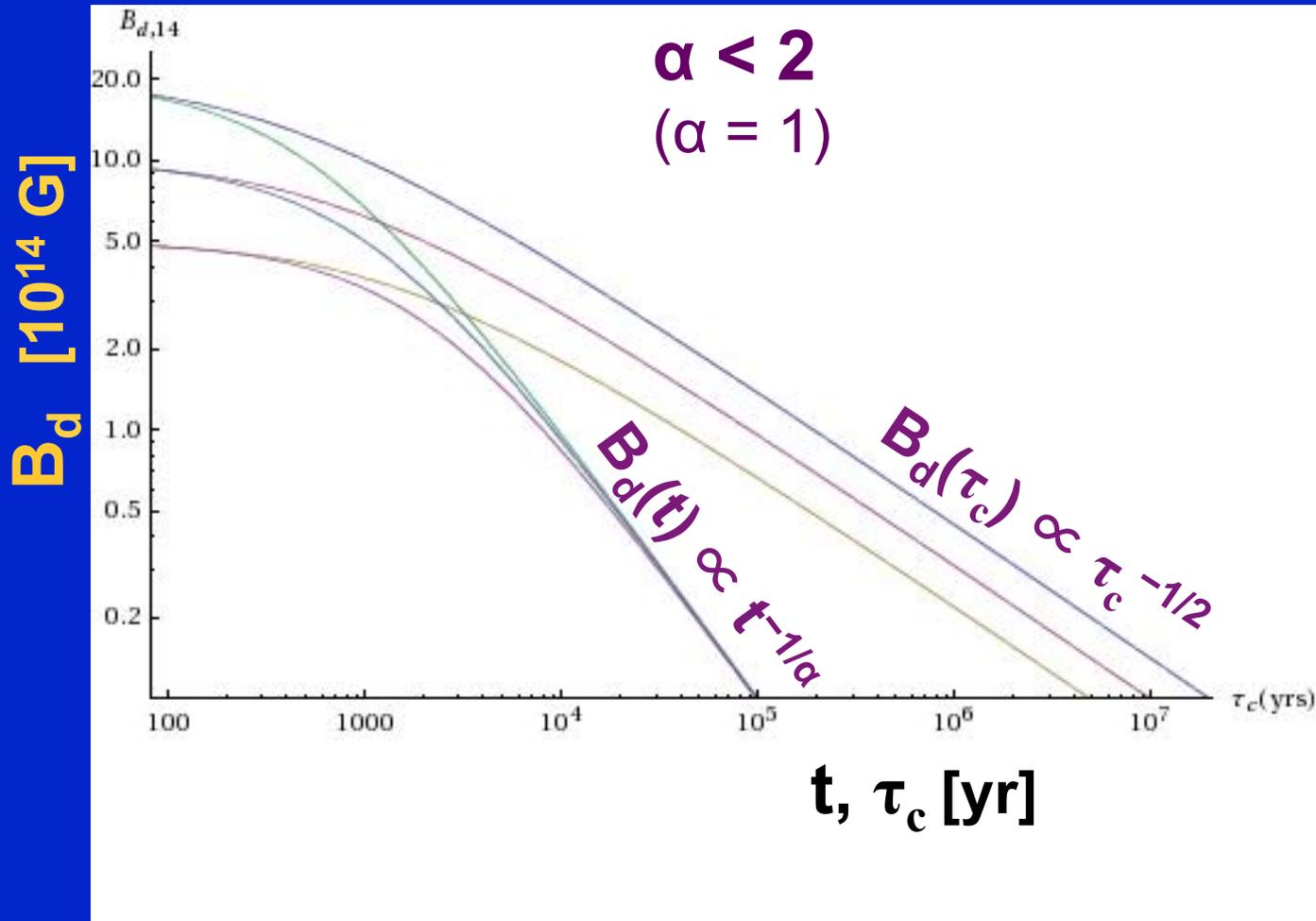
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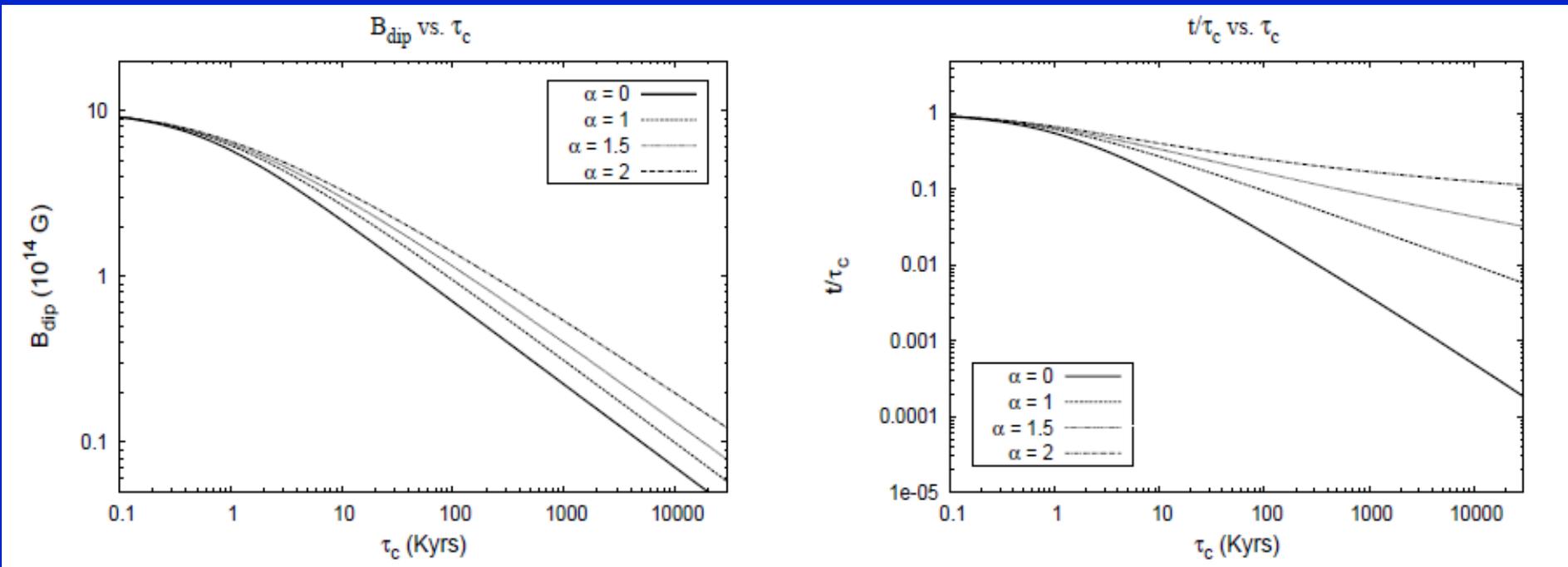
# Dipole field decay: parameterized model

$$\dot{B}_d \propto B_d / \tau_d \propto B_d^{1+\alpha}, \quad \tau_d \propto B_d^{-\alpha} \text{ with } \alpha < 2$$



# Dipole field decay: parameterized model

Once  $B_{\text{dip}}$  decays significantly, true age  $t \ll$  spin-down age  $\tau_c$



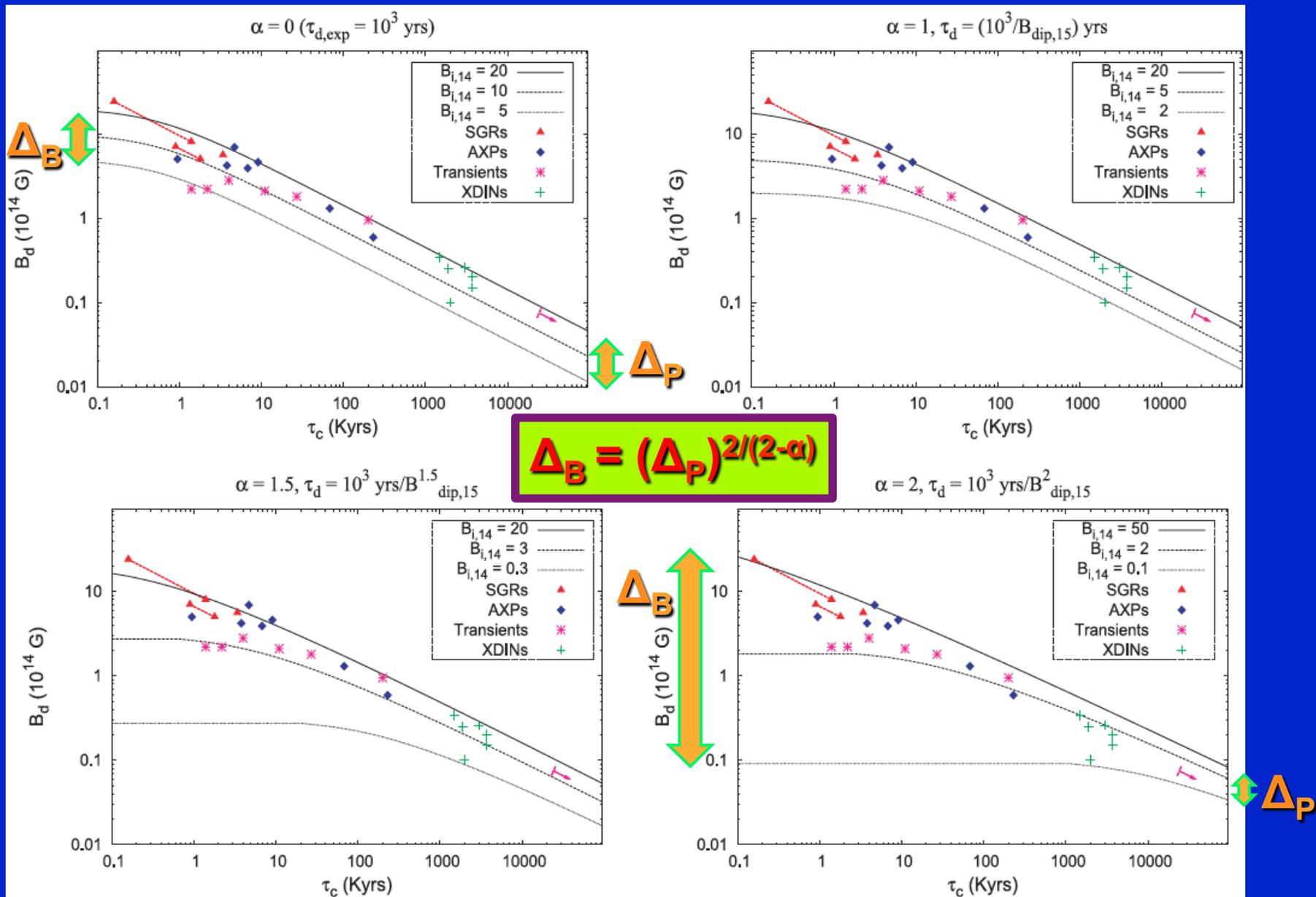
$$B_d \propto P \tau_c^{-1/2}$$

For  $t \gg \tau_{d,i}$ ,  $\alpha < 2$ :

$$B_d \propto t^{-1/\alpha} \propto \tau_c^{-1/2}$$

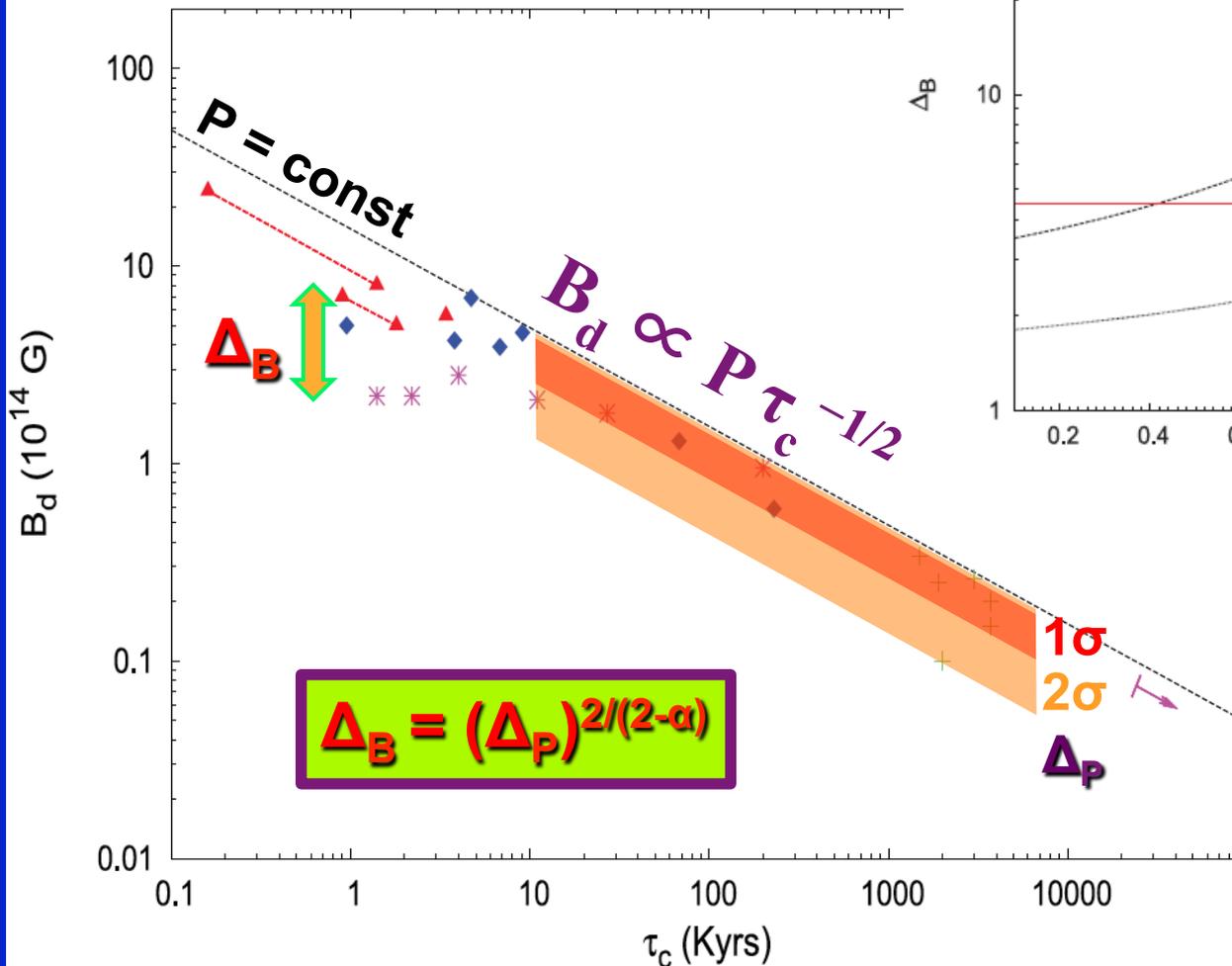
$$t/\tau_c \propto \tau_c^{(\alpha-2)/2}$$

# Dipole decay modes vs. Observations

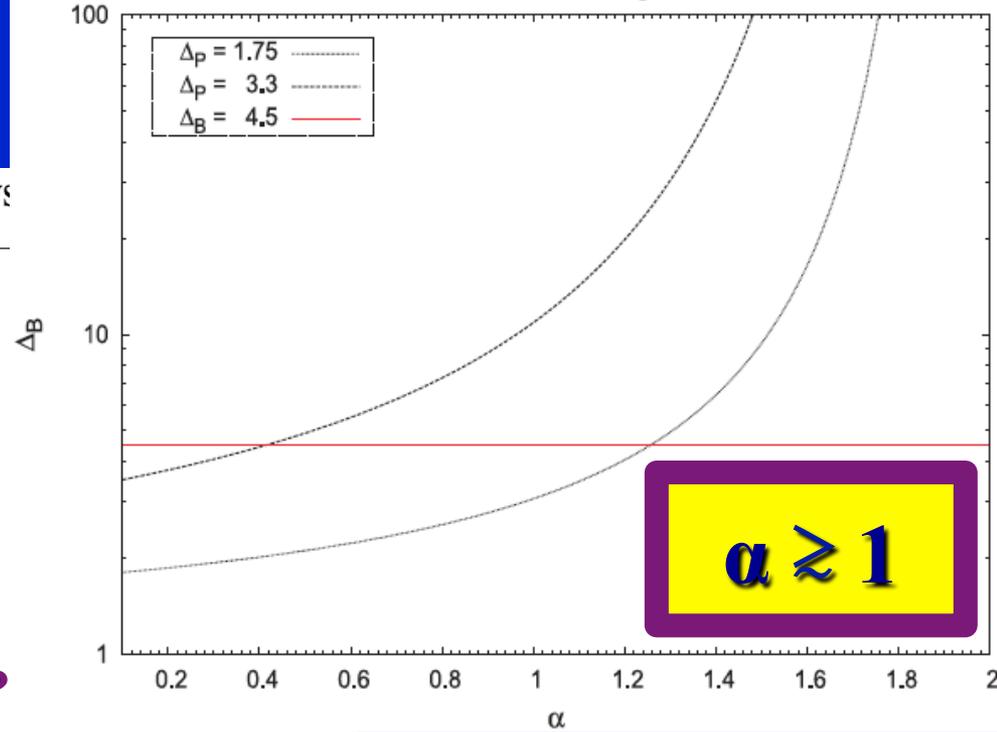


# Dipole decay modes vs. Observations

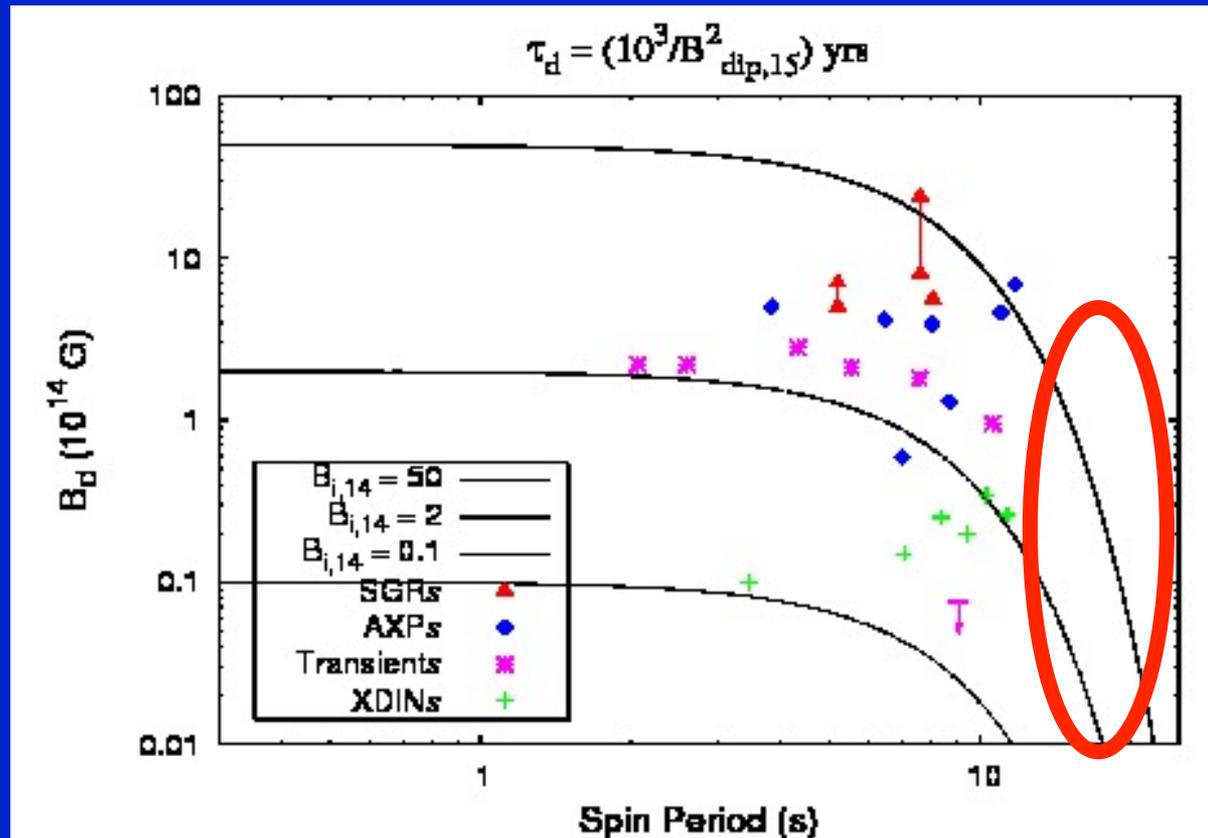
Observed Distribution of  $B_d$  vs  $\tau_c$



Allowed spread in  $B_{dip,i}$  vs.  $\alpha$

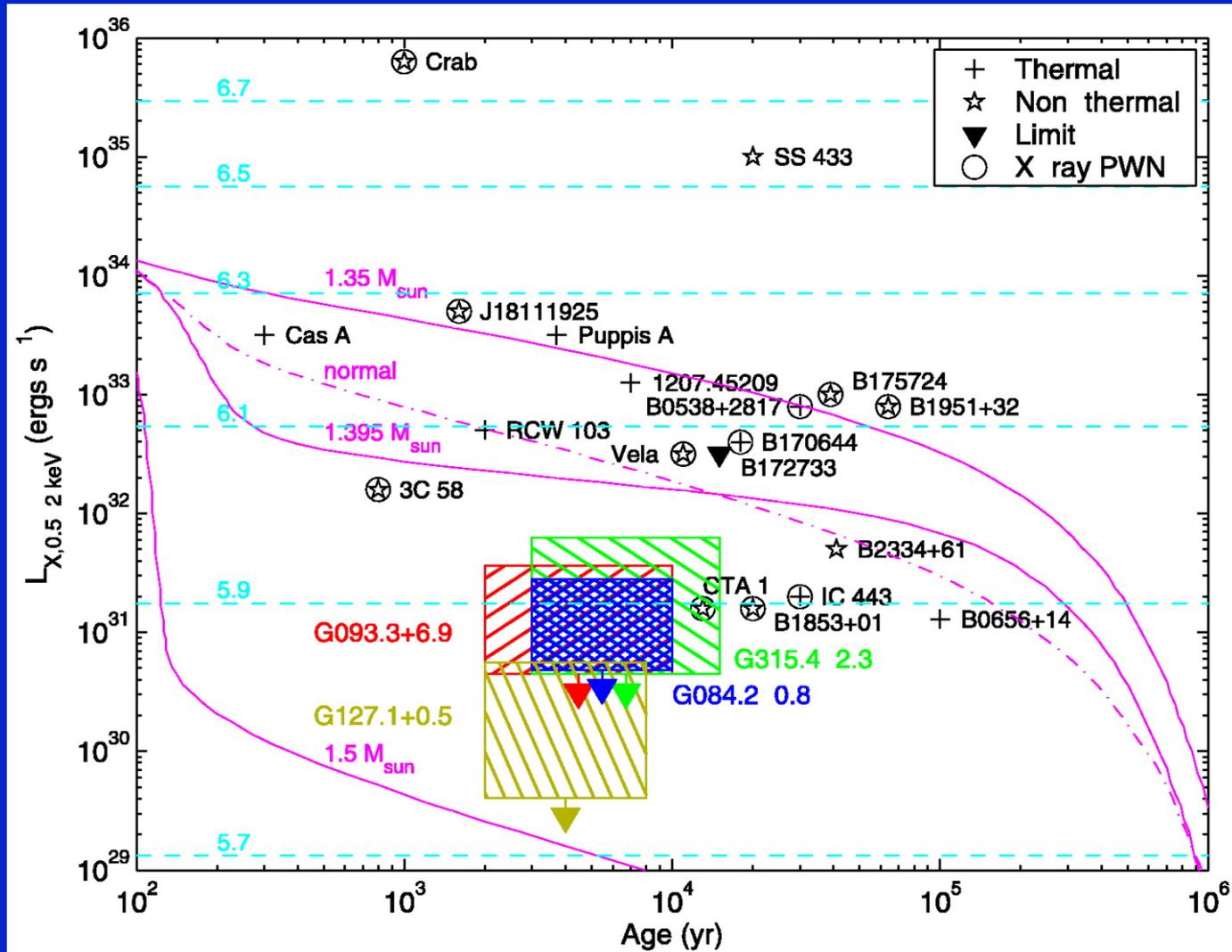


# Dipole decay modes vs. Observations



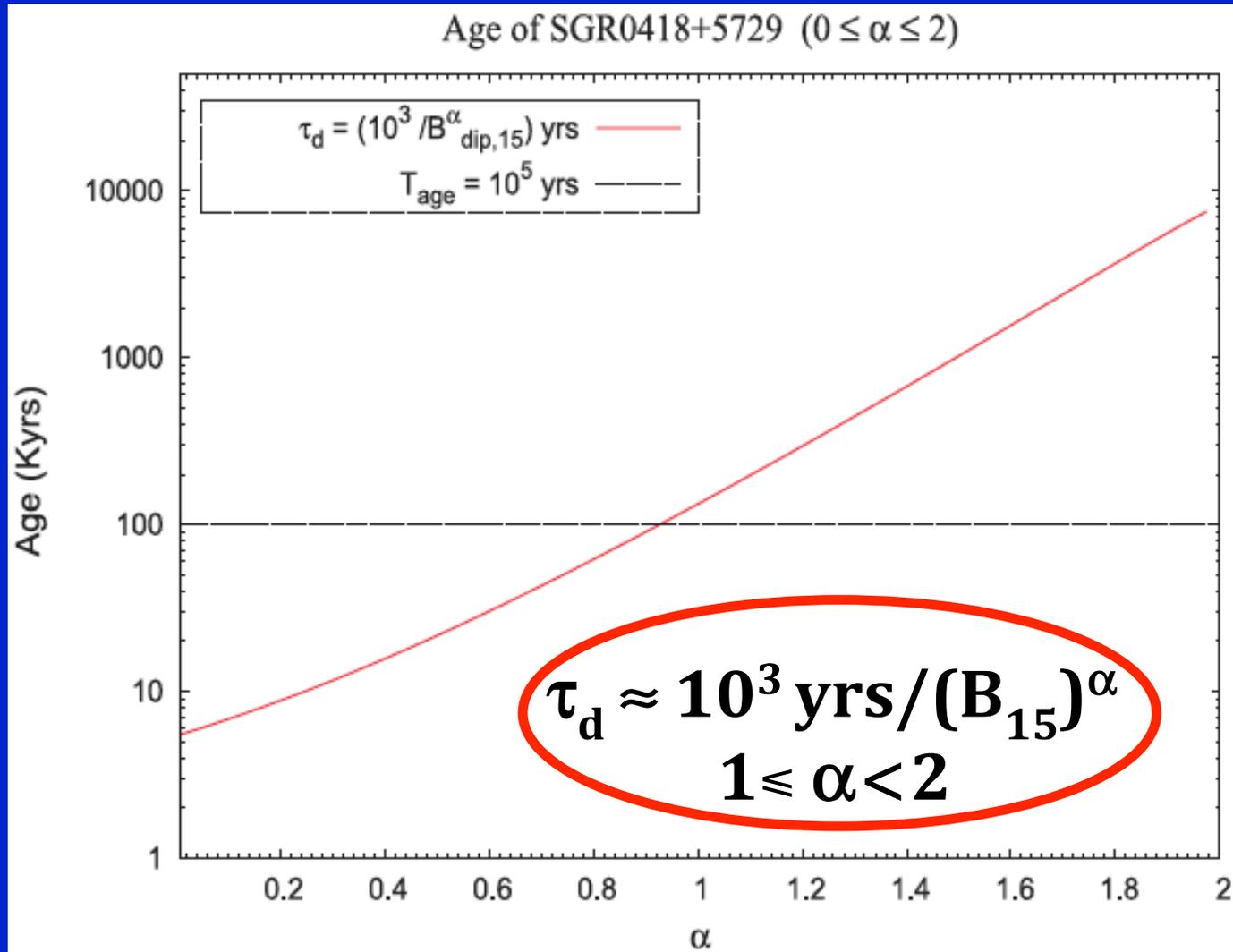
$$a < 2$$

# Age constraints on the field decay



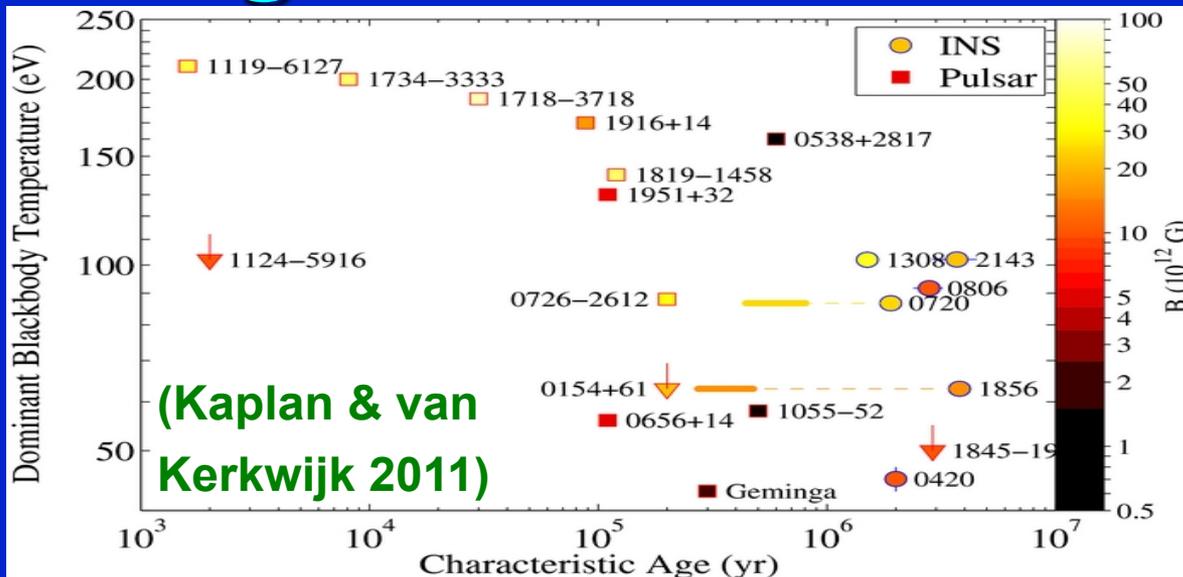
(Kaplan 2004)

# Age constraints on the field decay



$L_X(0.5-10 \text{ keV}) < 6 \times 10^{31} \text{ erg/s} \rightarrow t > 10^5 \text{ yr}$

# Age constraints on $\alpha$ from XDINs



$L_X \approx \text{few} \times 10^{31} \text{ erg/s}$

$kT < 0.1 \text{ keV} \approx 10^6 \text{ K}$

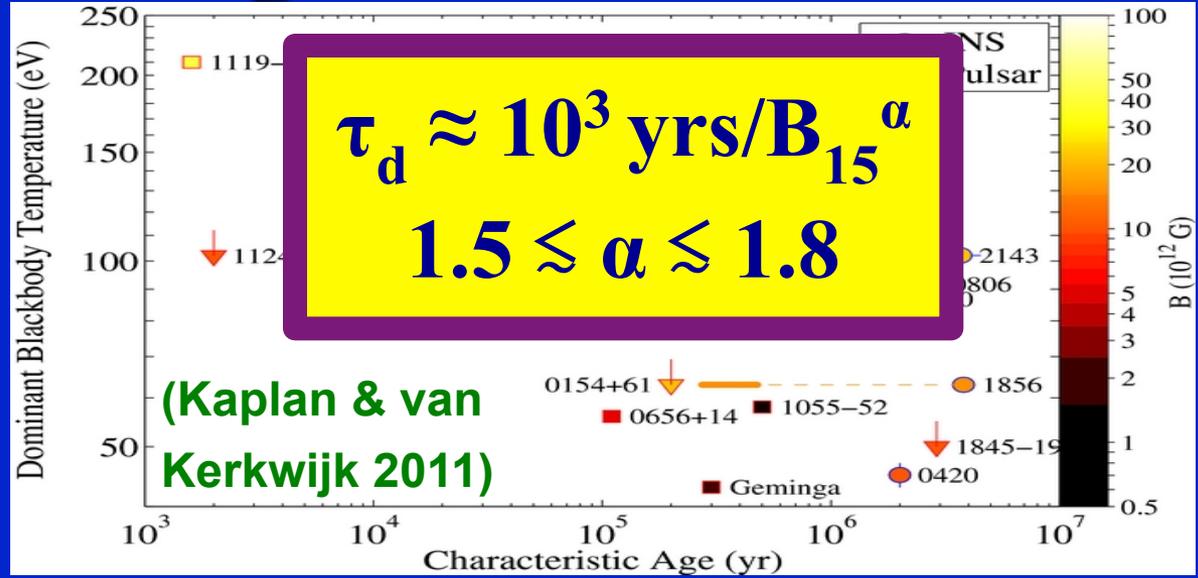
→ ages  $t \lesssim 10^5 \text{ yrs}$

( $t/\tau_c \lesssim 0.1$ )

→  $\alpha \gtrsim 1.5$

Name	$\tau_c$ [kyr]	$t_{\text{age}}$ [kyr] ( $\alpha = 1$ )	$t_{\text{age}}$ [kyr] ( $\alpha = 1.5$ )	$kT_{\text{eff}}$ [eV]
RXJ 1308	1500	30	100	102
RXJ 0806	3000	40	160	96
RXJ 0720	1900	42	160	90
RXJ 2143	3700	50	200	100
RXJ 1856	3700	66	350	63
RXJ 0420	2000	100	580	44

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$L_x \approx \text{few} \times 10^{31} \text{ erg/s}$   
 $kT < 0.1 \text{ keV} \approx 10^6 \text{ K}$   
 $\rightarrow$  ages  $t \lesssim 10^5 \text{ yrs}$   
 $(t/\tau_c \lesssim 0.1)$   
 $\rightarrow \alpha \gtrsim 1.5$

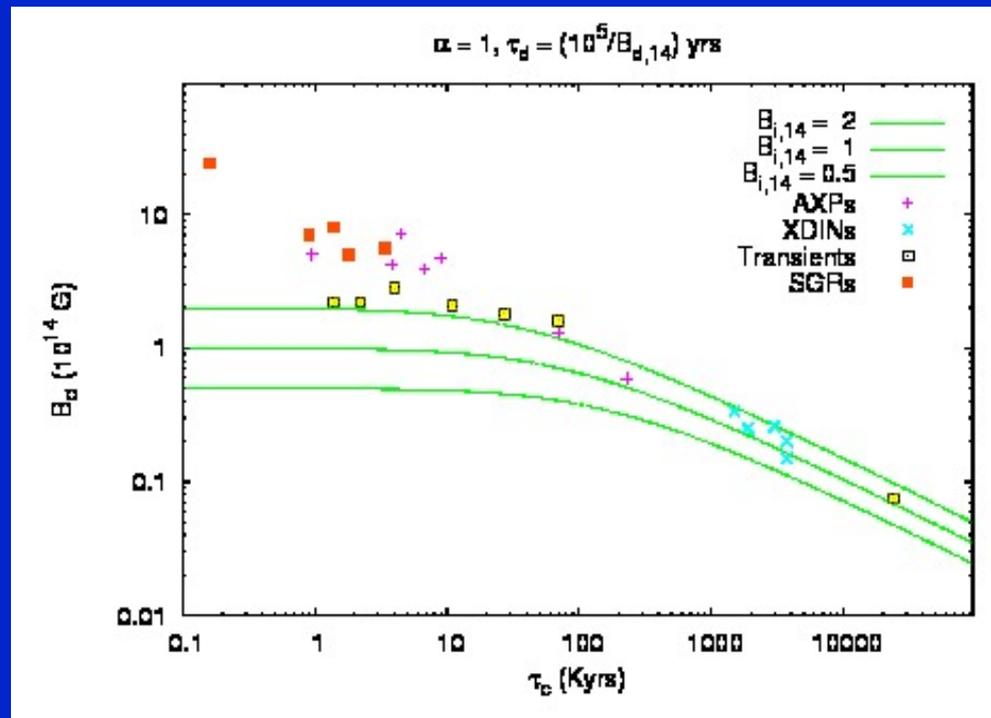
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# Dipole field decay: physical models

Hall Decay of the magnetic field in the NS crust is expected to operate fast enough at large B

$$\tau_d \sim 10^4 \text{ yrs} / B_{15} \quad (\alpha = 1)$$

(Goldreich & Reisenegger 1992; Cumming et al. 2004)

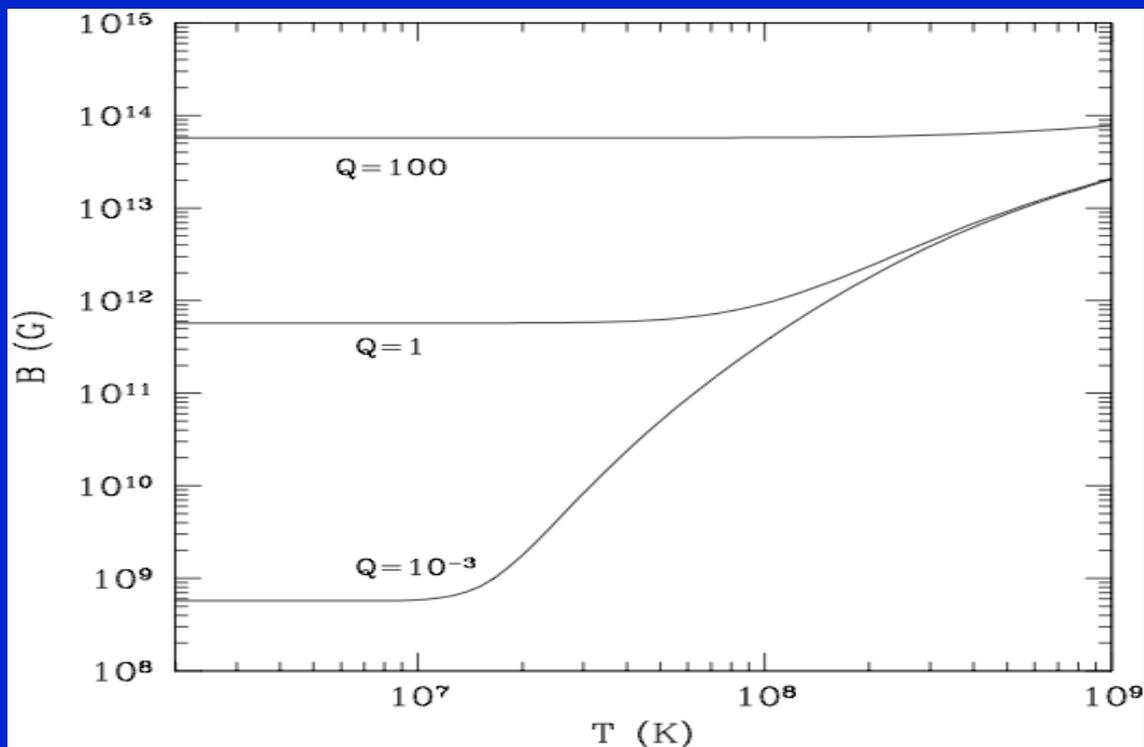


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B-field above which the Hall effect dominates over direct Ohmic decay

Q = crust impurity parameter

(Cumming et al. 2004)

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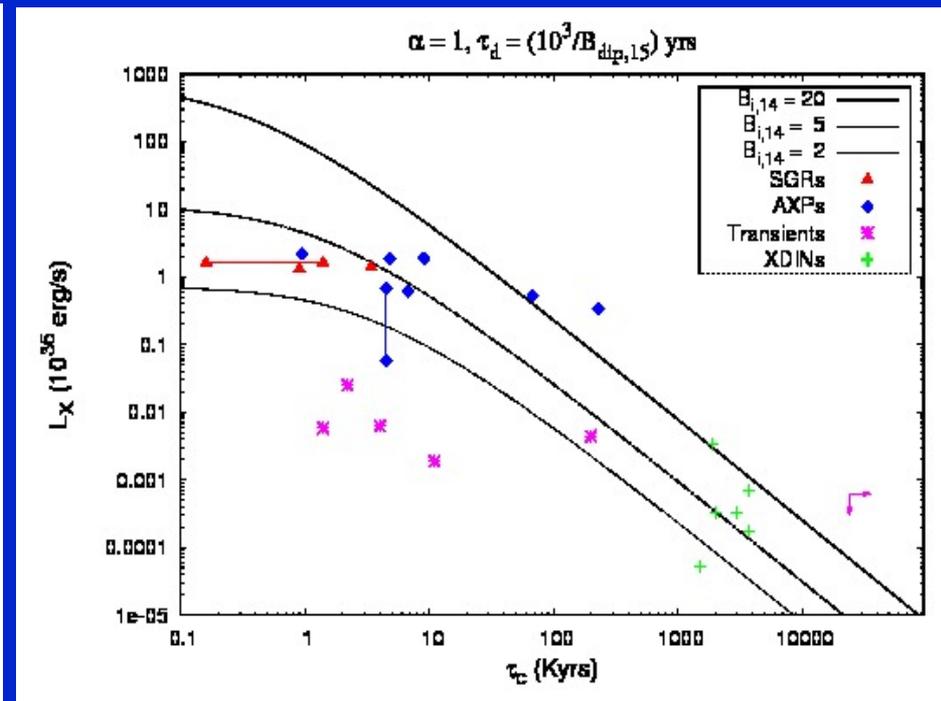
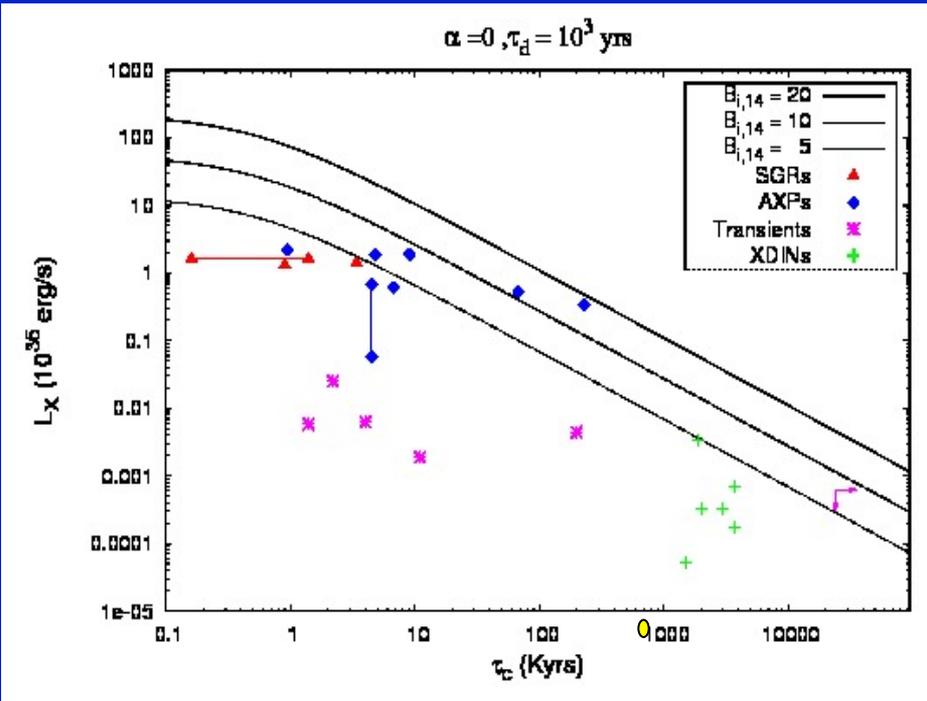
(Goldreich & Reisenegger 1992; Cumming et al. 2004)

B-independent Ohmic decay + diffusion in deep crust gives a power-law decay,  $B \propto t^{-1/\alpha}$  with  $1.5 \lesssim \alpha \lesssim 1.8$  and  $\tau_{d,i} \approx 10^4 \rho_{i,12} \text{ yrs} \propto B_i^{-\alpha} \rightarrow \alpha_{\text{eff}} = 0, \Delta_B = \Delta_P \rightarrow$   
problematic!

(Urpin, Changmugam & Sang 1994 [...] → Urpin & Yakovlev 2008)

# Luminosity Evolution vs. Dipole Decay

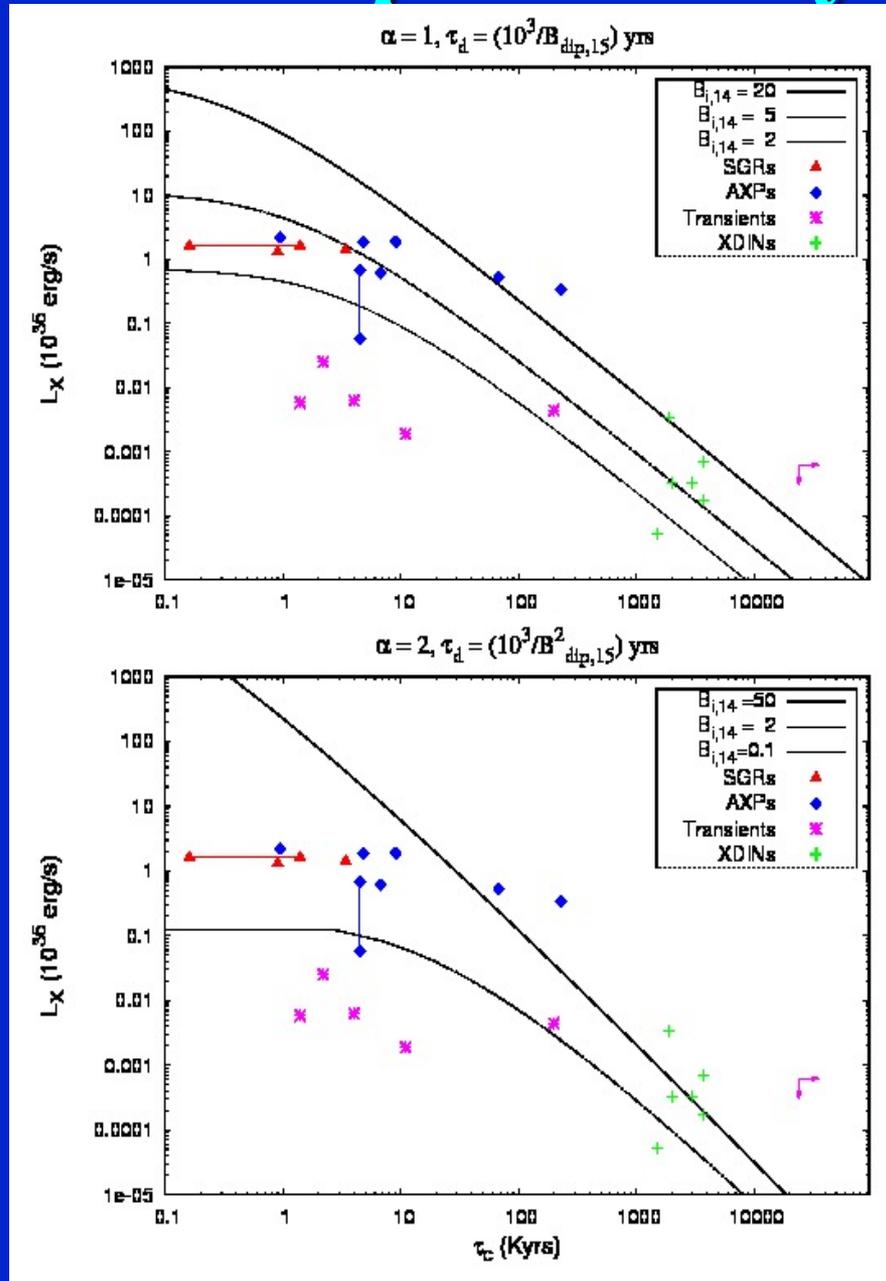
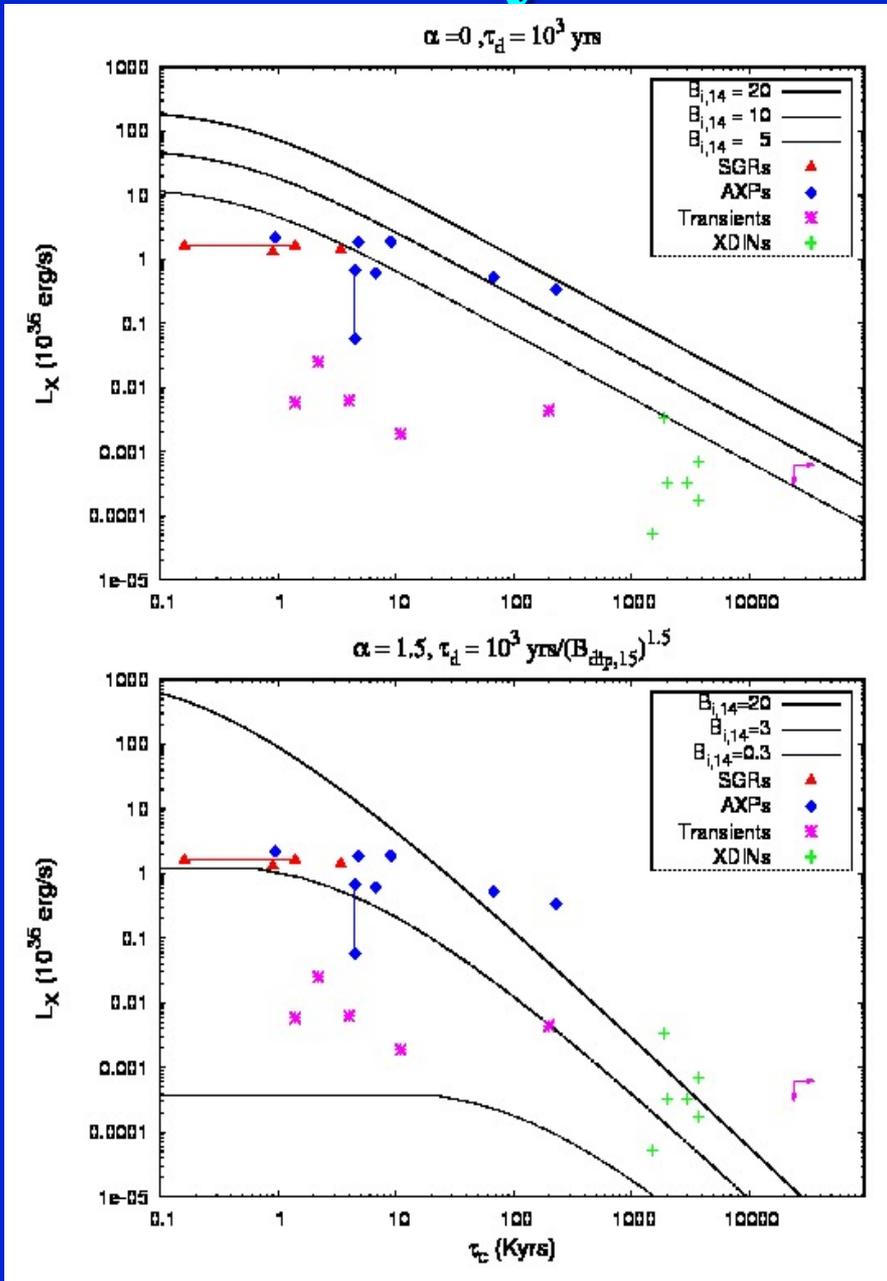
New observable:  $L_X(2-10 \text{ keV})$



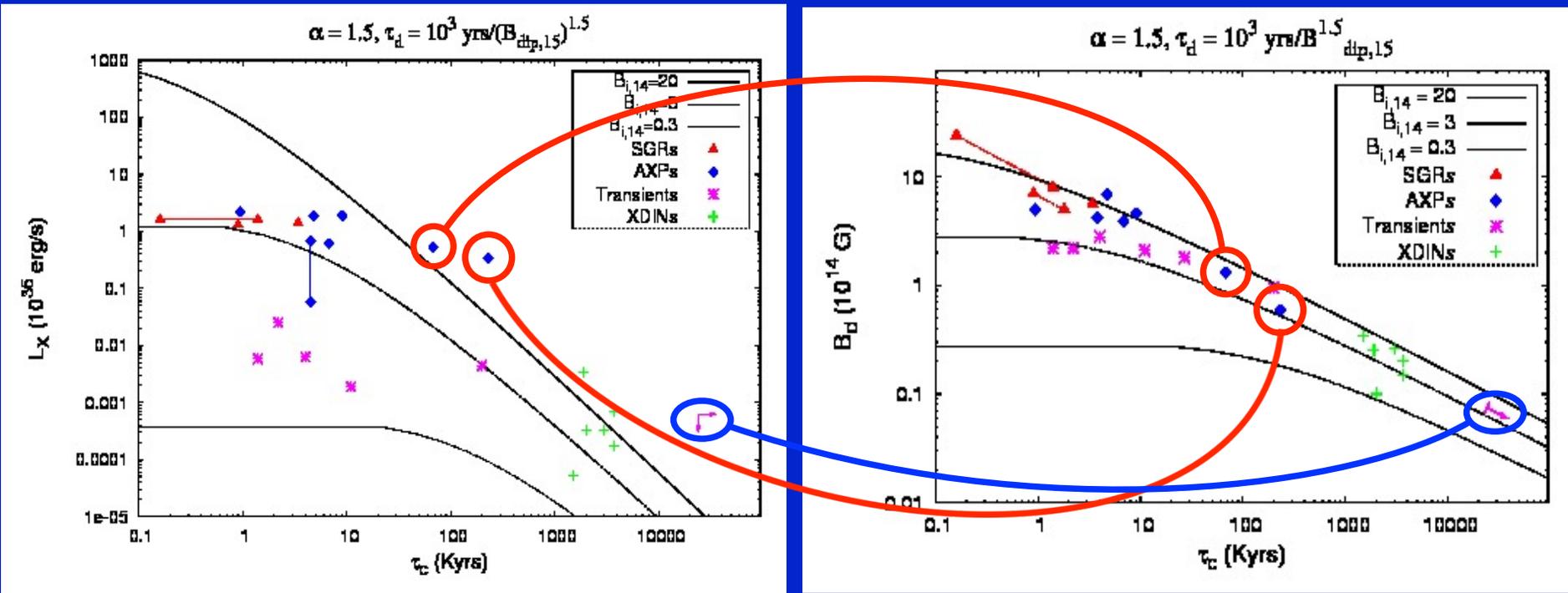
$$E_{B,d} \sim R^3 B_d^2$$

$$|\dot{E}_{B,d}| = L_B \sim R^3 B_d^{2+\alpha} / A$$

# Luminosity Evolution vs. Dipole Decay



# Luminosity Evolution vs. Dipole Decay

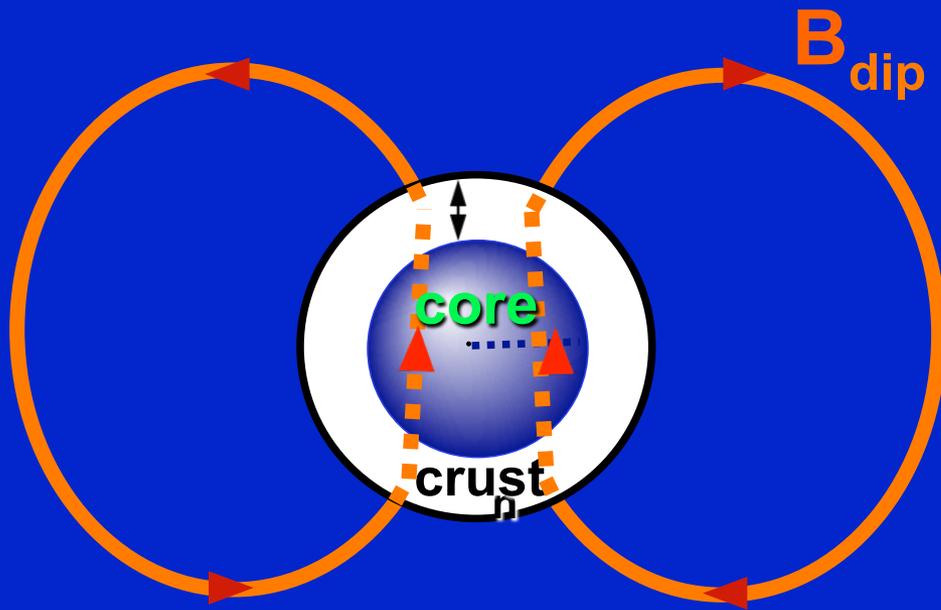


- We use  $L_X(2-10 \text{ keV}) \sim (0.3-0.5) L_{X,bol} \rightarrow$  bolometric + GR corrections imply  $L_{X,intrinsic} \sim \text{several} \times L_X(2-10 \text{ keV})$
- $L_B$  also goes to other channels (bursts,  $\nu$ 's)  $\rightarrow$  conservative

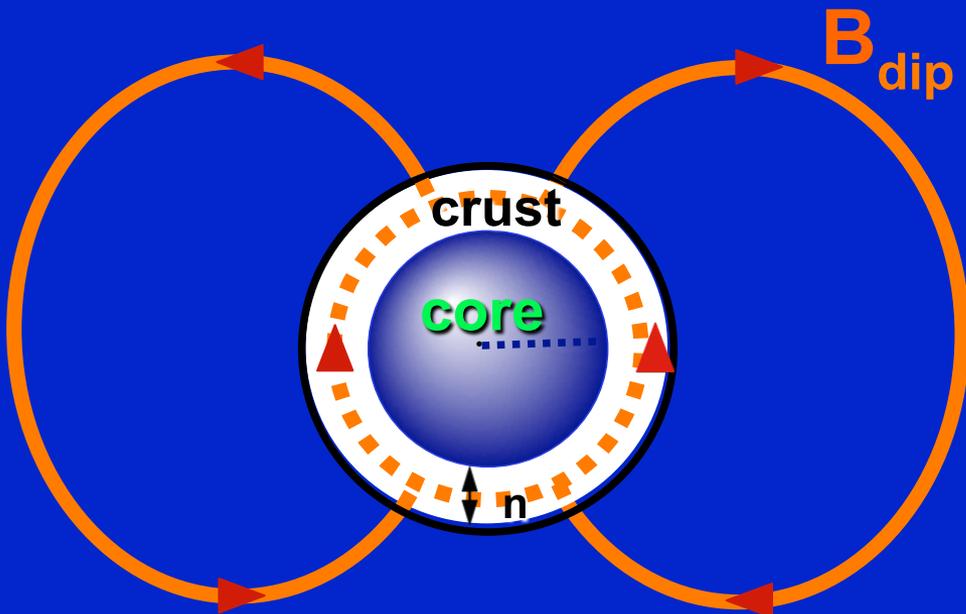
**Decay of the Dipole Field does not match  $L_X$  evolution**

**In particular it cannot power sources @  $\tau_c \gtrsim 10^5 \text{ yrs}$**

# Interior Field



1. Decay of the interior field releases heat in the core, which is conducted to the surface, producing a quasi-thermal emission



2. Interior field is confined to the crust, and/or core field does not decay (condensed phase?)

# Decay of Internal Field: toy models

## Ambipolar diffusion in the NS core

$$L_{B,\text{int}} = (R_*^3/3)B_{\text{int}}^2/\tau_{d,\text{int}} \approx 10^{38} B_{\text{int},16}^{16/5} R_{*,6}^3 \text{erg s}^{-1}$$

$$\frac{dU^-}{dt} \simeq 9.6 \times 10^{20} T_9^8 \rho_{15}^{2/3} \text{erg cm}^{-3} \text{s}^{-1}$$

$$T_{8,\text{eq}}^{(s)} \simeq 2.7 B_{15}^{2/5} \left(\frac{\rho_{15}}{0.7}\right)^{-2/3}$$

$$T_s \simeq 1.17 \times 10^6 \text{ K} \left[ (7\zeta)^{9/4} + \left(\frac{\zeta}{3}\right)^{5/4} \right]^{1/4}$$

where  $\zeta \equiv T_{c,9} - 0.001 g_{14}^{1/4} (7 T_{c,9})^{1/2}$  and  $g_{14} \approx 1.87$  is the surface gravity.

$$L_{X,\infty} = \left(1 - \frac{2GM_*}{R_*c^2}\right) L_X \approx 10^{33} \text{ erg s}^{-1} \left[ (7\zeta)^{9/4} + \left(\frac{\zeta}{3}\right)^{5/4} \right]$$

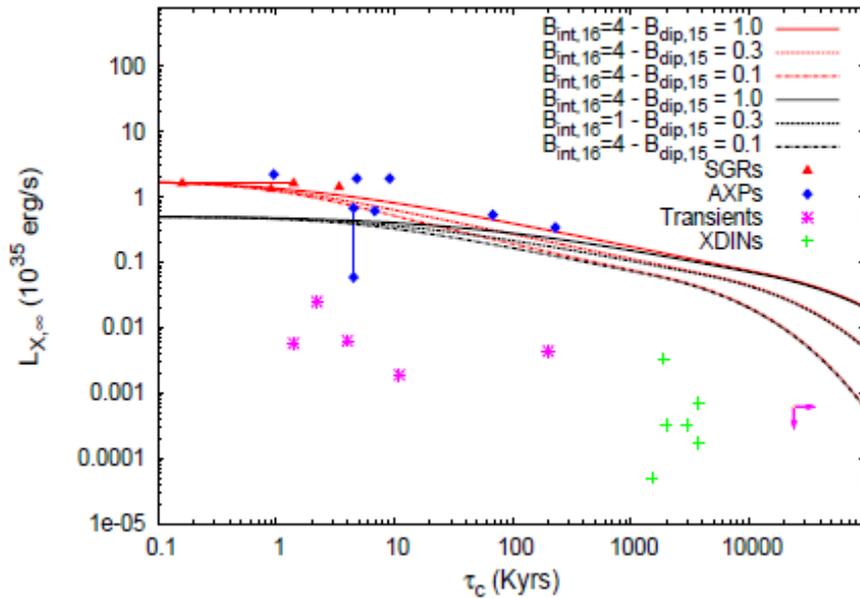
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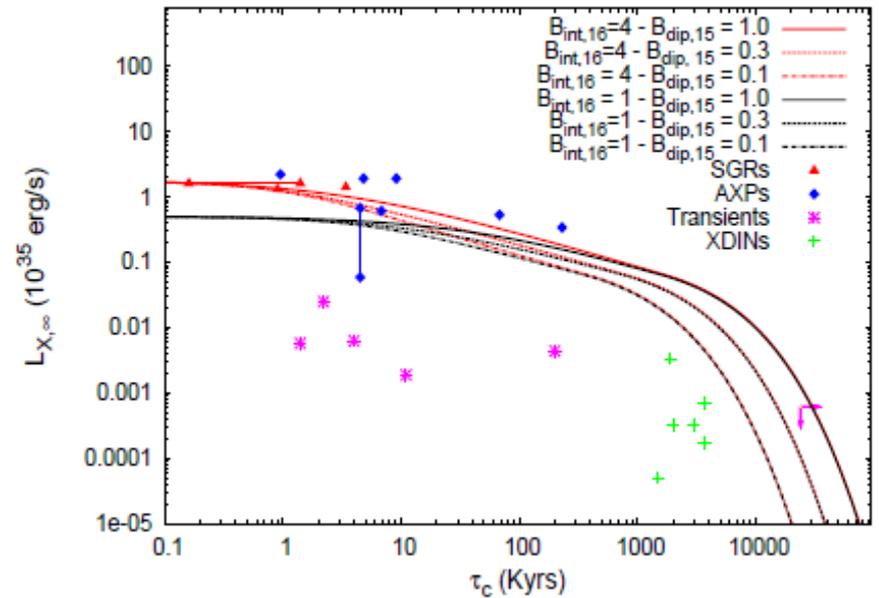
$$L_{B,int} = (R_*^3/3)B_{int}^2/\tau_{d,int} \approx 10^{38} B_{int,16}^{16/5} R_{*,6}^3 \text{ erg s}^{-1}$$

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$$\tau_{d,int} = 9.6 \times 10^3 \text{ yrs} / (B_{int,16})^{6/5} - \alpha = 1$$



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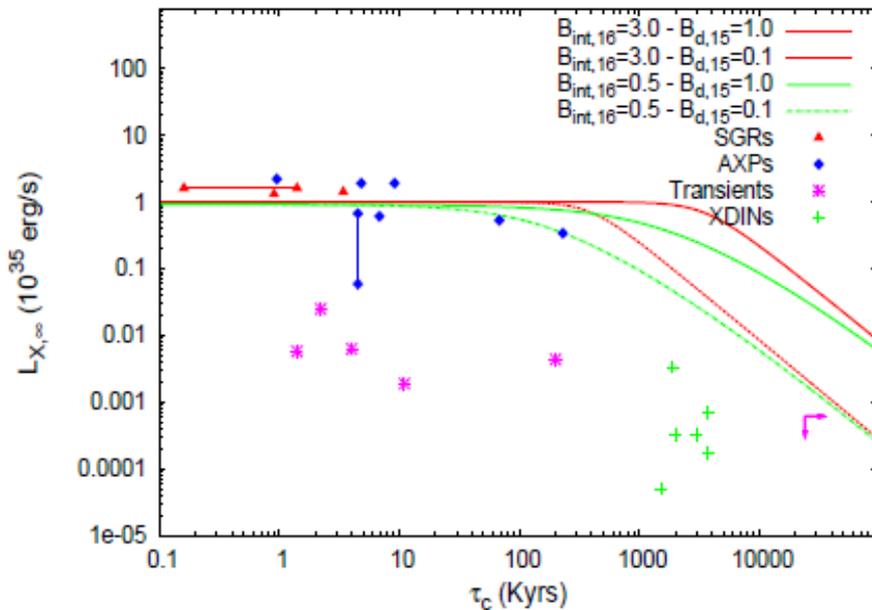
# Decay of Internal Field: toy models

## Hall decay of internal field in the deep crust

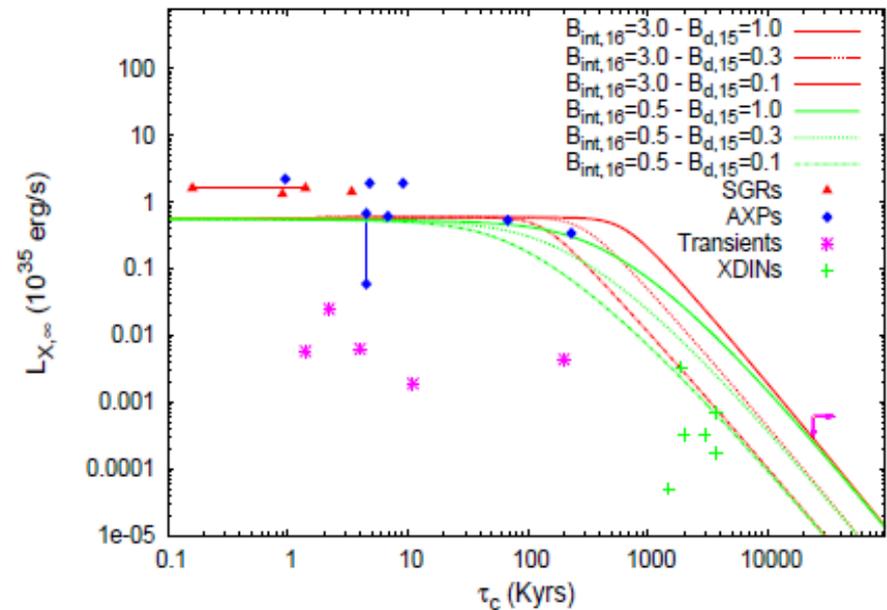
$$L_{B,int}^{(I)} \approx 2 \times 10^{36} B_{16}^3 R_{*,6} \text{ erg s}^{-1}$$

(Arras, Cumming & Thompson 2004)

$$\tau_{d,int} = 2.4 \times 10^4 / B_{int,16} \text{ yrs} - \alpha = 1$$



$$\tau_{d,int} = 2.4 \times 10^4 / B_{int,16} \text{ yrs} - \alpha = 1.5$$



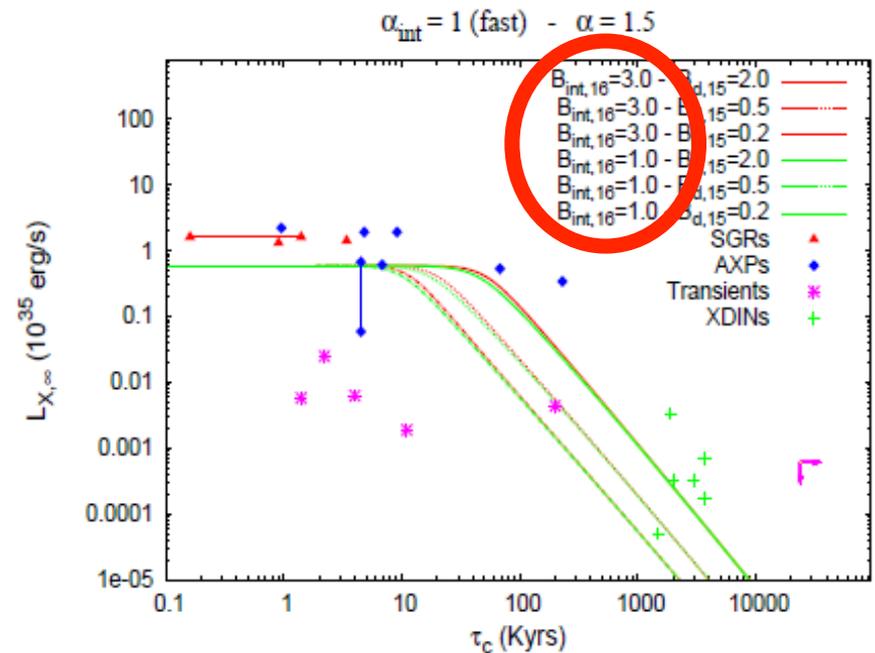
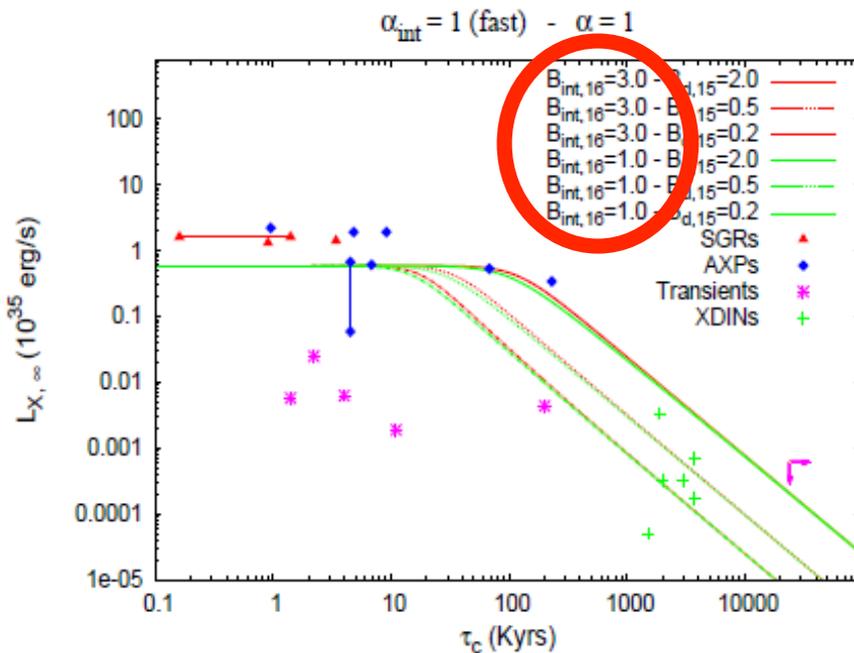
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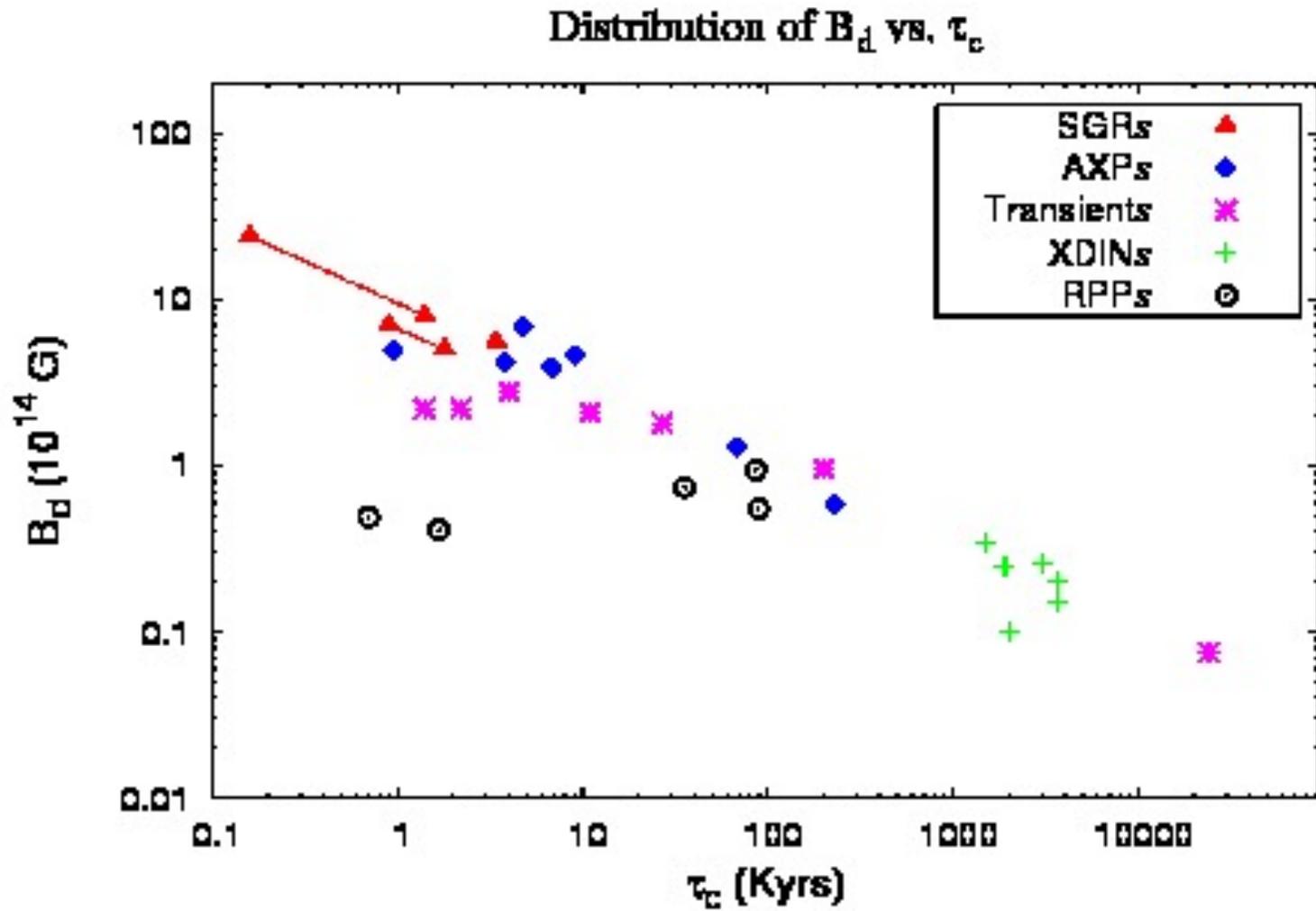
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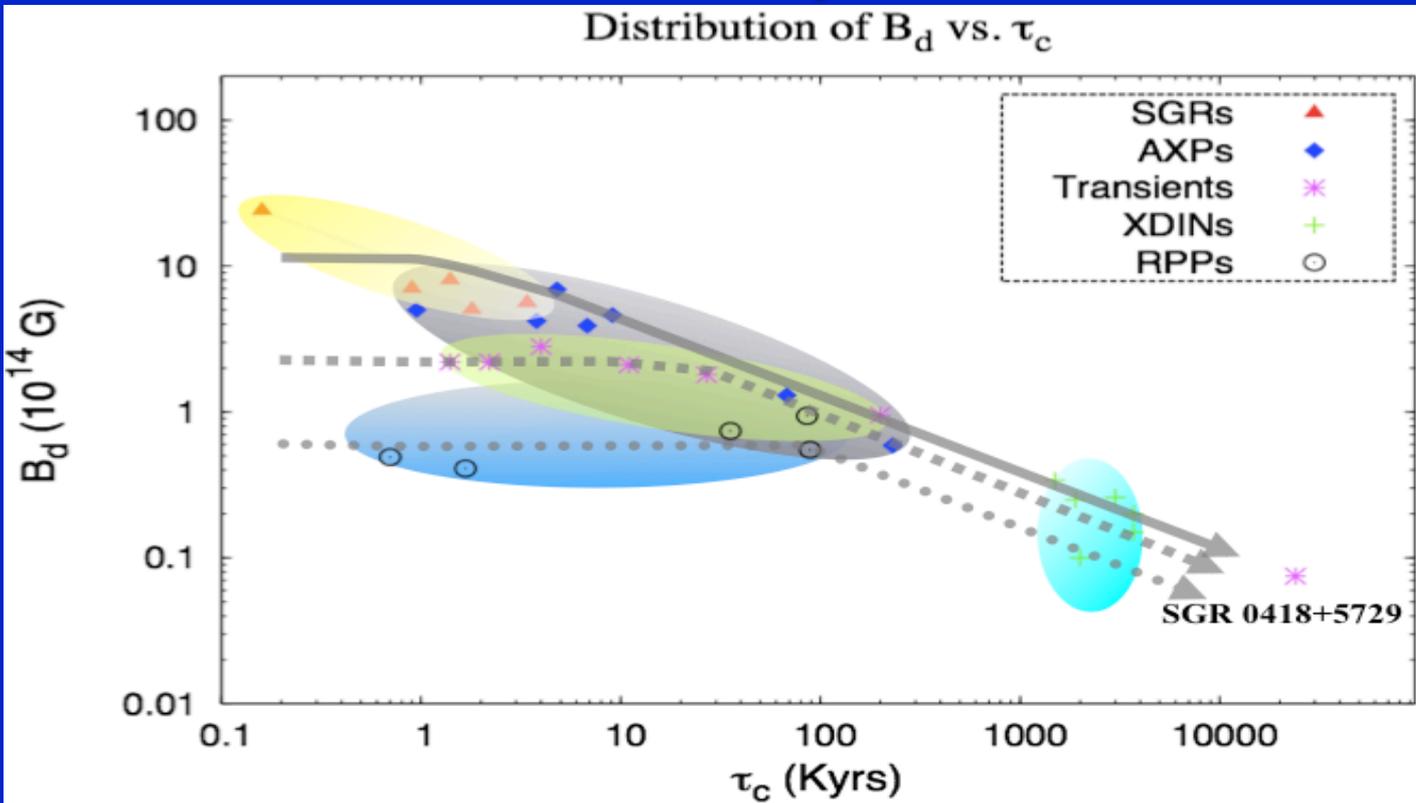
$$B_{int,i} \gtrsim 10^{16} \text{ G}$$



# Adding Sources

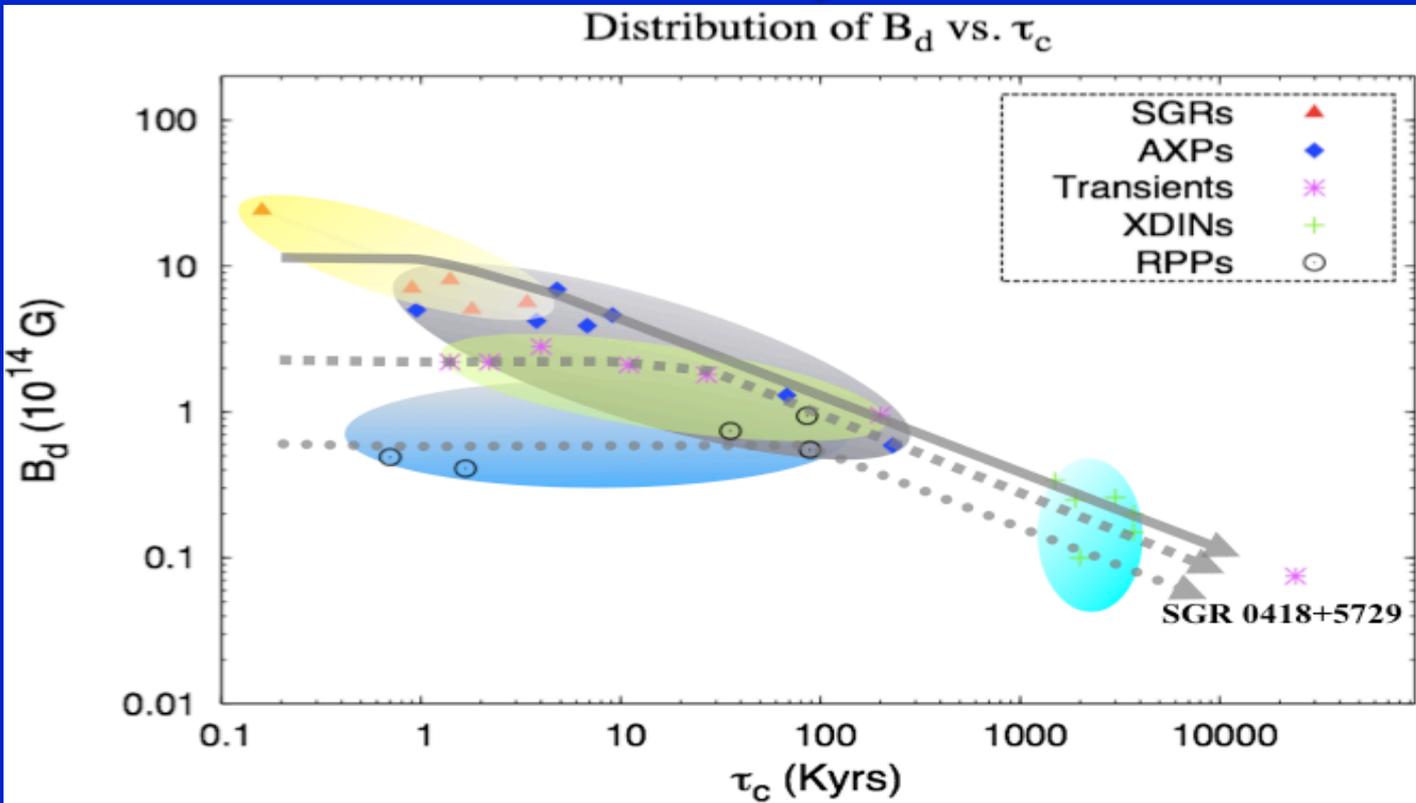


# Evolutionary Links:



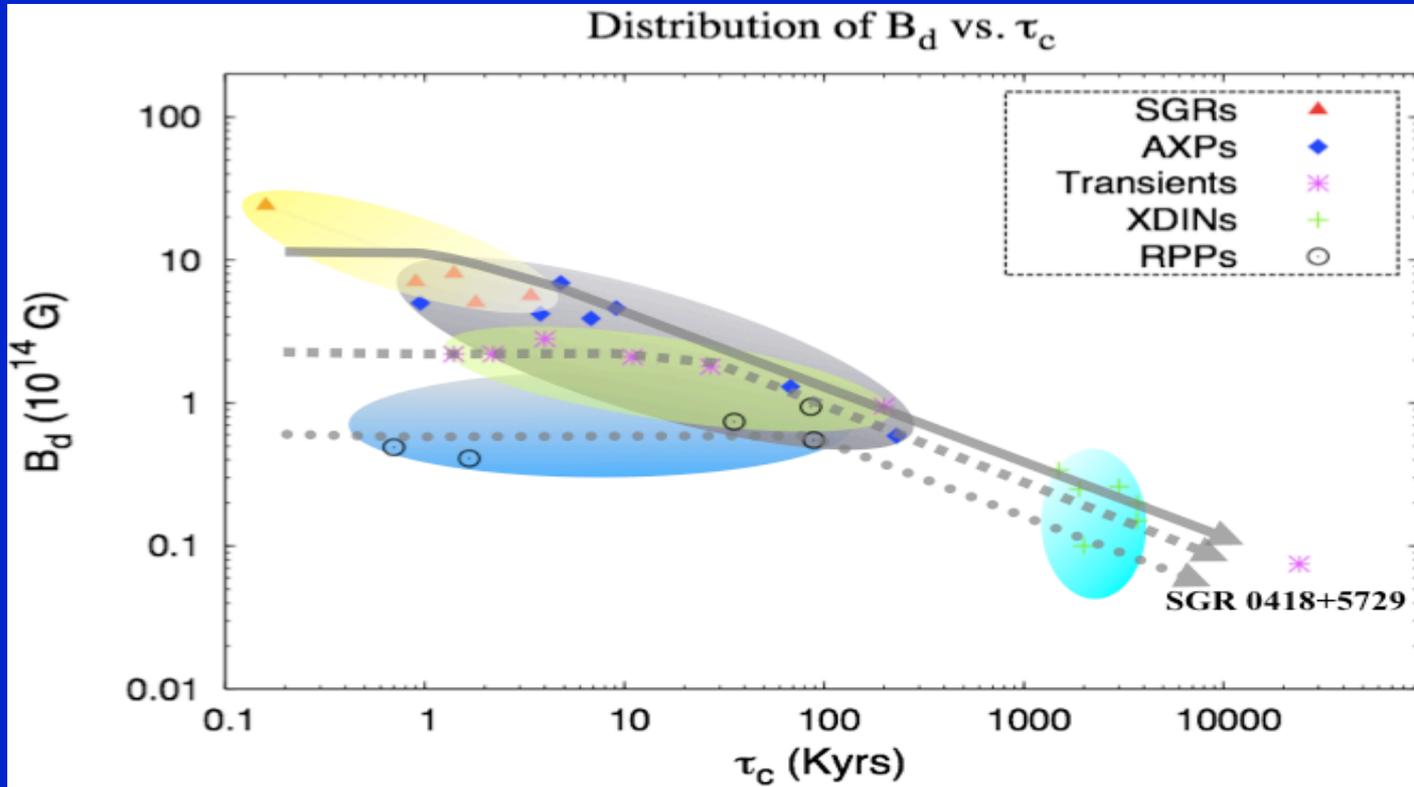
1. SGR/AXP branch (Kouveliotou et al. 1998): SGR  $\rightarrow$  AXP  $\rightarrow$  XDIN?  
 $B_{d,i} \sim 10^{15}$  G,  $B_{int,i} \gtrsim 10^{16}$  G; early  $L_X \sim 10^{35}$  erg/s  $\ll L_{B,int}$  ( $\nu$ -limited)
2. Transient branch: Transient SGR/AXP  $\rightarrow$  ???  
 $B_{d,i} \sim 2 \times 10^{14}$  G,  $L_{X,quiescent} \ll L_{B,dip} \rightarrow B_{int,i} \sim ?$ ,  $L_{X,outburst} \sim L_{X,SGR/AXP}$   
 $\rightarrow ?$  ordered  $B_{int,i} \gtrsim 10^{16}$  G suppresses quiescent heat conduction to surface  
 $\blacklozenge B_{int,SGR/AXP} \rightarrow \alpha\text{-}\Omega$  dynamo,  $B_{int,transient} \rightarrow$  ??? (remnant field?)

# Evolutionary Links:



1. SGR/AXP branch (Kouveliotou et al. 1998): SGR  $\rightarrow$  AXP  $\rightarrow$  XDIN?  
 $B_{d,i} \sim 10^{15}$  G,  $B_{int,i} \gtrsim 10^{16}$  G; early  $L_X \sim 10^{35}$  erg/s  $\ll L_{B,int}$  ( $\nu$ -limited)
2. Transient branch: Transient SGR/AXP  $\rightarrow$  ???  
 $B_{d,i} \sim 2 \times 10^{14}$  G,  $L_{X,quiescent} \ll L_{B,dip} \rightarrow B_{int,i} \sim ?$ ,  $L_{X,outburst} \sim L_{X,SGR/AXP}$   
 $\rightarrow ?$  ordered  $B_{int,i} \gtrsim 10^{16}$  G suppresses quiescent heat conduction to surface
3. High-B RPP  $\rightarrow$  XDIN?

# Sources of interest:



**SGR J0418+5729:**  $B_{d,i} \sim (3-5) \times 10^{14}$  G, while currently  $t_{age} \sim 1-2$  Myr,  
 $B_d \sim (4-7) \times 10^{12}$  G,  $B_{int} \sim (1-2) \times 10^{14}$  G (for  $B_{int,i} \gtrsim 10^{16}$  G) &  
 $L_{B,int} \sim L_X \sim (4-10) \times 10^{30}$  erg/s

**XDINs:**  $t_{age} \sim 0.1-0.6$  Myr,  $B_{d,i} \sim (0.3-20) \times 10^{14}$  G,  $L_X$  likely remnant heat  
( $L_{B,int}$  might contribute a little); no evidence for  $B_{int} \rightarrow ?$  related to  
other high-B NSs without bursting activity (RPP, CCO?)

# Conclusions:

- We find strong observational evidence of dipole field decay on  $\sim 10^3$  yr timescale for strongest magnetars,  $B_{\text{dip},i} \sim 10^{15}$  G
- Dipole field decay index  $\alpha$ , defined by  $\dot{B}_d \propto B_d^{1+\alpha}$ , is in the range  $1 \leq \alpha < 2$ , and more likely  $1.5 \lesssim \alpha \lesssim 1.8$
- Once  $B_{\text{dip}}$  decays significantly, true age  $\ll$  spin-down age
- $L_{\text{X,persistent}} > |\dot{E}_{B,\text{dipole}}| \rightarrow$  another energy source is required  
 $\rightarrow$  likely internal field decay: requires  $B_{\text{internal},i} \gtrsim 10^{16}$  G
- Evolutionary tracks:
  - ◆ SGR/AXP branch
  - ◆ Transient branch
  - ◆ High-B RPP  $\rightarrow$  XDIN branch?