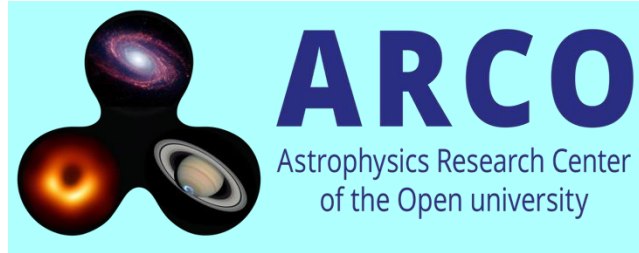


Explosive Astrophysical Transients

in the Multi-Messenger Era: Some Highlights



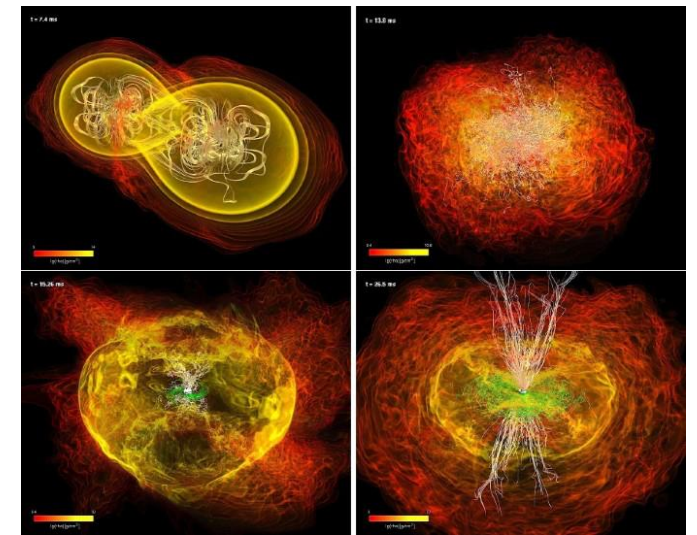
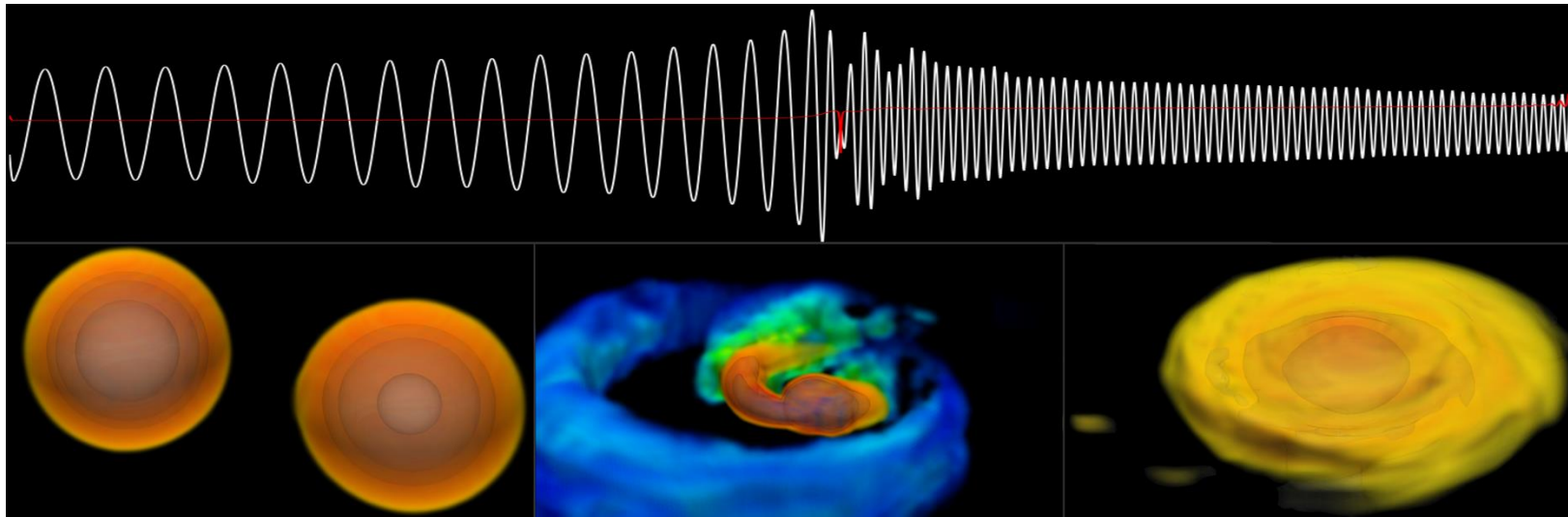
**Jonathan
Granot**



THE GEORGE
WASHINGTON
UNIVERSITY
WASHINGTON DC



Open University of Israel & **George Washington University** & **YITP**



YITP long-term workshop: Multi-Messenger Astrophysics in the Dynamic Universe

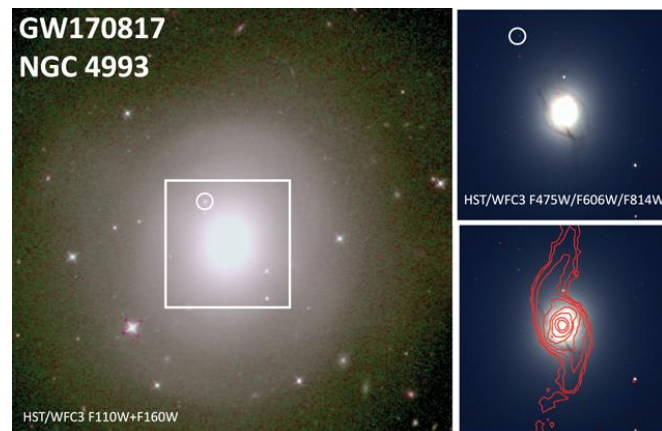
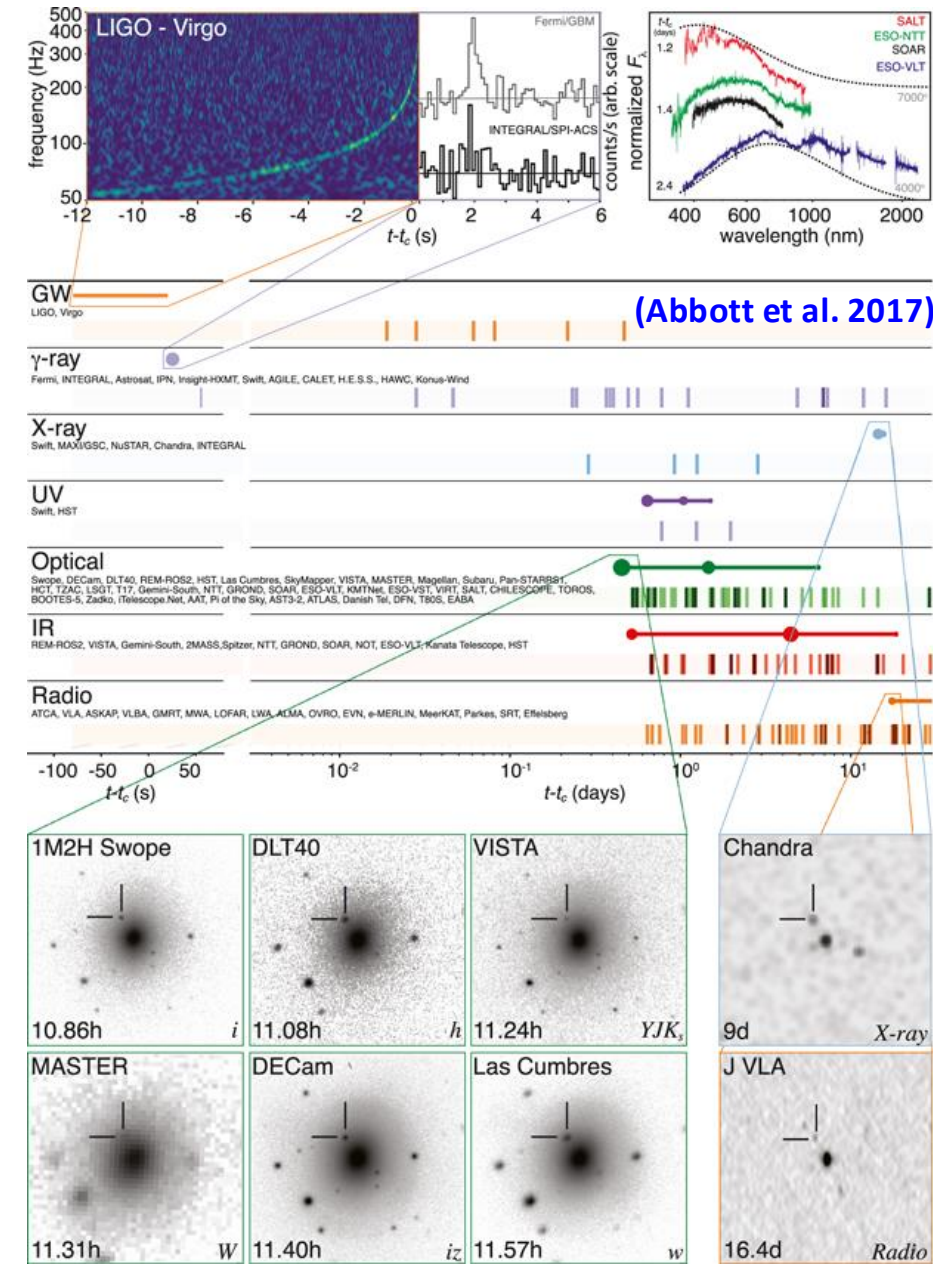
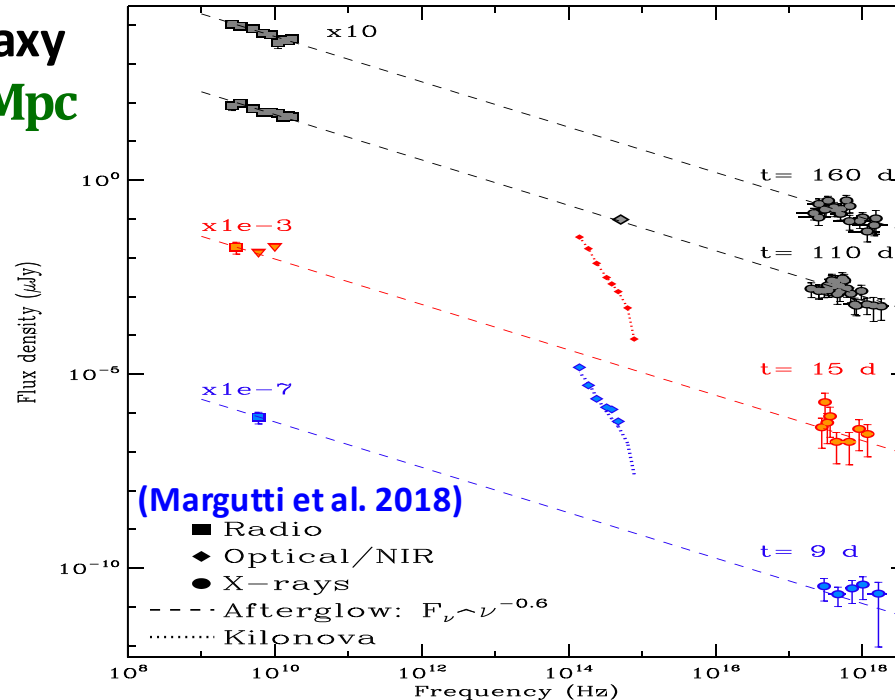
Yukawa Institute for Theoretical Physics, Kyoto University, Kyoto, Japan, 26 January 2026

Outline of the Talk: what will or will not be covered

- **GRBs:** select highlights from the recent decade
 - ◆ **GRB170817A/GW170817:** 1st EM-GW counterpart (sGRB, kilonova, afterglow), r-process elements synthesis, jet angular structure & large off-axis viewing angle
 - ◆ **Polarization:** GW170817 UL \Rightarrow constraints on B-field structure in the afterglow shock; GRB180720B: reverse to forward shock transition; afterglow pol. from structured jets
 - ◆ **Prompt GRB:** dissipation (IS/rec), emission (syn/Compt/SSC/had) (Rahaman)
 - ◆ **TeV emission:** from nearby GRBs – afterglow, reverse shock, prompt? (???)
 - ◆ **GRB 221009A (B.O.A.T):** bright in TeV, shallow jet, 6-12 MeV line (Salafia)
- **Einstein Probe** X-ray transients: mostly GRBs, start earlier, last longer (Hamidani)
- **ULGRB 250702B:** ~day long, unclear origin (some type of collapsar or TDE?)
- **Magnetars / FRBs:**
 - ❖ FRBs from a Galactic magnetar, SGR 1935+2154 (28.4.2020)
 - ◆ Persistent Radio Sources (**PRSs**) associated with repeating Fast Radio Bursts (**FRBs**)
 - ◆ **Extragalactic Magnetar Giant Flares**- current sample and prospects

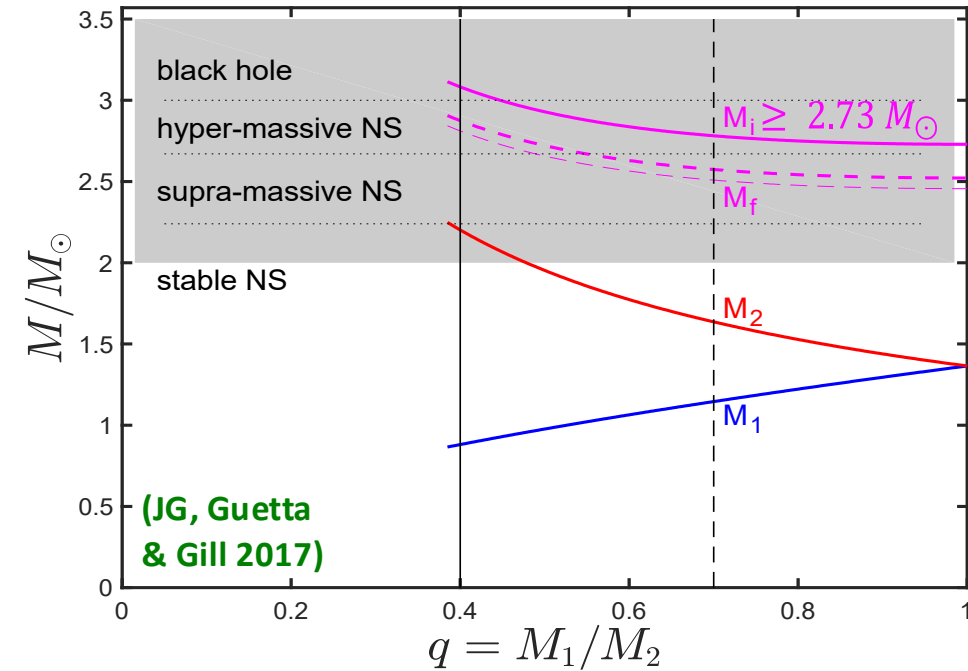
GW170817 / GRB170817A: NS-NS merger

- First electromagnetic counterpart to a GW event
 - The **short GRB** 170817A (very under-luminous, 1.74 s γ -GW delay)
 - Optical (IR to UV) **kilonova** emission (**1st clear-cut**) for a few weeks
 - X-ray (> 9 d; still barely detected) to radio (>16 d) **afterglow**
- First NS-NS merger detected in gravitational waves (GW)
- First direct sGRB - NS-NS merger association (Eichler+ 1989)
- The γ -GW 1.74 s delay constrains GW speed: $\left| \frac{v_{\text{GW}}}{c} - 1 \right| \lesssim 4 \cdot 10^{-16}$
- $D_{\text{GW}} = 43^{+2.9}_{-6.9}$ Mpc; host galaxy is elliptical: $D = 41.0 \pm 3.1$ Mpc ($z = 0.009783$) 2 kpc from host center in projection

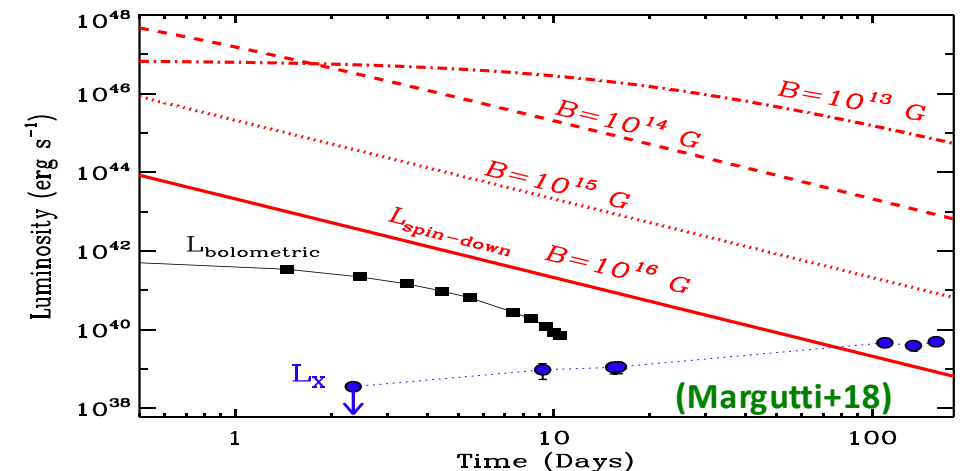


GW170817 / GRB170817A: Remnant Type

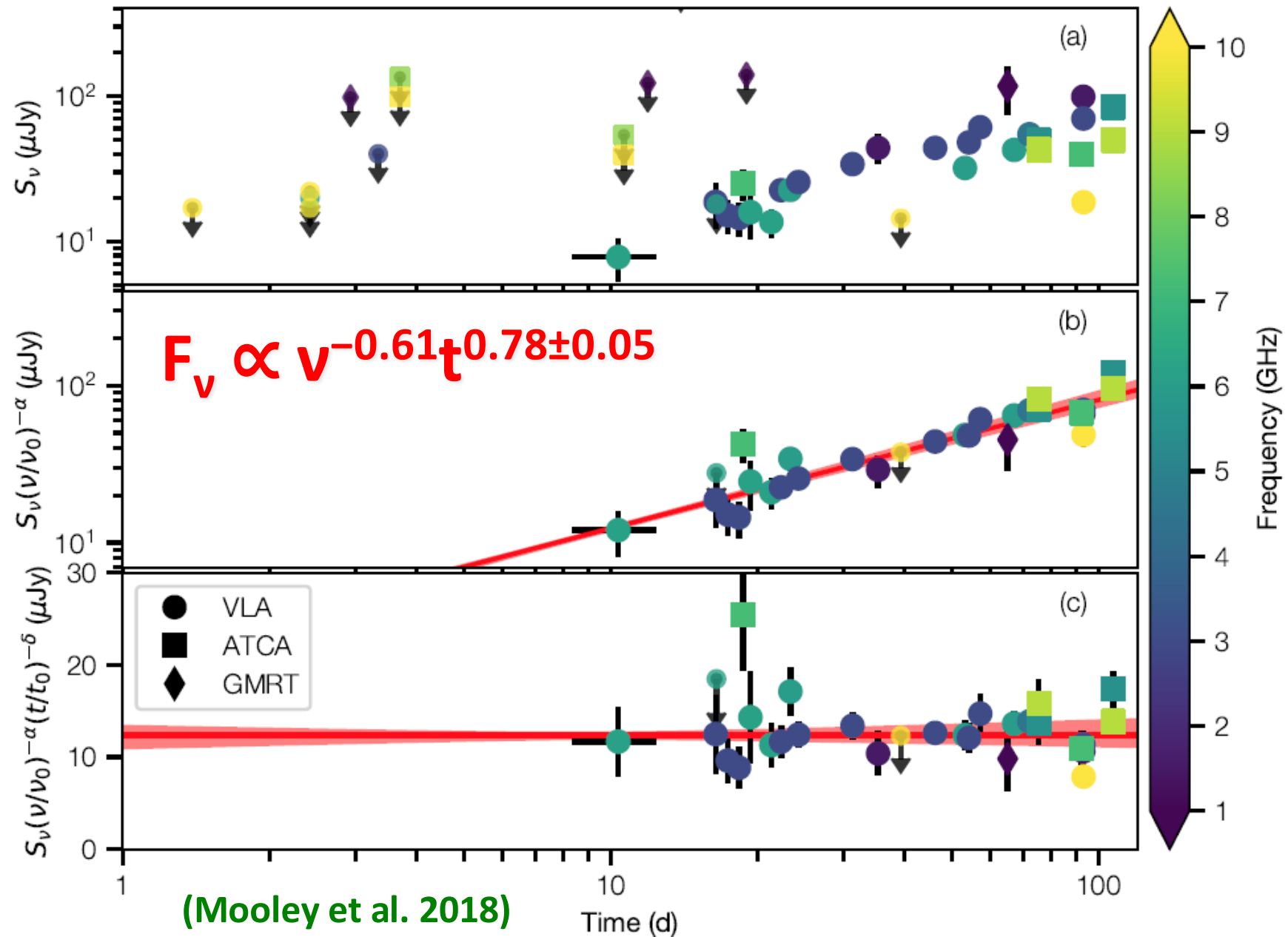
- $M_{1,2}$ = pre-merger NS $M_{\text{gravitational}}$
- post-merger total mass: $M_i = M_1 + M_2$
- Final mass $M_f \approx 0.93M_i$ due to:
 - ❖ GW & neutrino energy losses
 - ❖ Mass ejection during the merger
- A stable NS or SMNS $\Rightarrow P_0 \approx 1 \text{ ms} \Rightarrow E_{\text{rot}} \gtrsim 10^{52.5} \text{ erg}$,
 $\tau_{\text{sd}} \approx 20B_{13}^{-2} \text{ days} \Rightarrow$ would contradict afterglow observations (also what produces the GRB/afterglow?)
- The argument can be reversed to constrain NS EoS & $M_{\text{max}} \lesssim 2.17M_{\odot}$ (Margalit & Metzger 2017; Rezzolla et al. 2018)



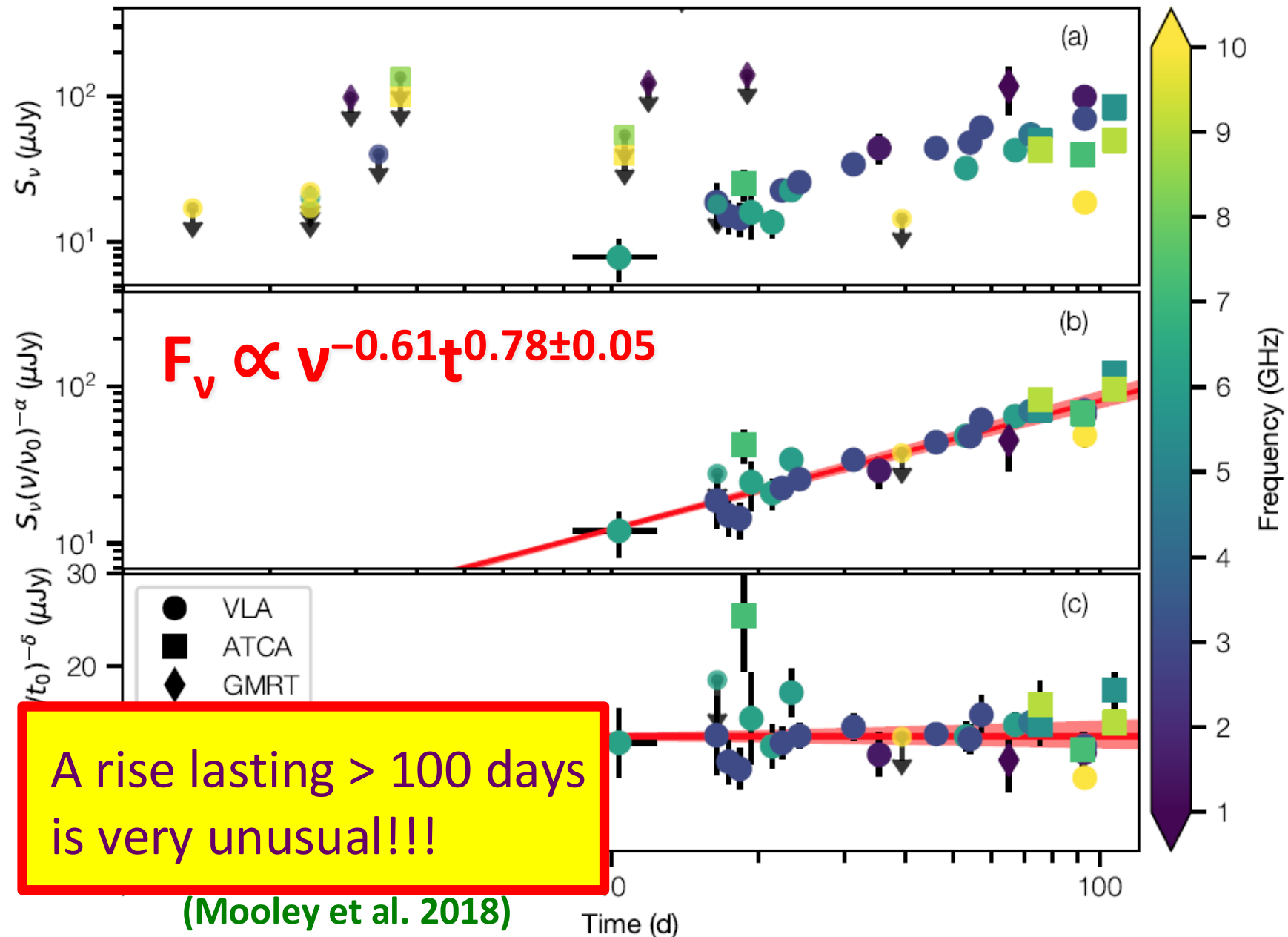
Chirp mass: $\mathcal{M} = \left(\frac{M_1^3 M_2^3}{M_1 + M_2} \right)^{1/5} = 1.188^{+0.004}_{-0.002} M_{\odot}$ (Abbott+ 2017)



GRB170817A: Afterglow Observations



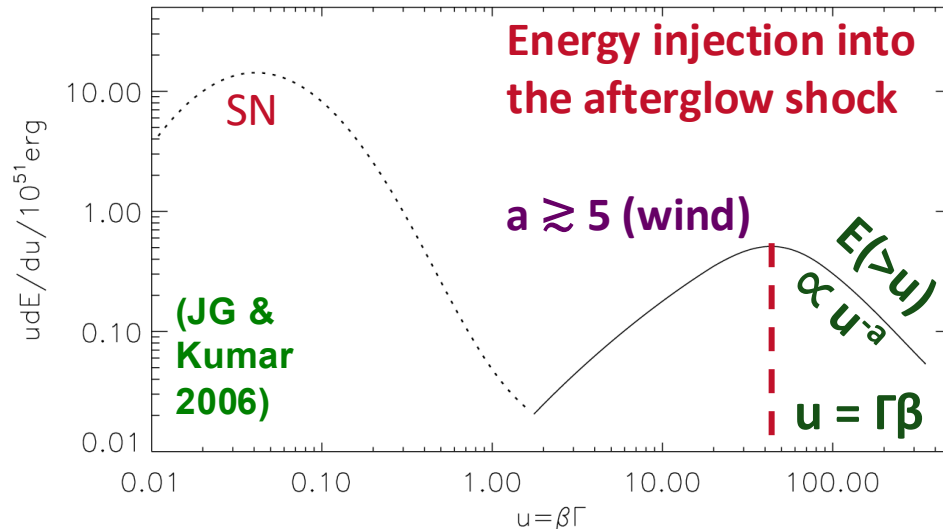
GRB170817A: Afterglow Observations



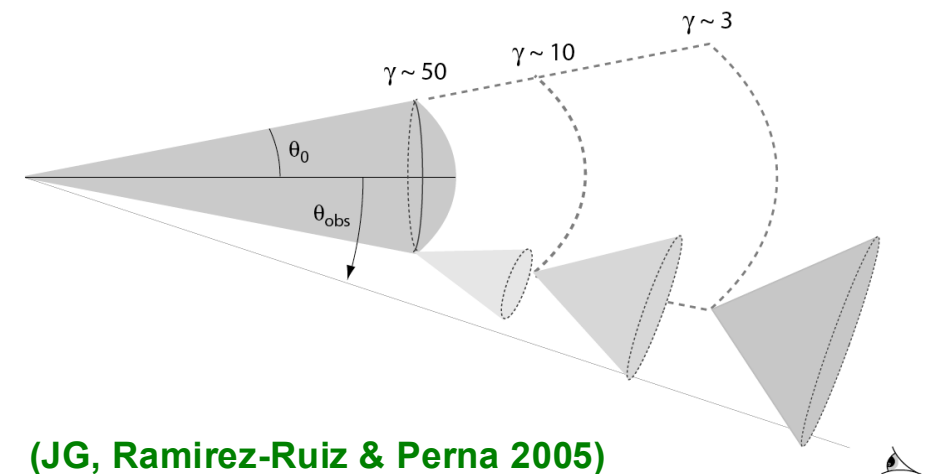
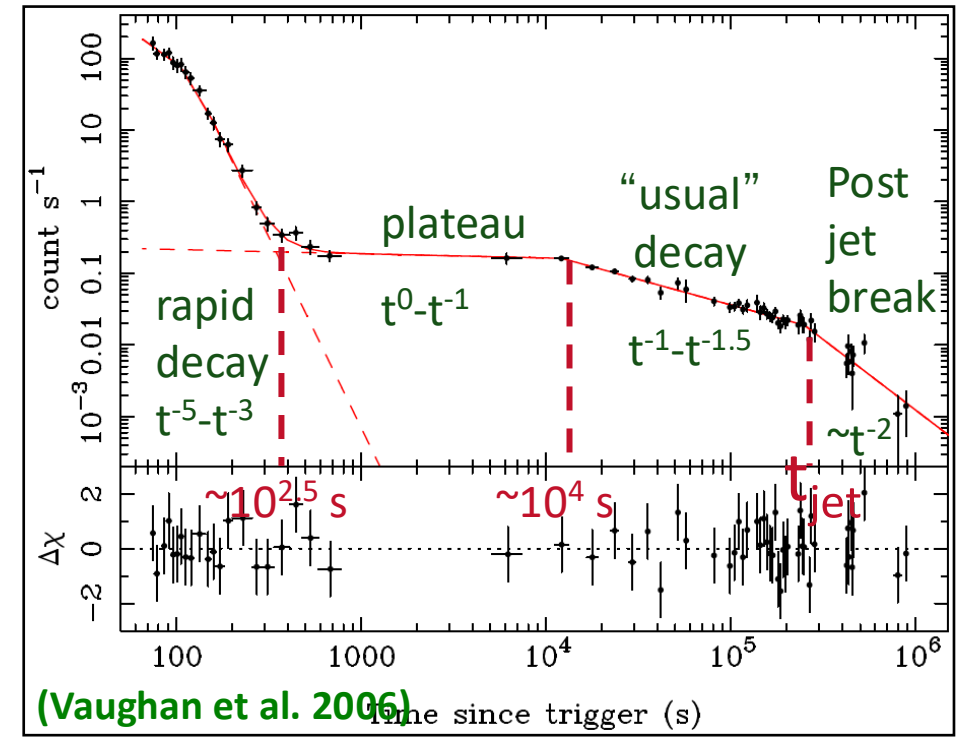
Analogy to rising F_ν : X-ray Plateaus

■ Possible solutions:

- ◆ Evolution of shock microphysical parameters (JG, Konigl & Piran 2006)
- ◆ Energy injection into external shock:
 1. long-lived relativistic wind
 2. slower ejecta catching up
(Sari & Meszaros 00; Nousek+ 06; JG & Kumar 06)



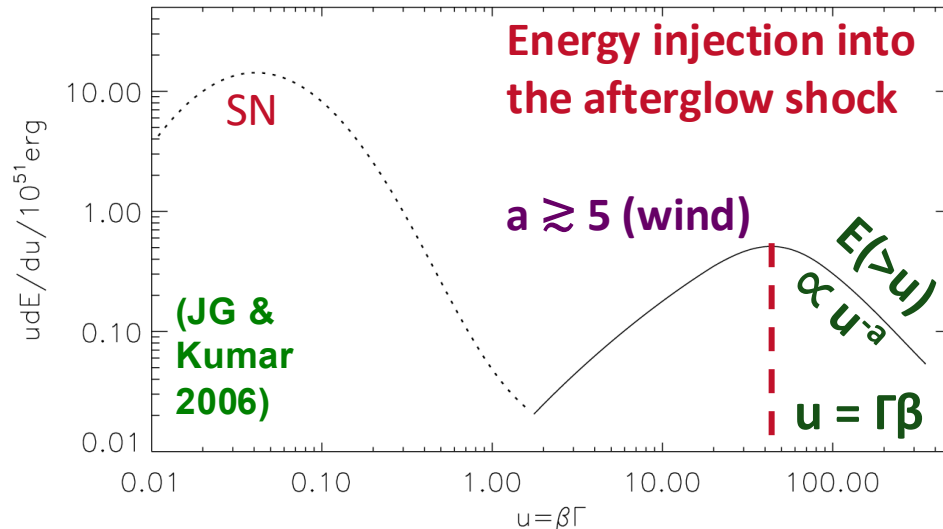
- ◆ Viewing angle effects



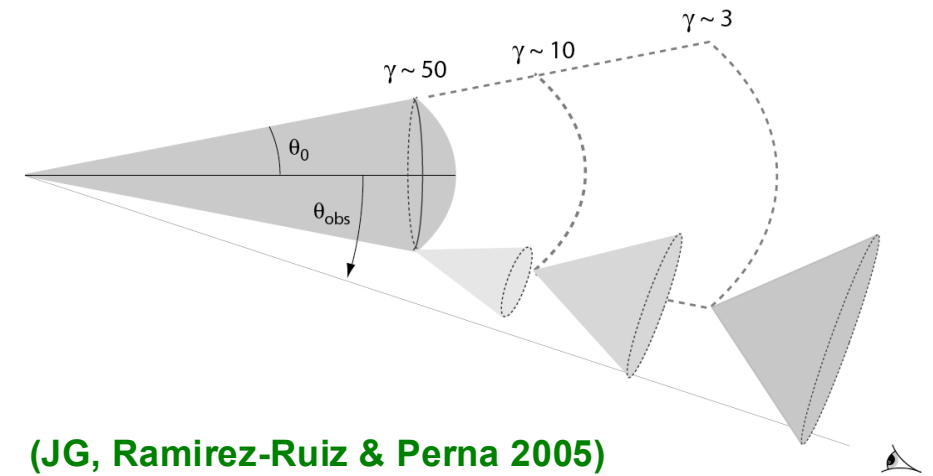
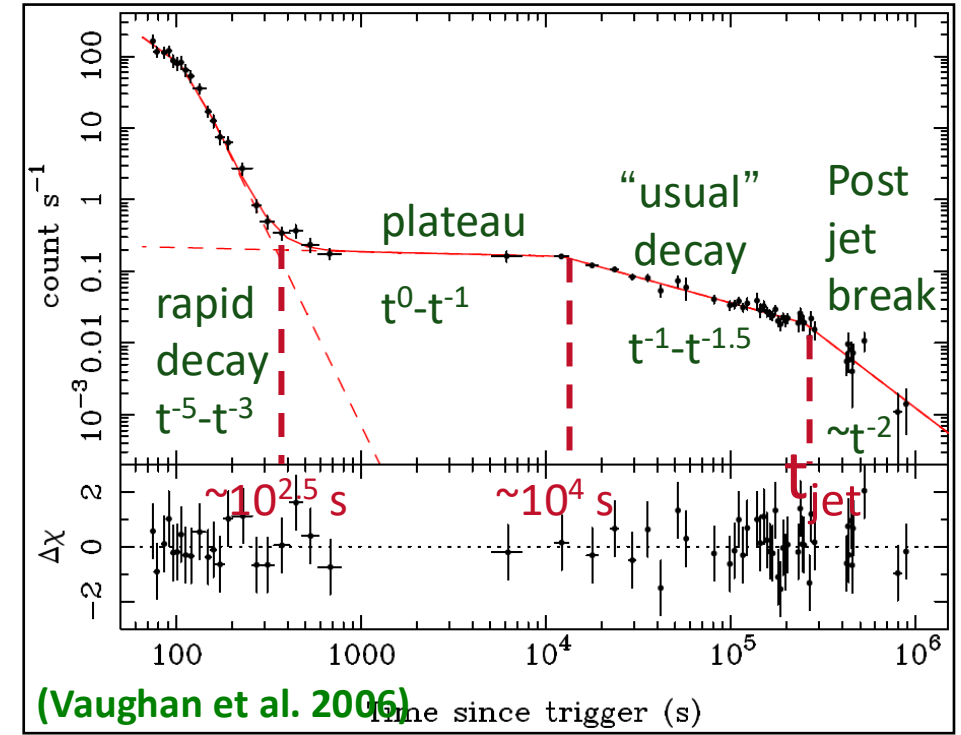
Analogy to rising F_ν : X-ray Plateaus

■ Possible solutions:

- ◆ Evolution of shock microphysical parameters (JG, Konigl & Piran 2006)
- ◆ Energy injection into external shock:
 1. long-lived relativistic wind
 2. slower ejecta catching up **radial**
(Sari & Meszaros 00; Nousek+ 06; JG & Kumar 06)



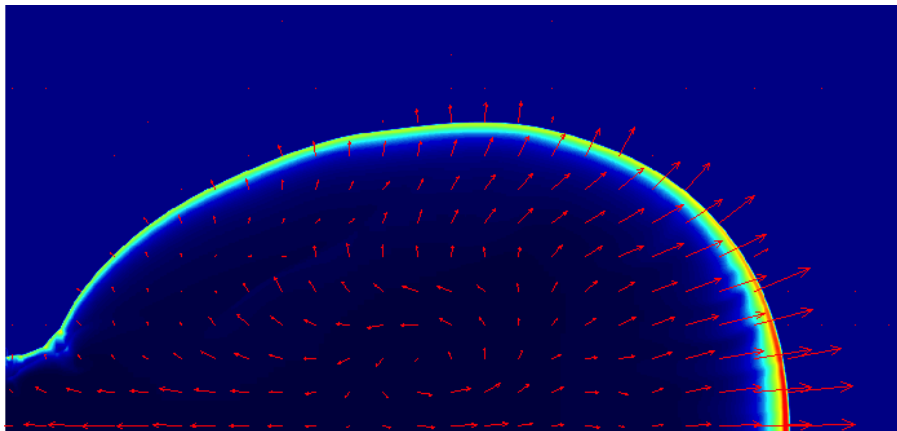
- ◆ **Viewing angle effects** **angular**



(JG, Ramirez-Ruiz & Perna 2005)

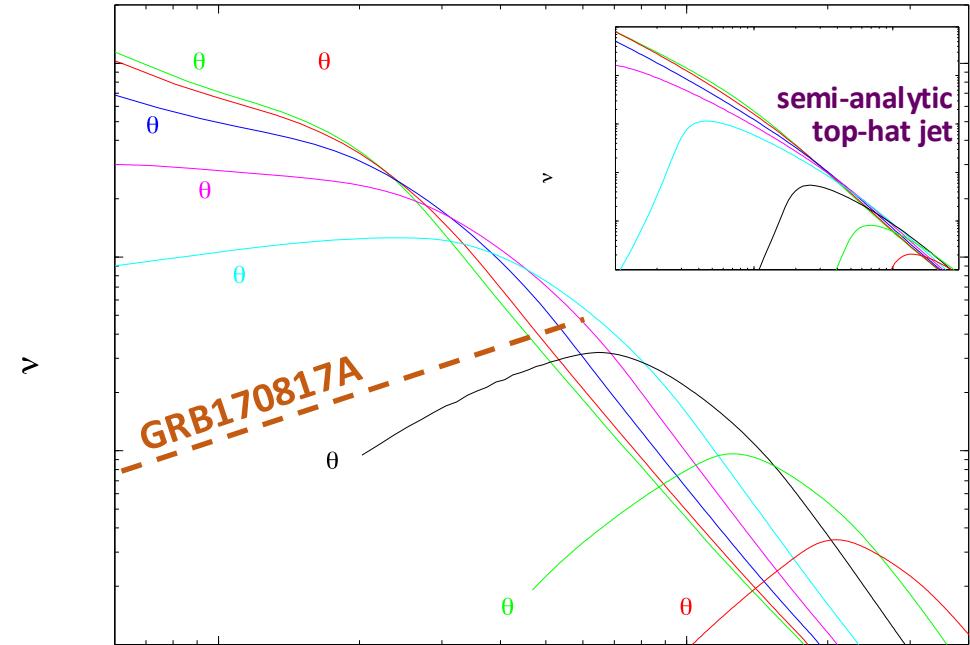
Analogy to rising F_ν : Off-Axis Viewing

- The emission is initially strongly beamed away from our L.o.S
- F_ν rises as beaming cone widens
- When beaming cone reaches LoS F_ν peaks & approaches on-axis F_ν
- The rise is much more gradual for hydrodynamic simulations due to slower matter at the jet's sides with non-radial velocities

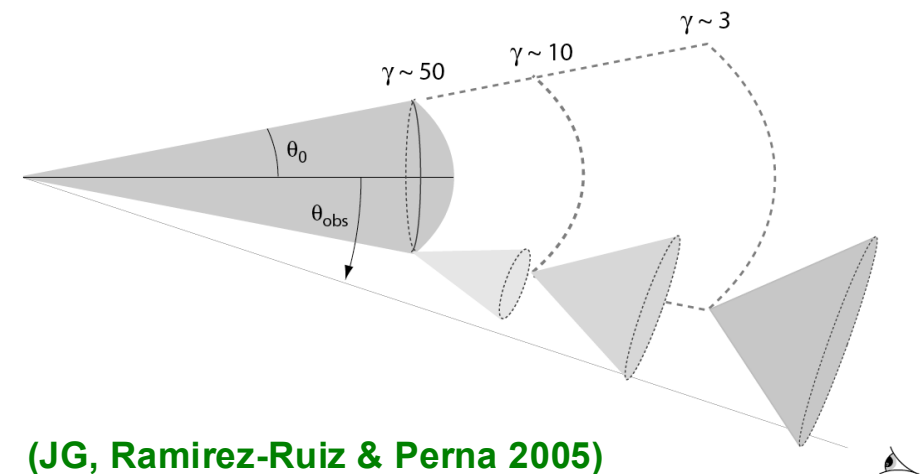


(JG et al. 2001)

0.311E-12 0.163E-06



(JG et al. 2002)

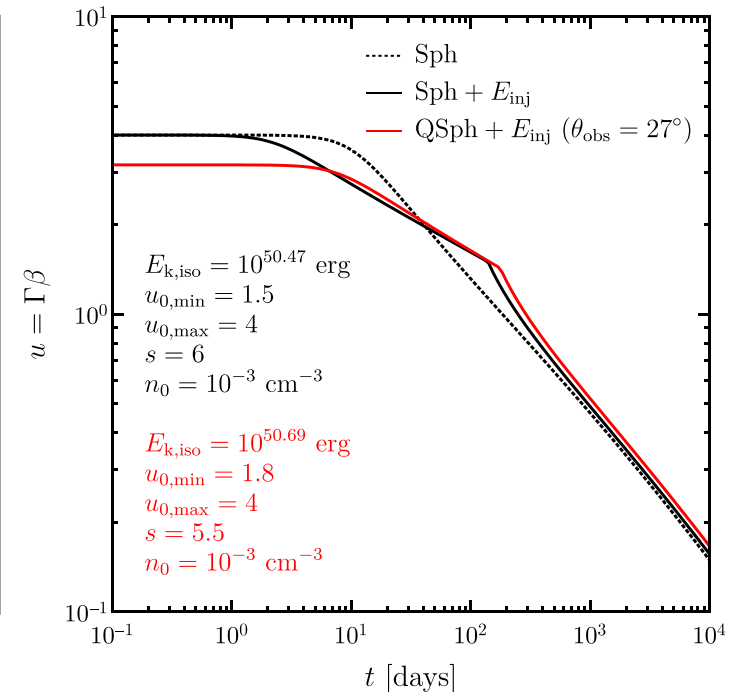
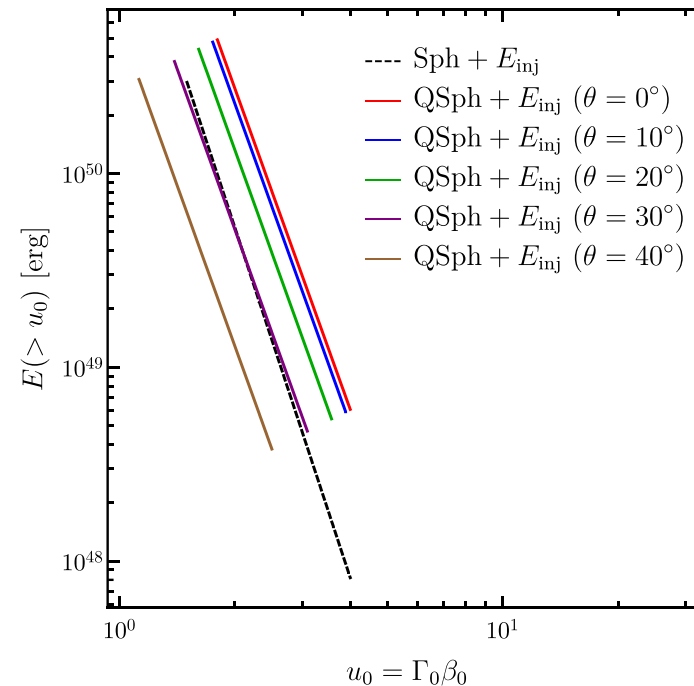


(JG, Ramirez-Ruiz & Perna 2005)

Outflow Structure: Breaking the Degeneracy (Gill & JG 18)

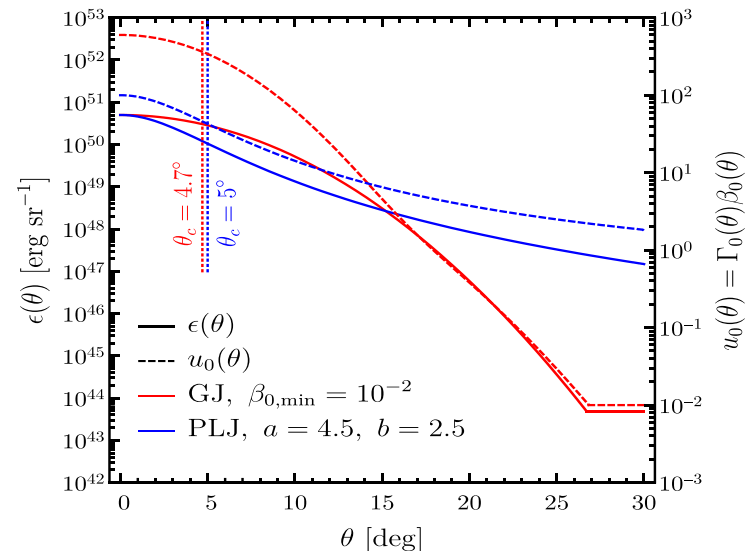
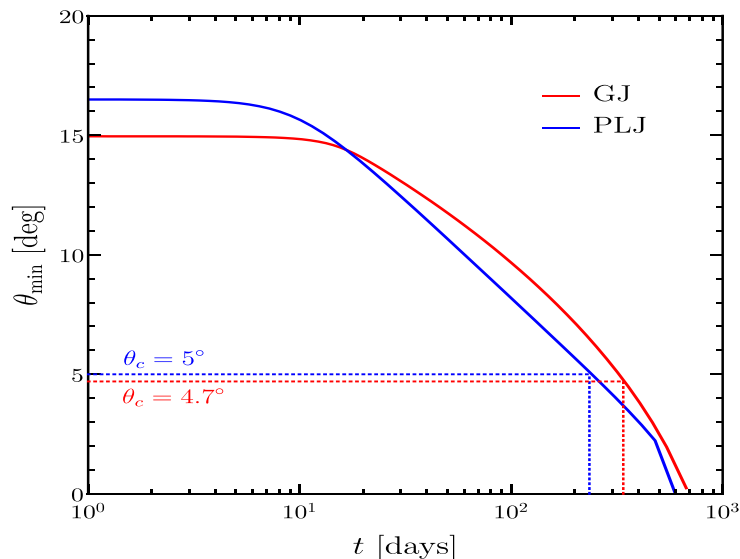
- The lightcurves leave a lot of degeneracy between models
- The degeneracy may be lifted by calculation the afterglow **images** & **polarization** (e.g. Nakar & Piran 2018; Nakar et al. 2018)
- We considered 4 different models including both main types
 - ◆ Sph+E_{inj}: Spherical with energy injection $E(>u) \propto u^{-6}$, $1.5 < u < 4$
 - ◆ QSph+E_{inj}: Quasi-Spherical + energy injection $E(>u) \propto u^{-s}$, $u_{\min,0} = 1.8$ $u_{\max,0} = 4$, $s = 5.5$, $\zeta = 0.1$

$$\frac{\epsilon(\theta)}{\epsilon_0} = \frac{u_{0,\min}(\theta)}{u_{\min,0}} = \frac{u_{0,\max}(\theta)}{u_{\max,0}} = \frac{\zeta + \cos^2 \theta}{\zeta + 1}$$



Outflow Structure: Breaking the Degeneracy (Gill & JG 18)

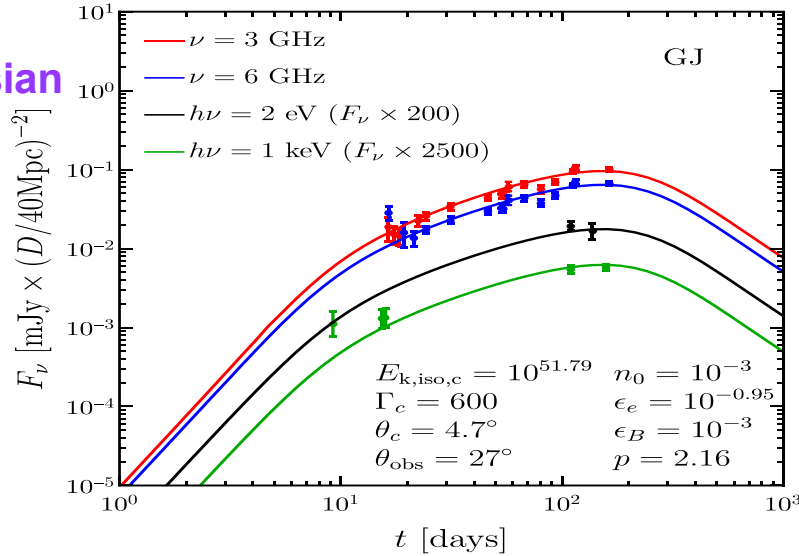
- The lightcurves leave a lot of degeneracy between models
- The degeneracy may be lifted by calculation the afterglow **images** & **polarization** (e.g. Nakar & Piran 2018; Nakar et al. 2018)
- We considered 4 different models including both main types
 - ◆ GJ: Gaussian Jet (in $\epsilon = dE/d\Omega$, Γ_0-1) $\Gamma_c = 600$, $\theta_c = 4.7^\circ$
 - ◆ PLJ: Power-Law Jet; $\epsilon = \epsilon_c \Theta^{-a}$, $\Gamma_0-1 = (\Gamma_c-1)\Theta^{-b}$, $\Theta = [1+(\theta/\theta_c)^2]^{1/2}$, $\Gamma_c = 100$, $\theta_c = 5^\circ$, $a = 4.5$, $b = 2.5$
- As there is a lot of freedom we fixed: $p=2.16$, $\epsilon_B = n_0 = 10^{-3}$, $\theta_{\text{obs}} = 27^\circ$



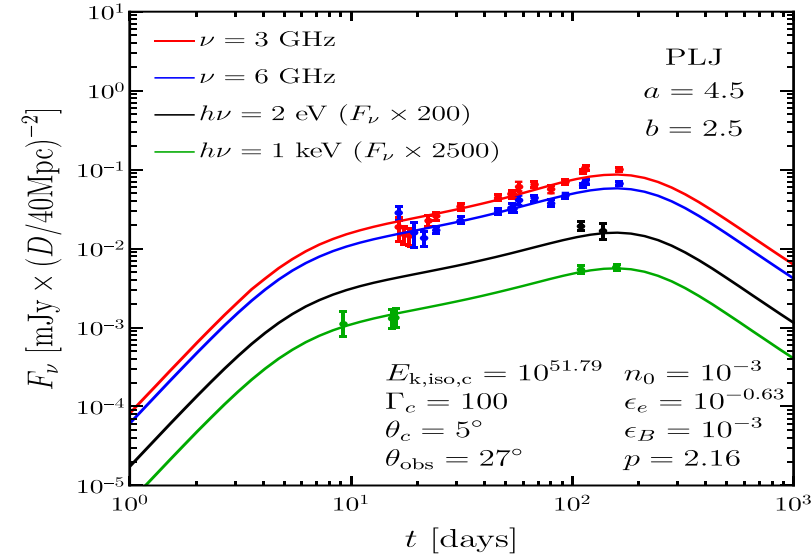
Outflow Structure: Breaking the Degeneracy (Gill & JG 18)

■ Tentative fit to GRB170817A afterglow data (radio to X-ray)

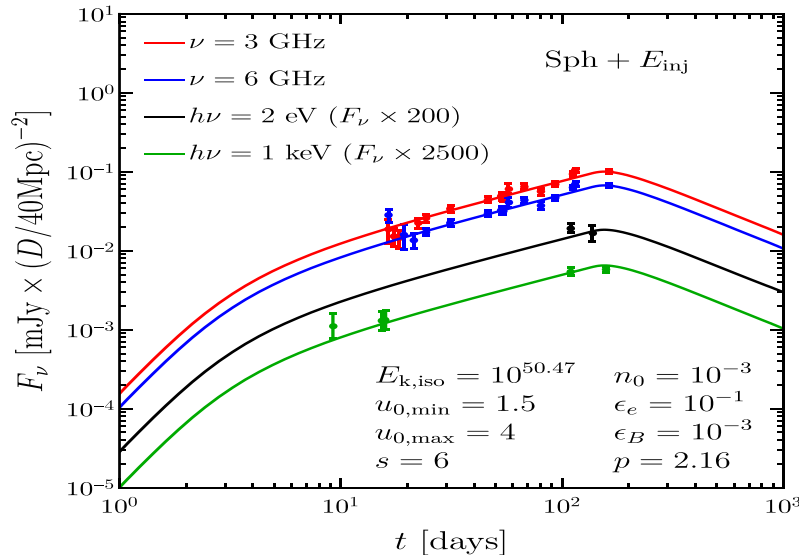
Gaussian
Jet



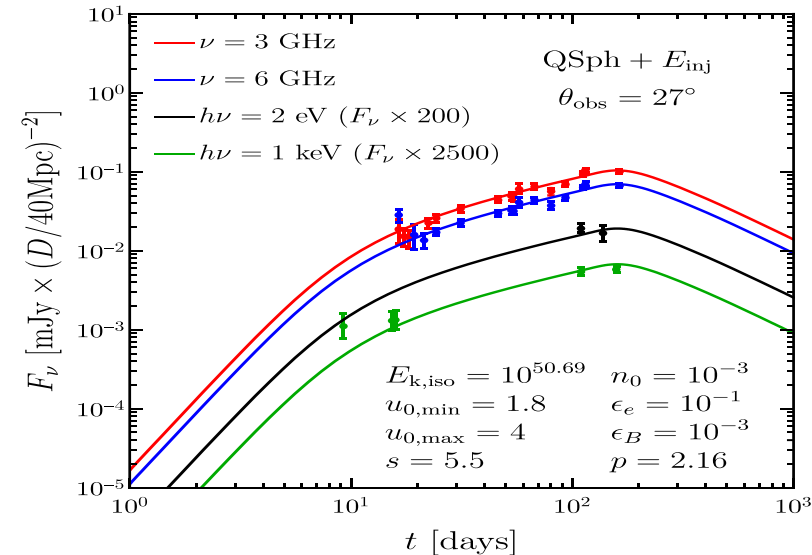
Power
Law
Jet



Spherical
+
Energy
Injection



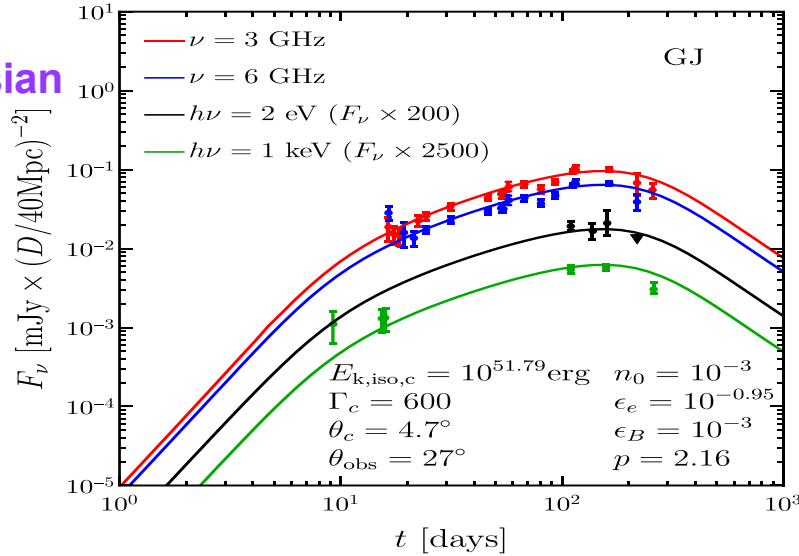
Quasi
Spherical
+
Energy
Injection



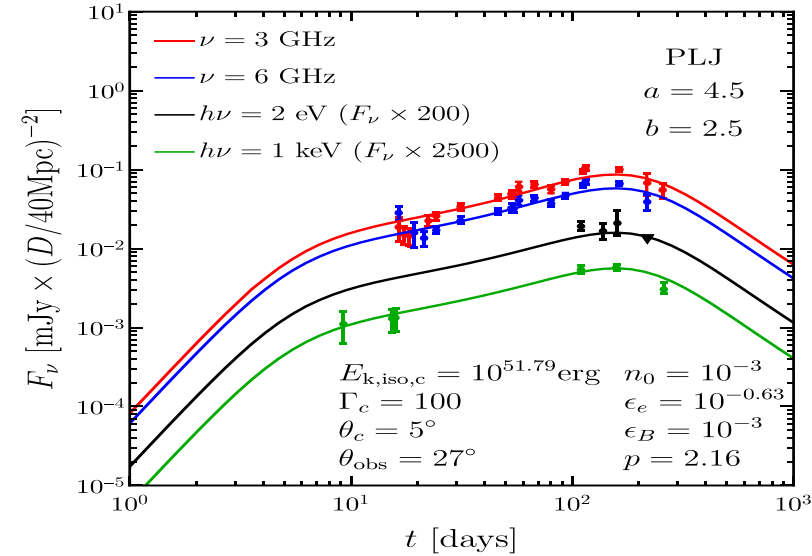
Outflow Structure: Breaking the Degeneracy (Gill & JG 18)

- New data that came out established a peak at $t_{\text{peak}} \sim 150$ days

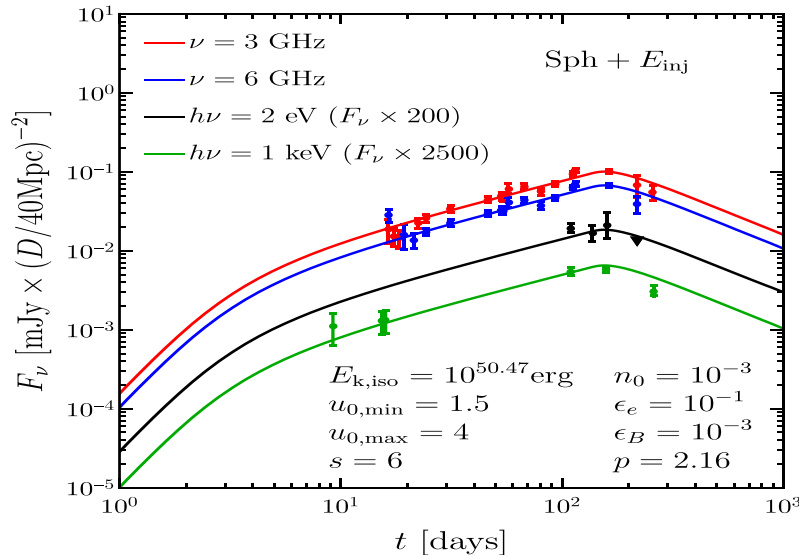
Gaussian
Jet



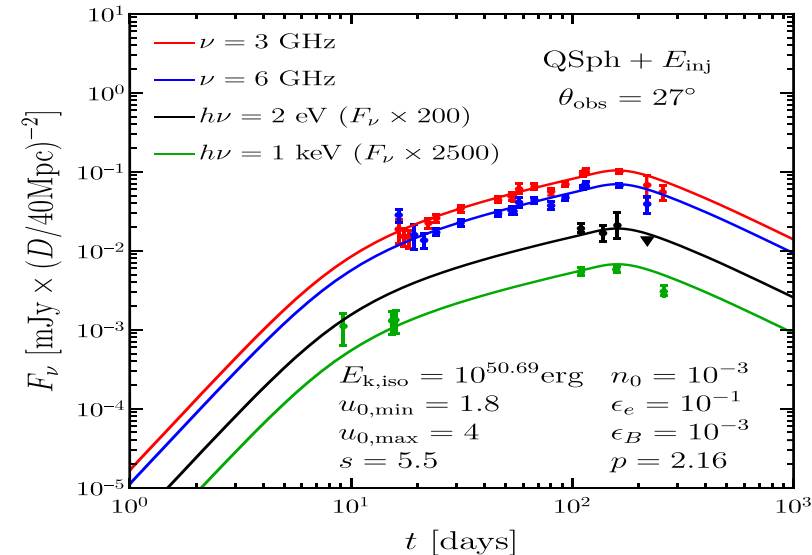
Power
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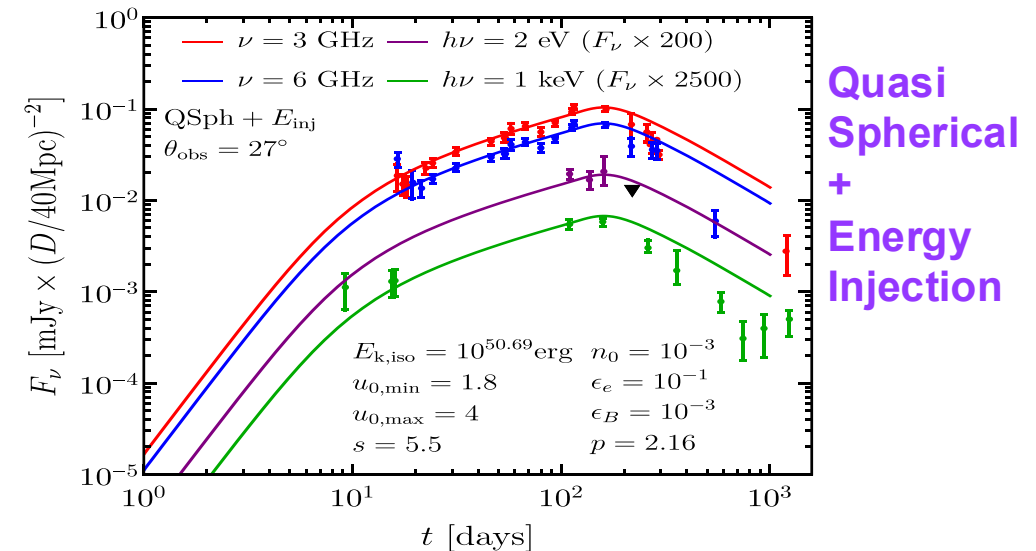
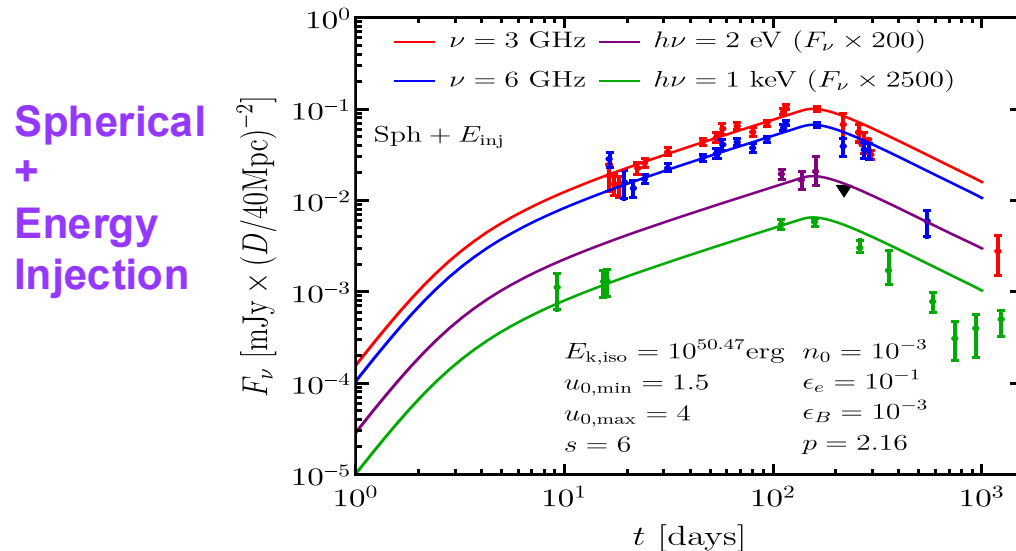
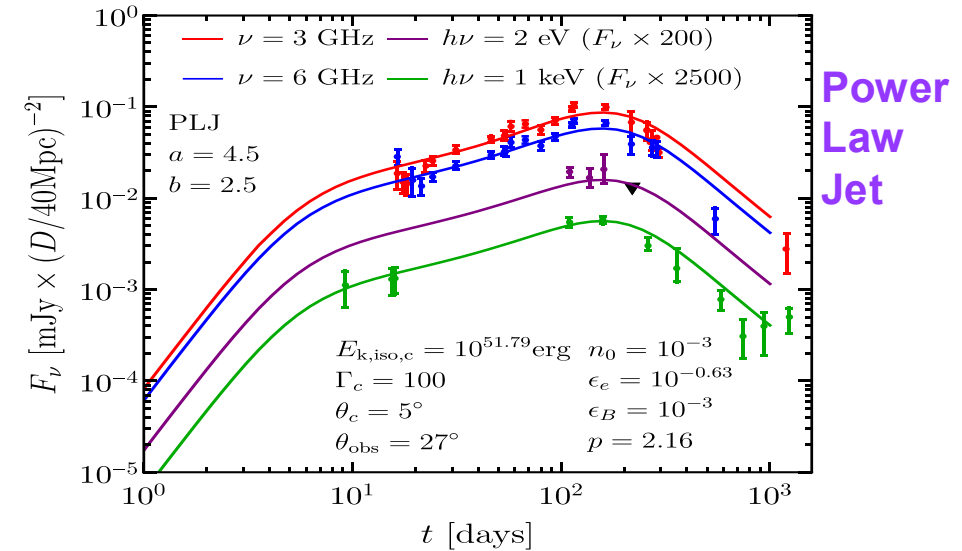
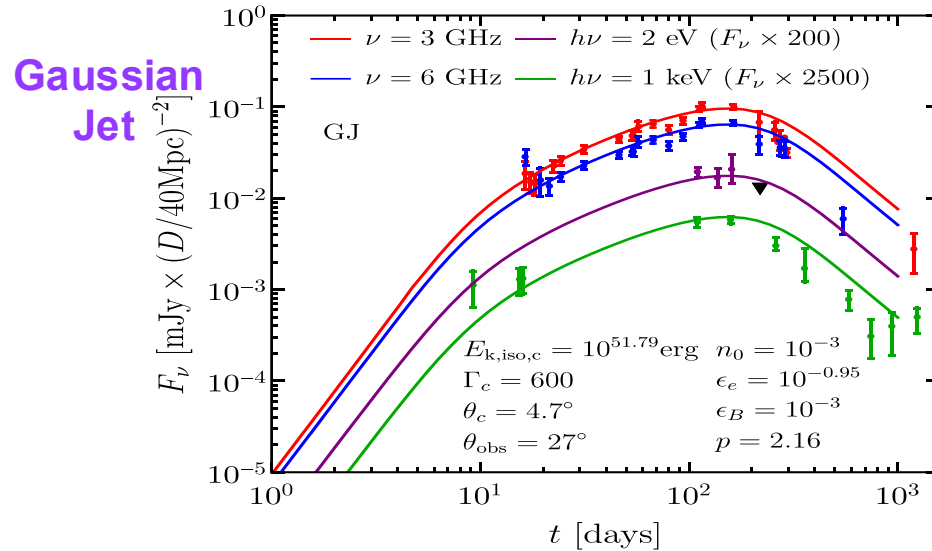


Quasi
Spherical
+
Energy
Injection



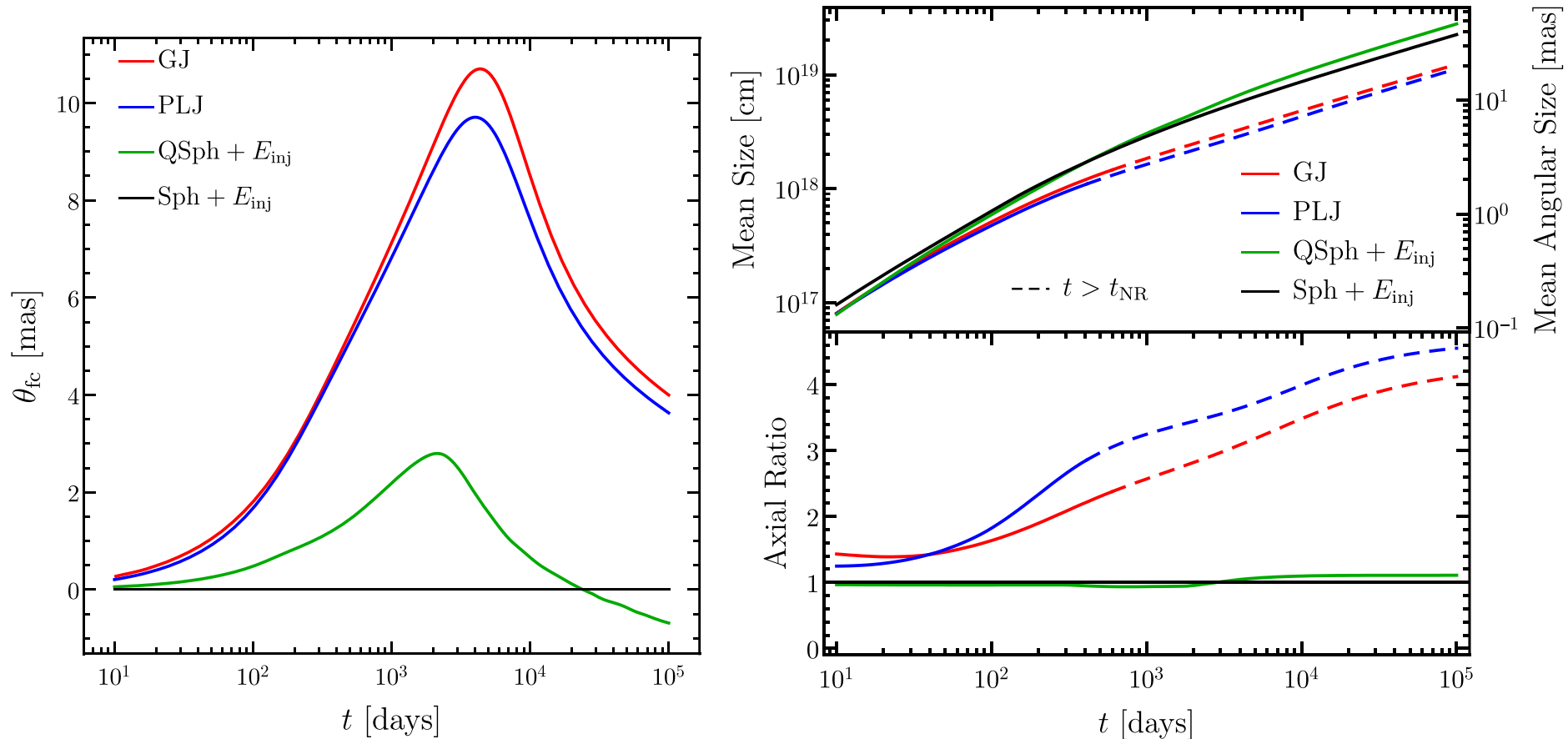
Outflow Structure: Breaking the Degeneracy (Gill & JG 18)

- The jet models decay faster (closer to post-peak data: $F_\nu \propto t^{-2.2}$)



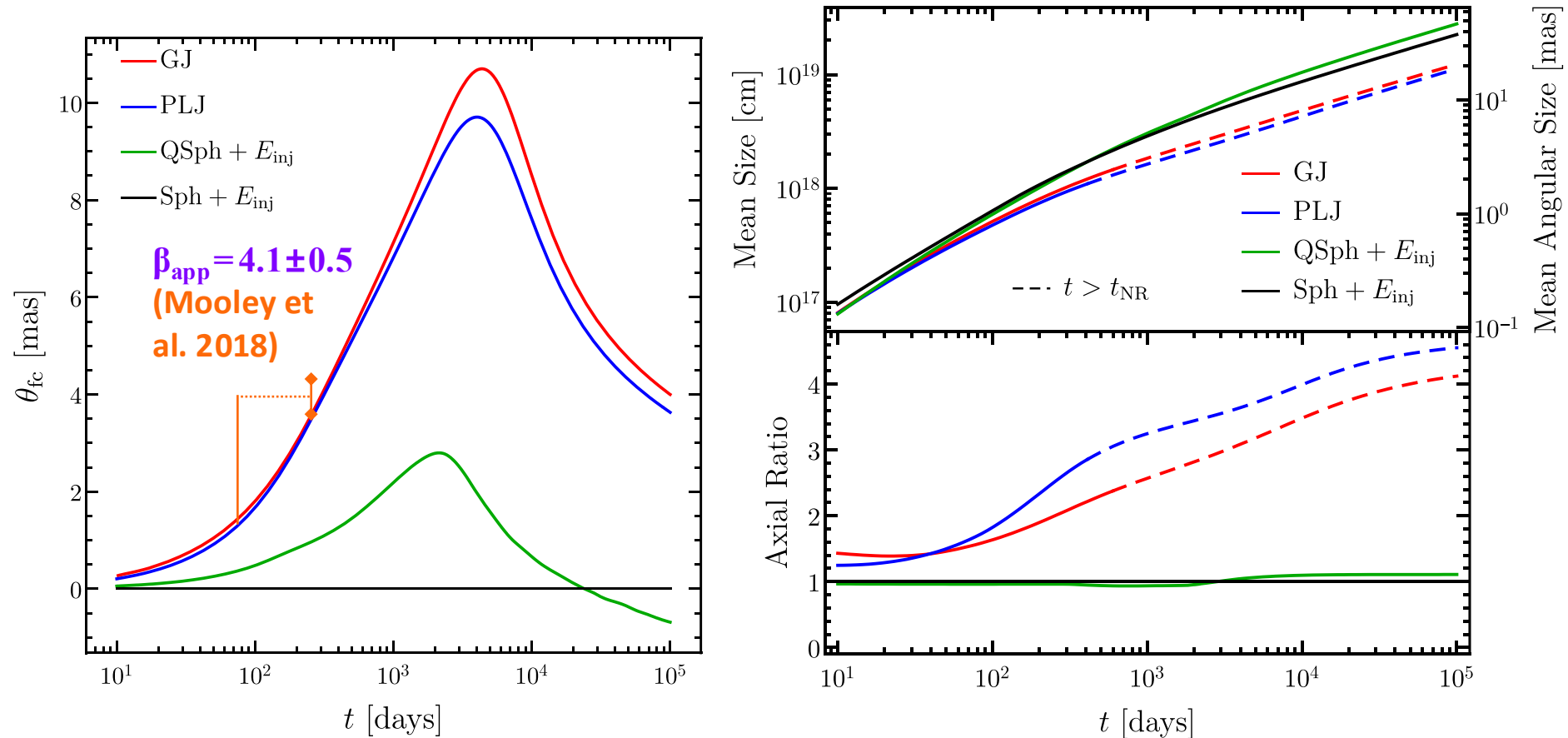
Afterglow Images: Flux Centroid, Size, Shape

- The flux centroid motion: a potentially powerful diagnostic
- It may be hard to tell apart models based on the image size alone, but a much higher axis-ratio is expected for jet models



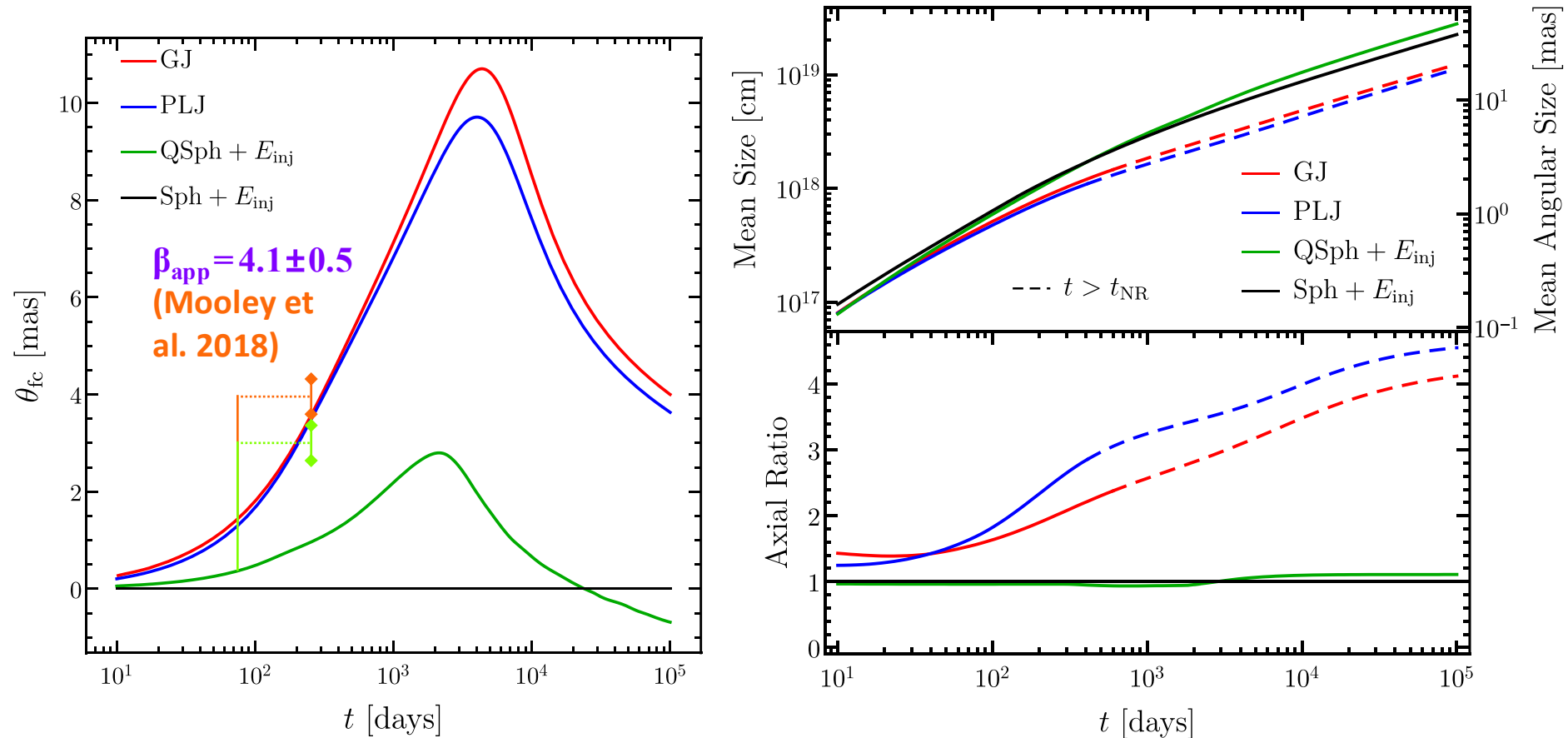
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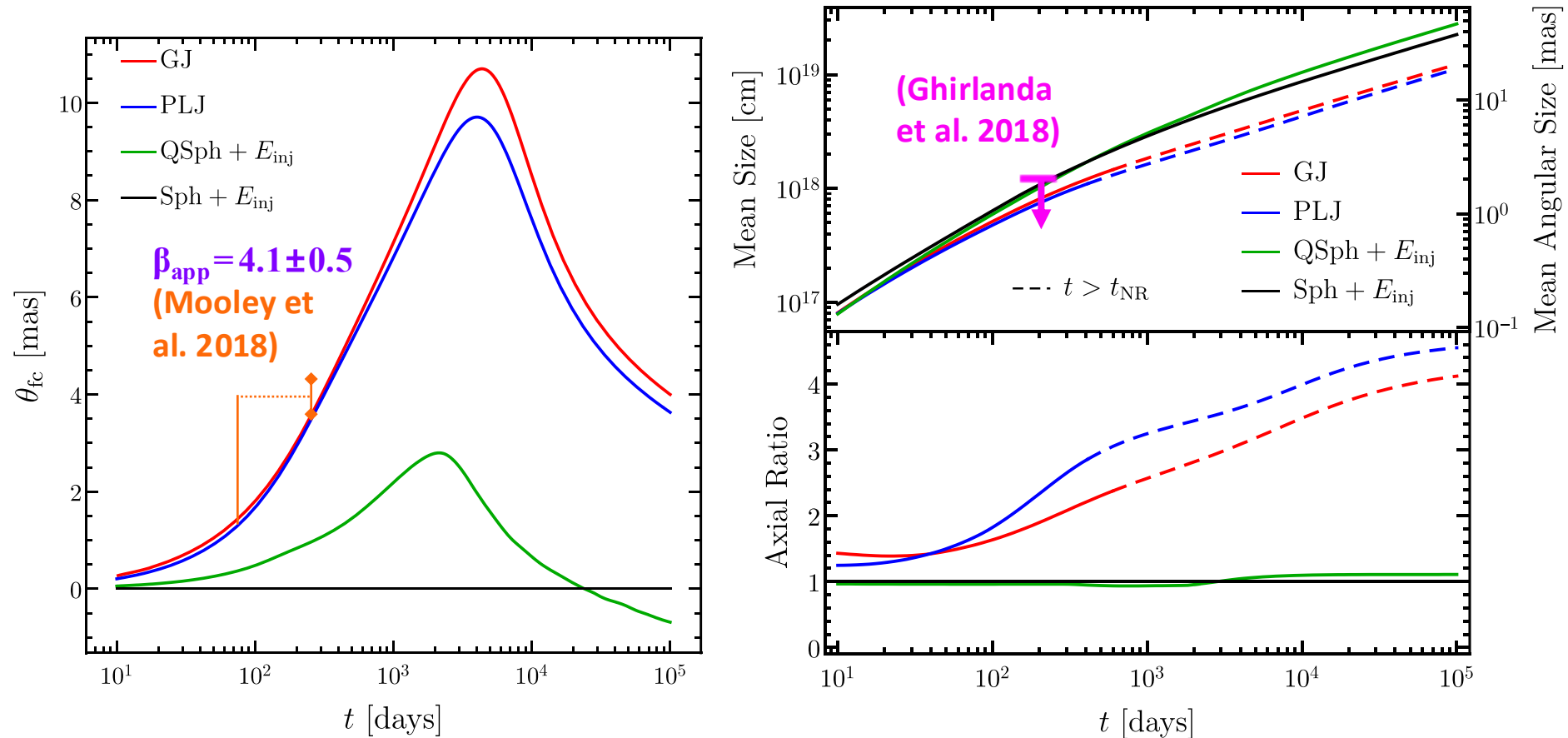
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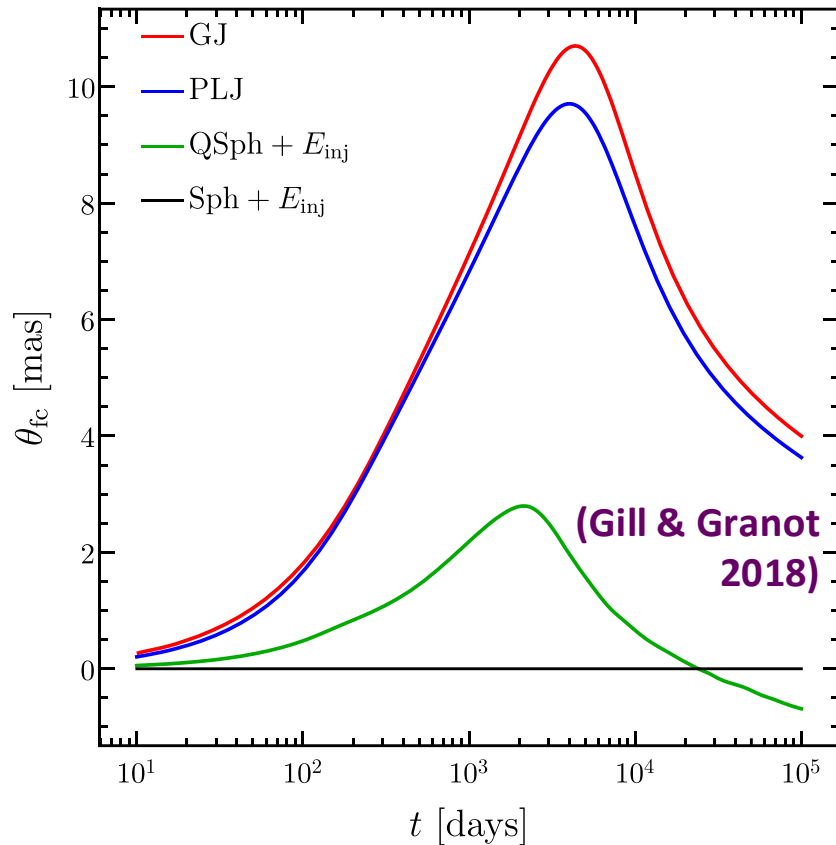
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Afterglow Images: Flux Centroid, Size, Shape

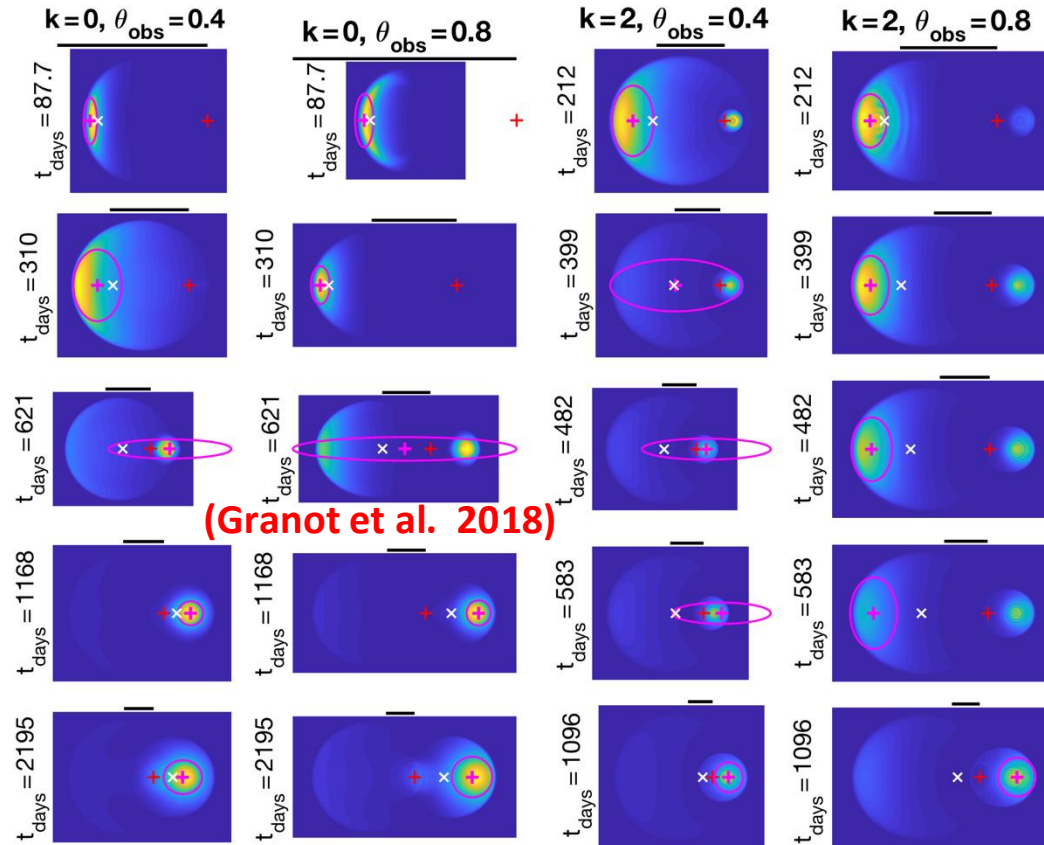
- The flux centroid motion: a potentially powerful diagnostic
- It may be hard to tell apart models based on the image size alone, but a much higher axis-ratio is expected for jet models

Radio flux centroid motion: semi-analytic



(Gill & JG 2018)

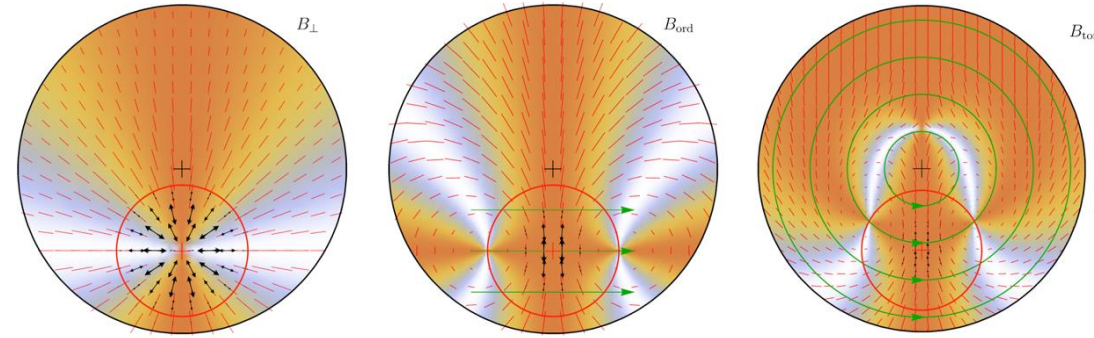
Agree with radio afterglow images from simulations



(JG, De Colle & Ramirez-Ruiz 2018)

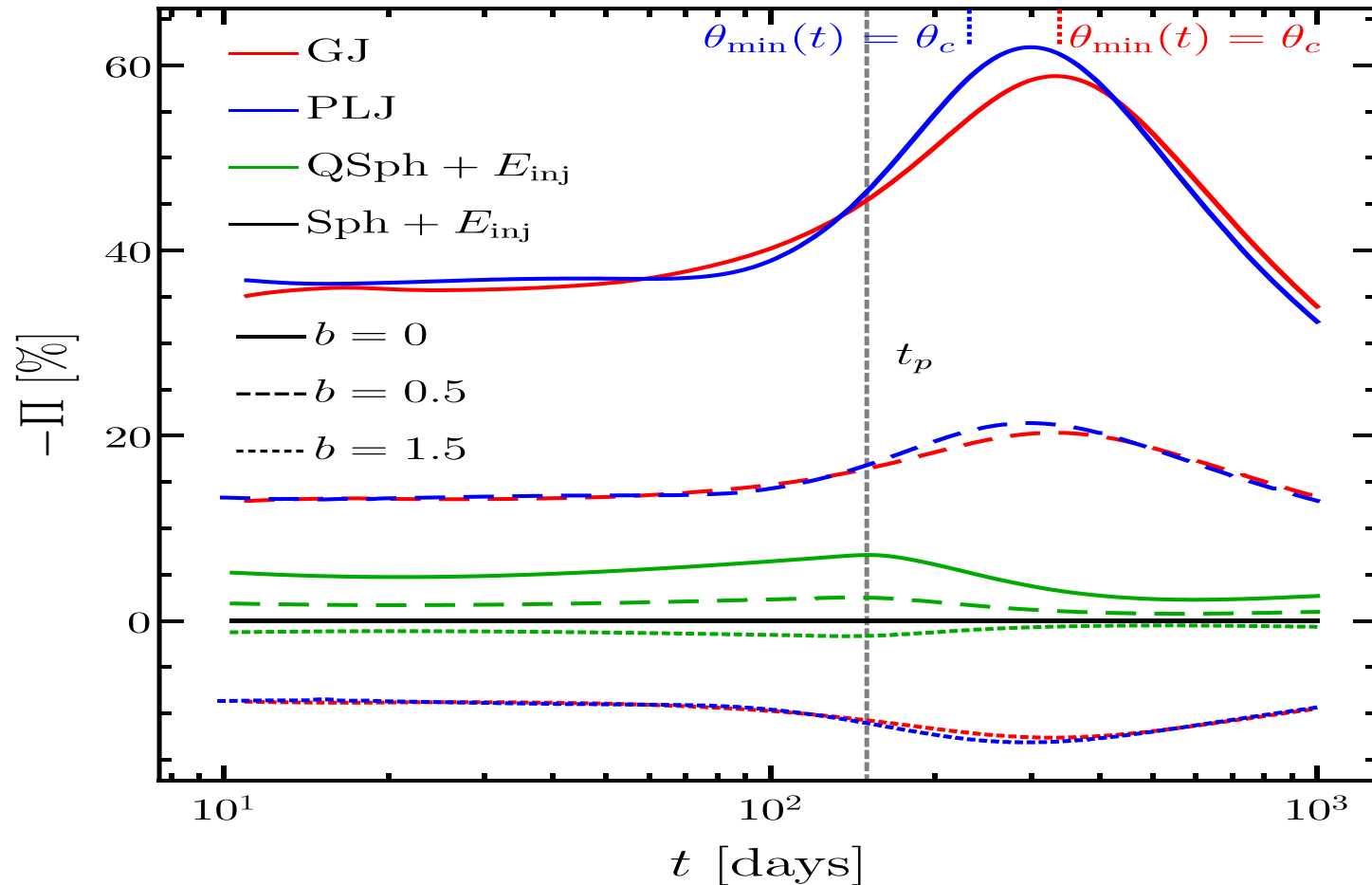
GRB Polarization probes the B-field & Jet structures:

- **Prompt GRB:** hard X-ray – soft γ -ray
 \Rightarrow hard to measure \Rightarrow no clear detections
(stay tuned: POLAR-2, LEAP, COSI, eXTP)
- **Reverse Shock:** also probes original ejecta, but in optical to radio \Rightarrow detections
 - ❖ Probes B-field structure & turbulence in the ejecta near the deceleration epoch
 - ❖ Radio UL $\Pi_{8.5\text{GHz}}(1 \text{ day}) < 7\% \Rightarrow$ rules out ordered B-field (coherent over $\theta_B \gtrsim 1/\Gamma$) or toroidal B-field + $dE/d\Omega \propto \theta^{-2}$ structured jet; allows $\theta_B \lesssim 10^{-2}$ patches (JG & Taylor 2005)
 - ❖ RINGO2 GRB120308A: $\Pi_{\text{opt}}(240 - 323 \text{ s}) = 28 \pm 4\% \Rightarrow$ ordered B-field (Mundell et al. 2013)
 - ❖ ALMA GRB190114C: $\Pi_{97.5\text{GHz}}(2.2 \rightarrow 5.2 \text{ hr}) = 0.9 \rightarrow 0.6\%$ with $\Delta\theta_p(2.2 \rightarrow 5.2 \text{ hr}) = 54^\circ$ (first GRB radio polarization) \Rightarrow favors patches with $\theta_B \sim 10^{-3}$ (Laskar et al. 2019)
- **Afterglow:** optical & radio – probes jet angular structure & B-field structure in collisionless relativistic shocks



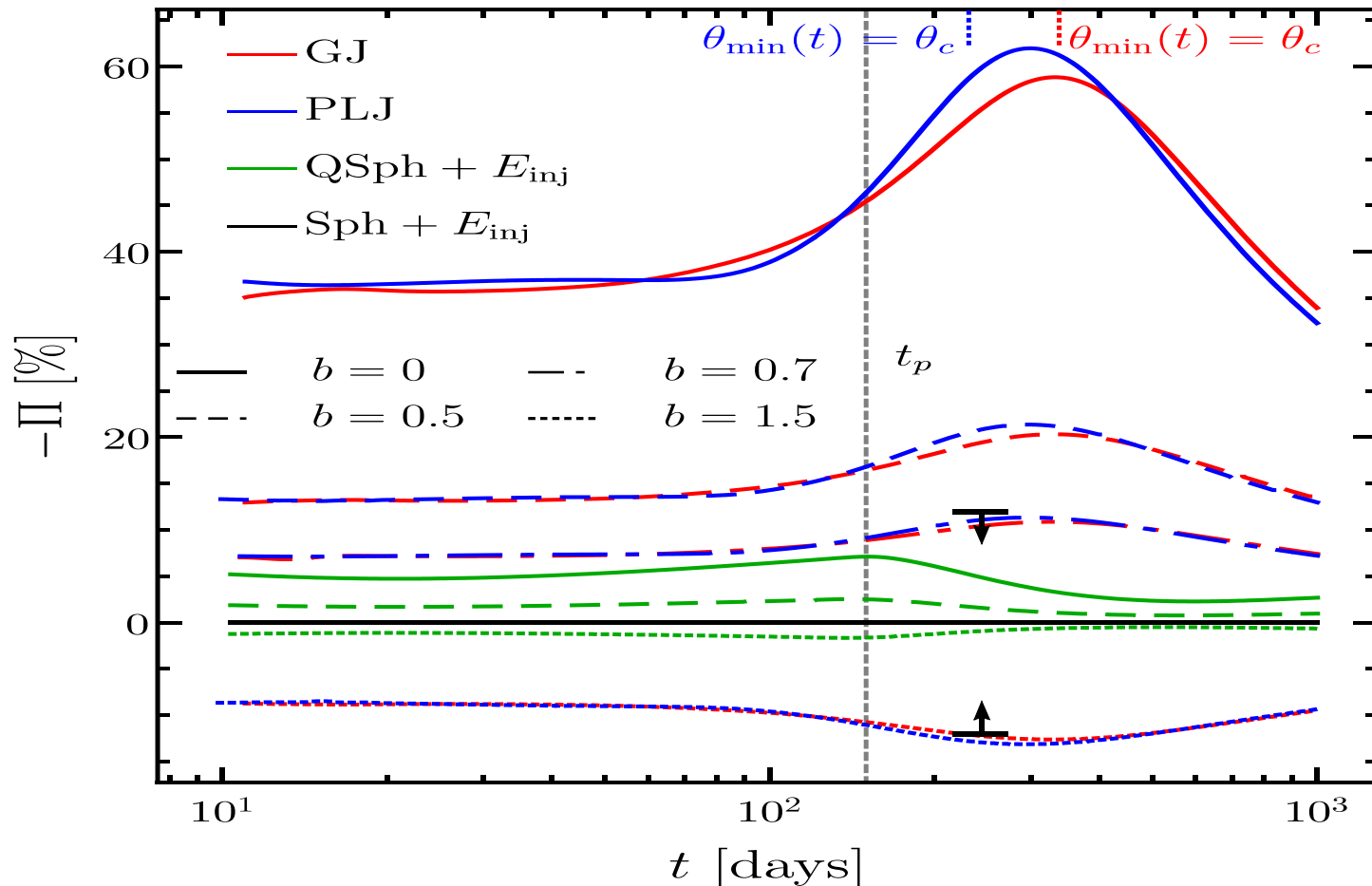
GRB 170817A: polarization UL \Rightarrow post-shock B-field

- Jet angular structure & θ_{obs} well constrained \Rightarrow breaks degeneracies
- Assuming a shock-produce B-field with $b \equiv 2\langle B_{\parallel}^2 \rangle / \langle B_{\perp}^2 \rangle$ (JG & königl 03; Gill & JG 18)



GRB 170817A: polarization UL \Rightarrow post-shock B-field

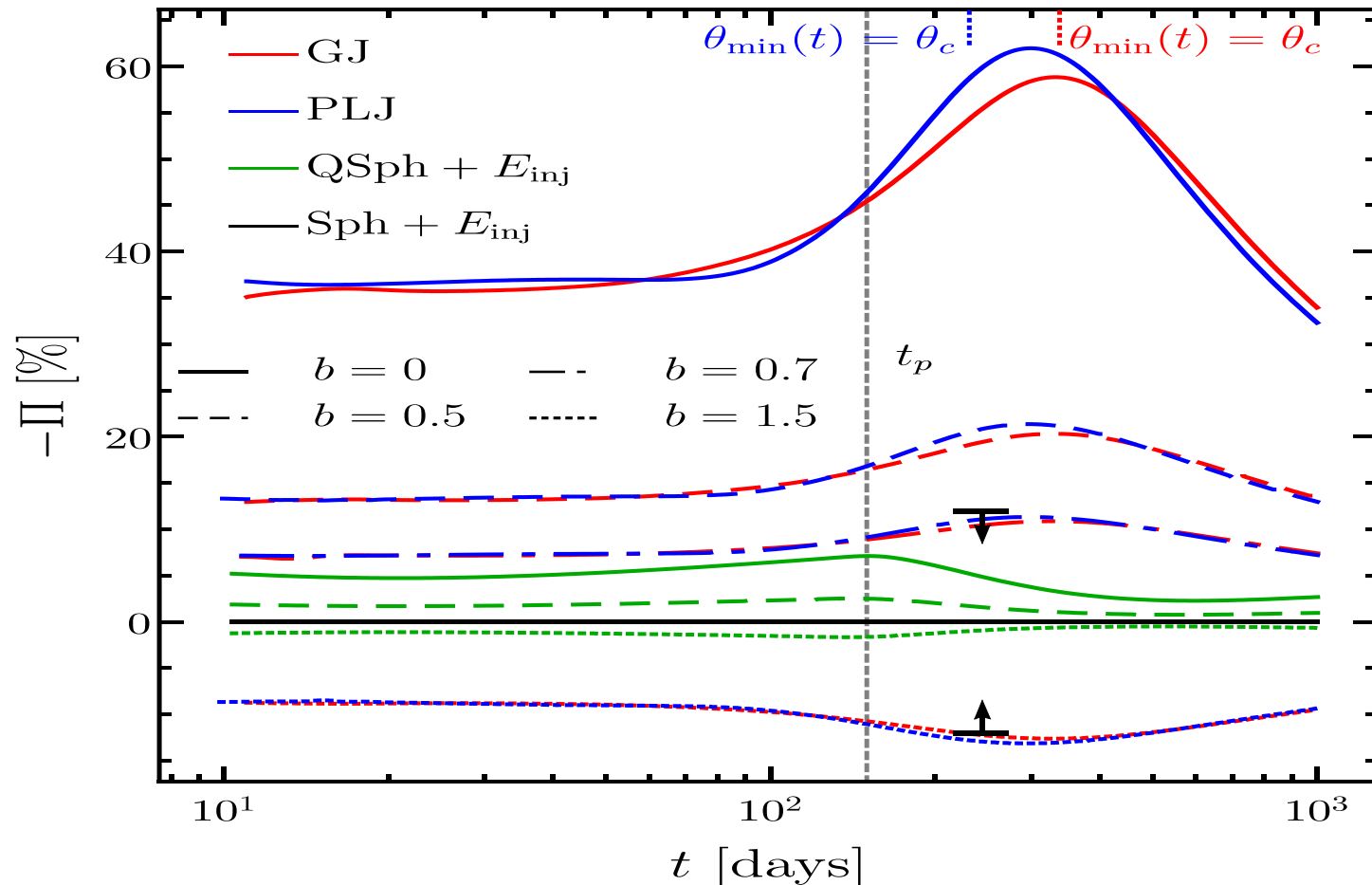
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Later: upper limit
 $P_{\text{lin}} < 12\%$ @
 $\nu = 2.8 \text{ GHz}$,
 $t = 244 \text{ days}$
(Corsi+ 2018)

GRB 170817A: polarization UL \Rightarrow post-shock B-field

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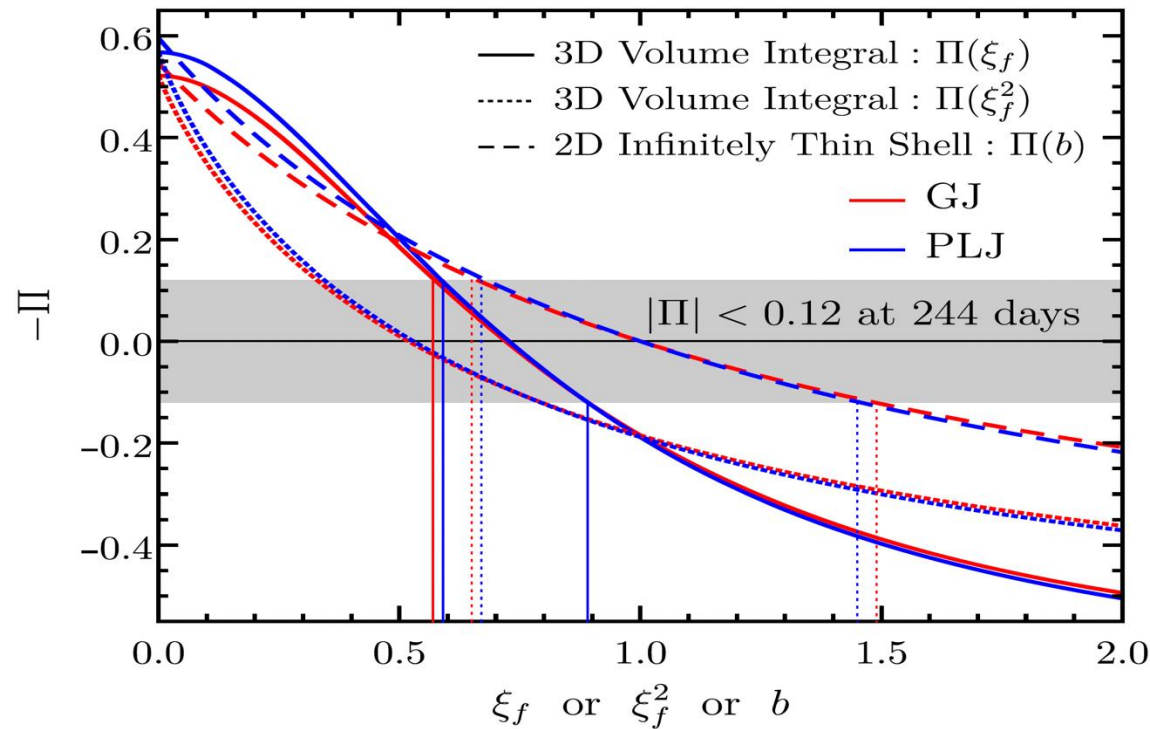
$0.7 \lesssim b \lesssim 1.5$
for jet models

Later: upper limit
 $P_{\text{lin}} < 12\%$ @
 $\nu = 2.8\text{GHz}$,
 $t = 244\text{ days}$
(Corsi+ 2018)

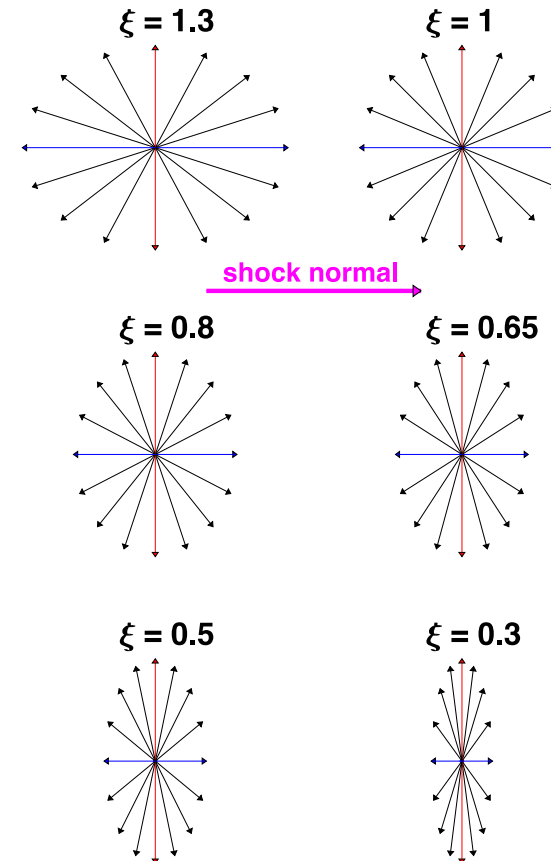
GRB 170817A: polarization UL \Rightarrow post-shock B-field

More realistic assumptions \Rightarrow B-field in collisionless shocks: (Gill & JG 2020)

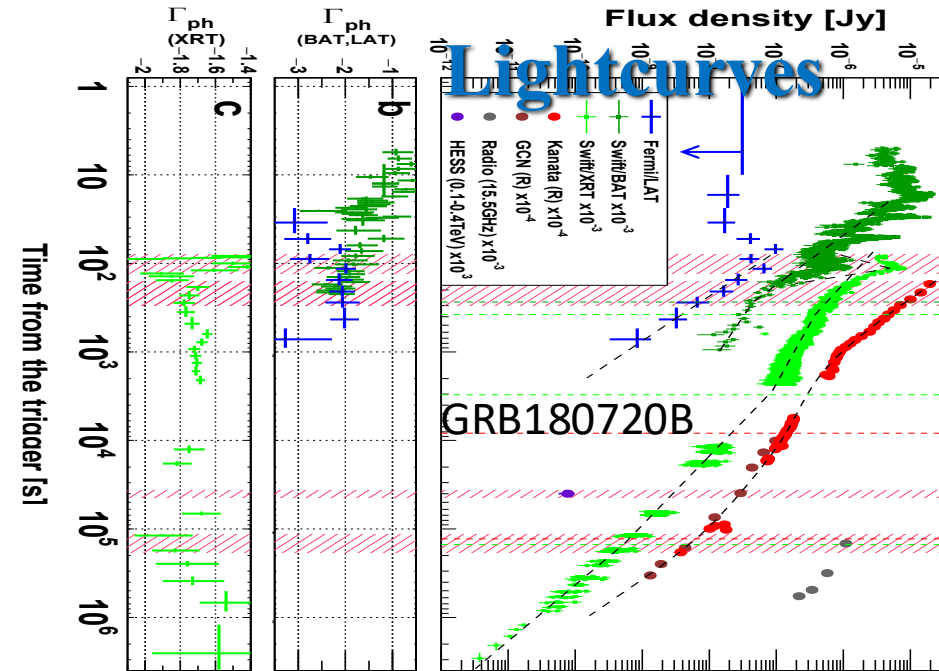
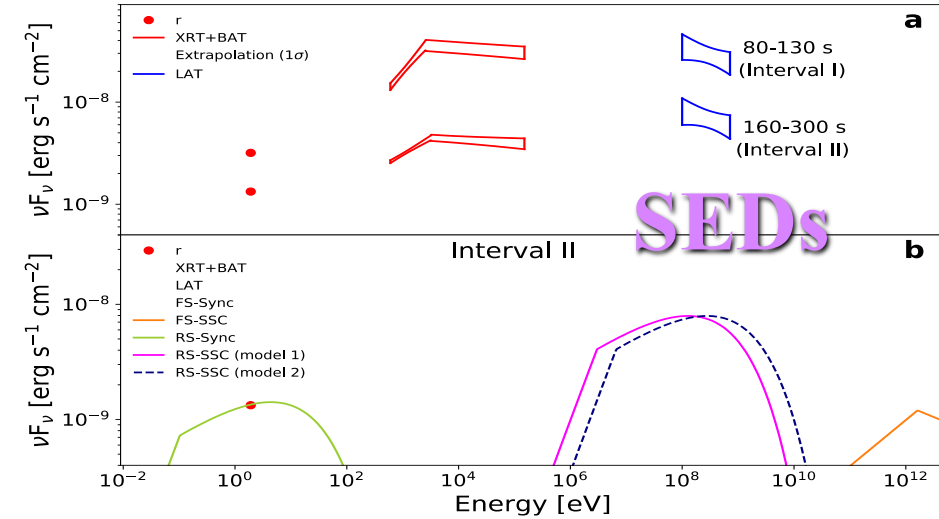
- 2D emitting shell \rightarrow 3D emitting volume (local BM76 radial profile)
- B-field evolution by faster radial expansion: $L'_r / L'_{\theta,\phi} \propto \chi^{(7-2k)/(8-2k)}$
- B-field isotropic in 3D with $B'_r \rightarrow \xi B'_r$ (Sari 1999); $\xi = \xi_f \chi^{(7-2k)/(8-2k)}$



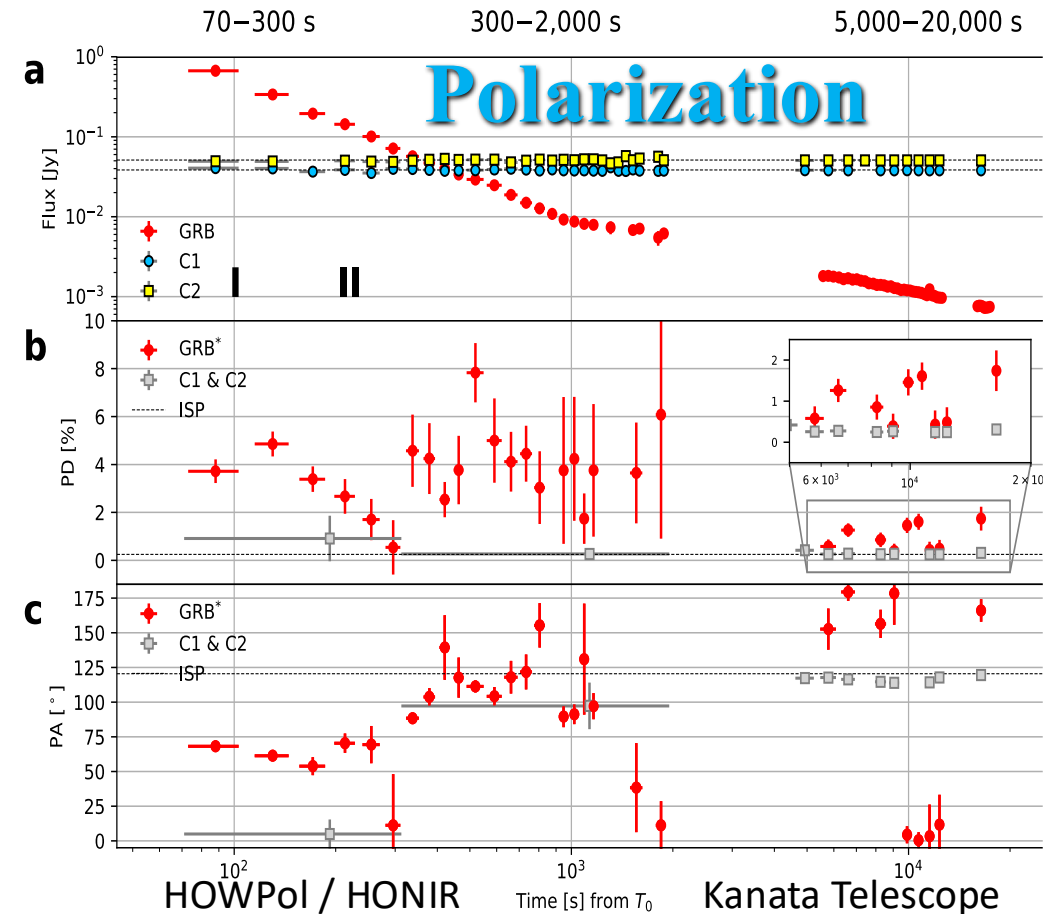
$$0.57 \lesssim \xi_f \lesssim 0.89$$



Reverse + Forward Shock Polarization: (Arimoto et al. 2023)

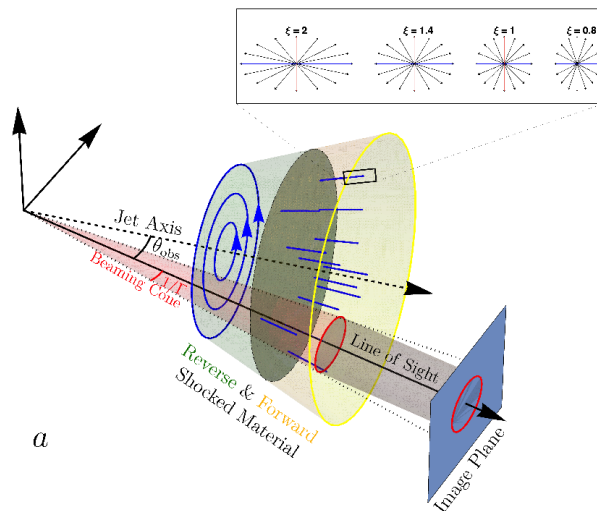
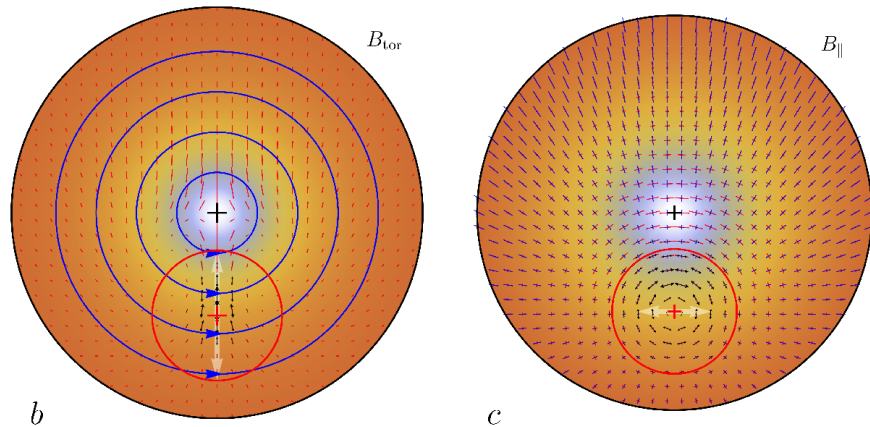


- **< 300 s: RS; P: 5 → 1%, $\theta_p \approx 70^\circ$**
- **0.3-2 ks: P ~ 2-8%, θ_p varies**
- **5-20 ks: FS; P ~ 0.5-2%, $\theta_p \approx 160^\circ$**
- **RS → FS dominance @ $\sim 10^3$ s**

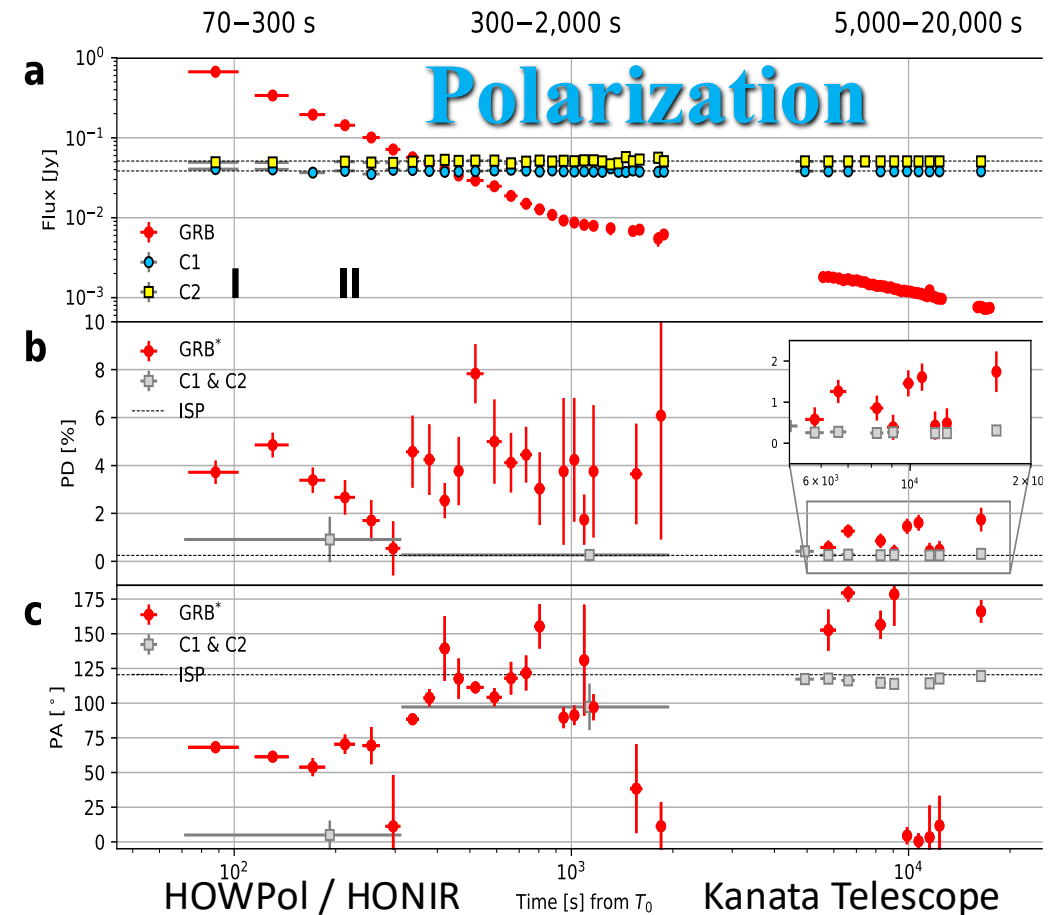


Reverse + Forward Shock Polarization: (Arimoto et al. 2023)

- < 300 s: ejecta; B_{tor} + turbulence
- 0.3-2 ks: turbulence-induced P
- 5-20 ks: CSM; radial stretching

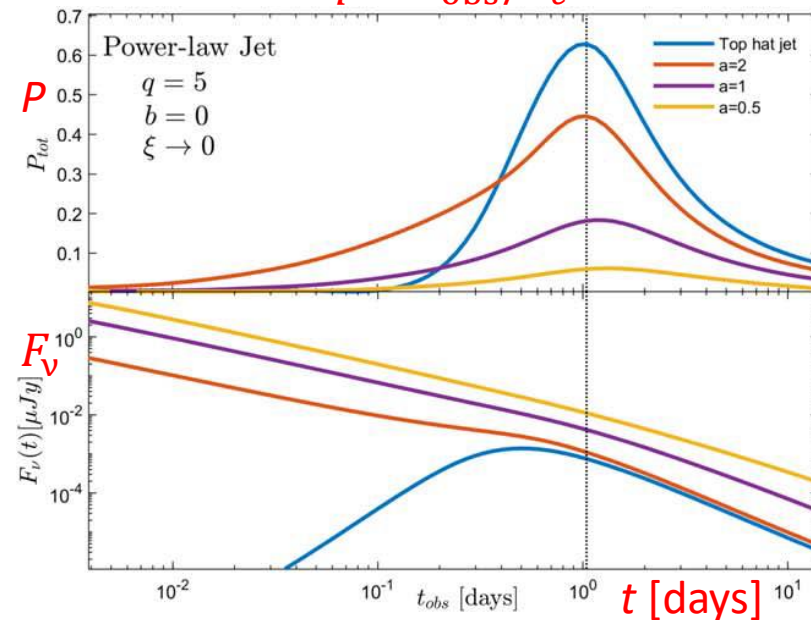
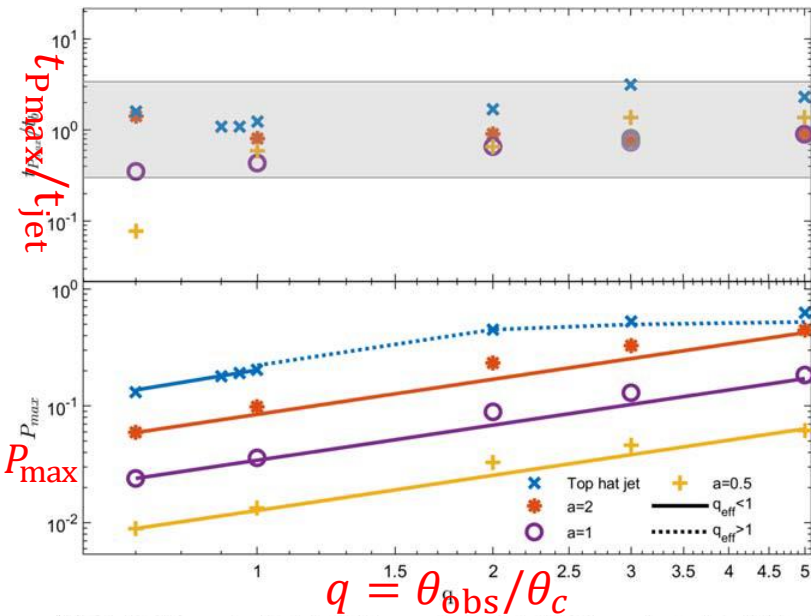


- < 300 s: RS; P: 5 → 1%, $\theta_p \approx 70^\circ$
- 0.3-2 ks: P ~ 2-8%, θ_p varies
- 5-20 ks: FS; P ~ 0.5-2%, $\theta_p \approx 160^\circ$
- RS → FS dominance @ $\sim 10^3$ s

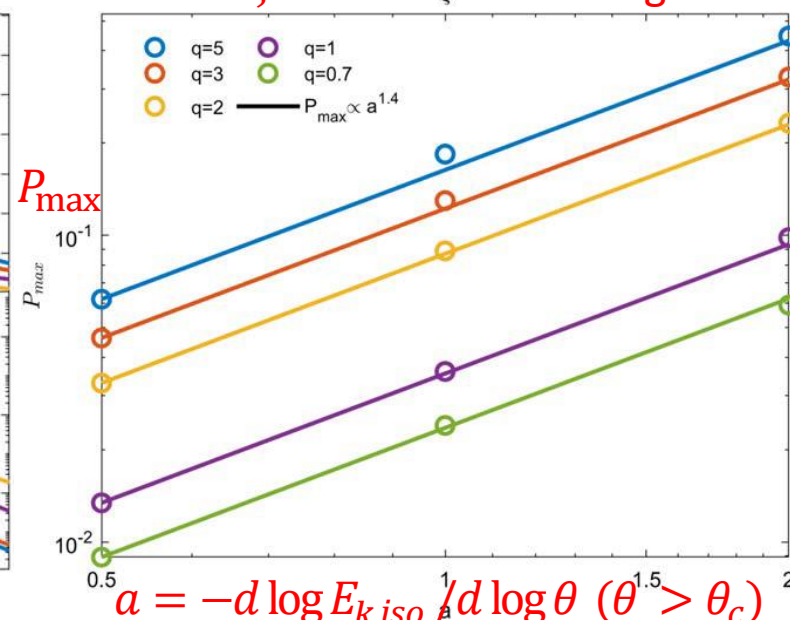
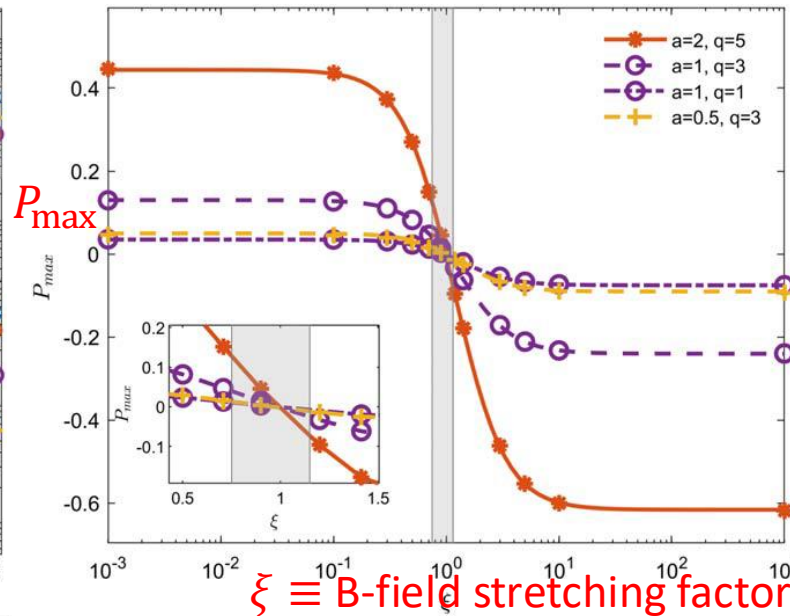


Afterglow Polarization from Shallow Jets (Birenbaum et al. 2024)

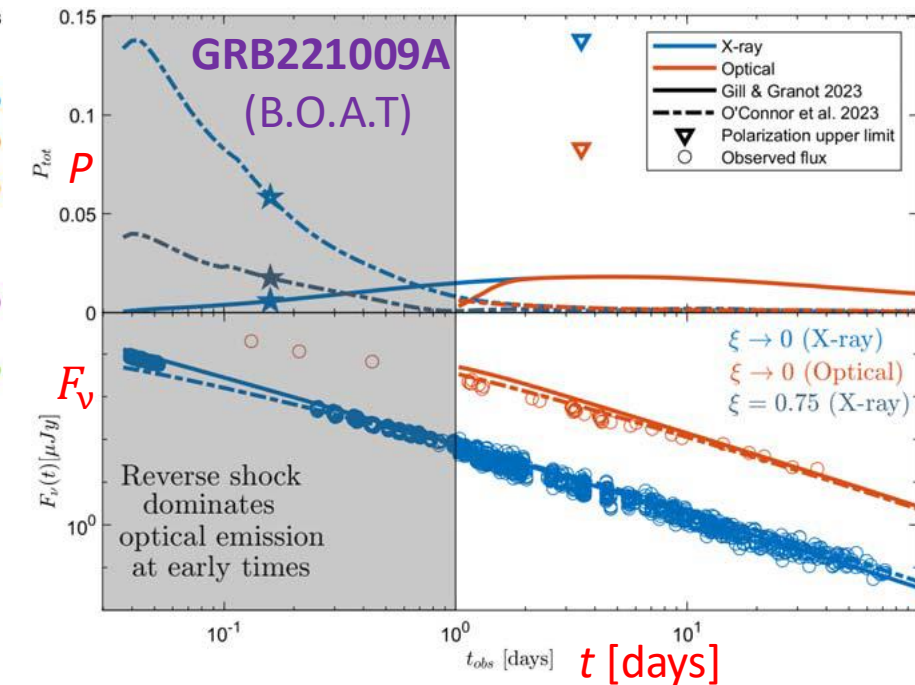
P peaks near the jet break time



$$P_{\max} \approx qa^{1.4}[0.055 \tanh(0.34 - 2.3 \log_{10} \xi) - 0.02]$$

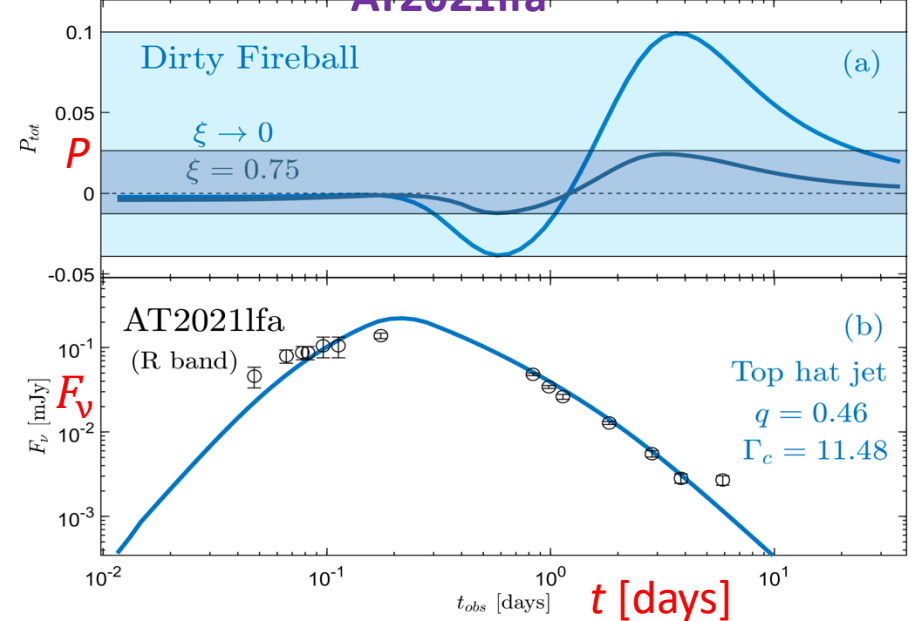
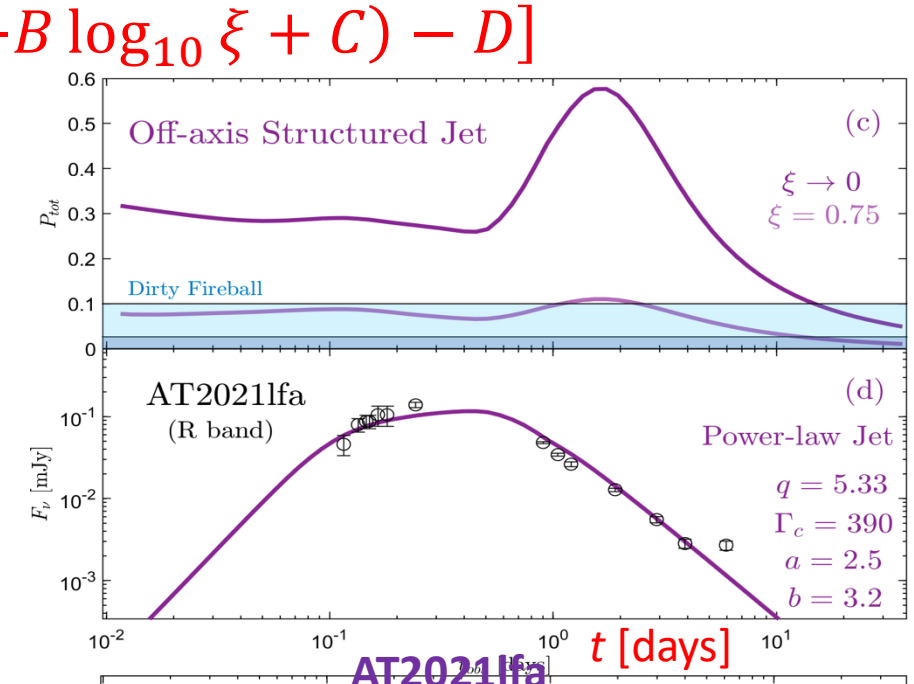
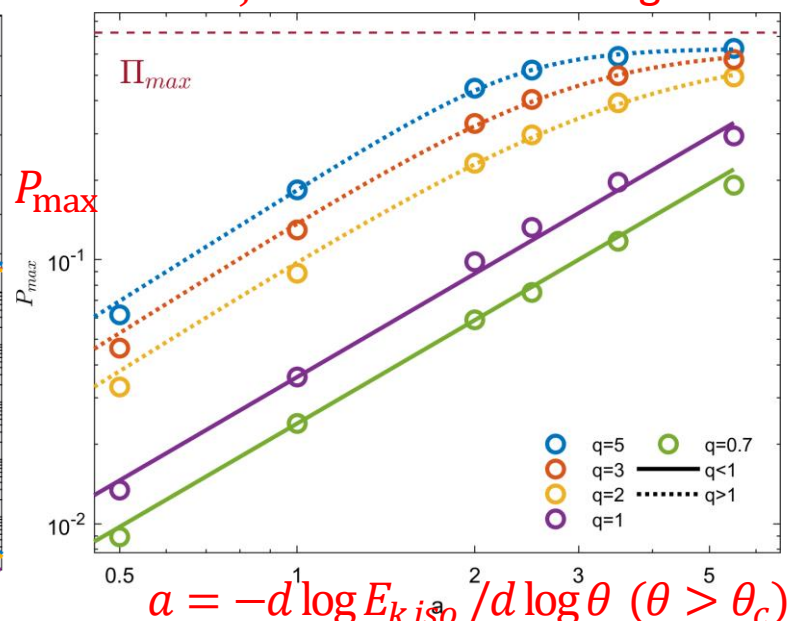
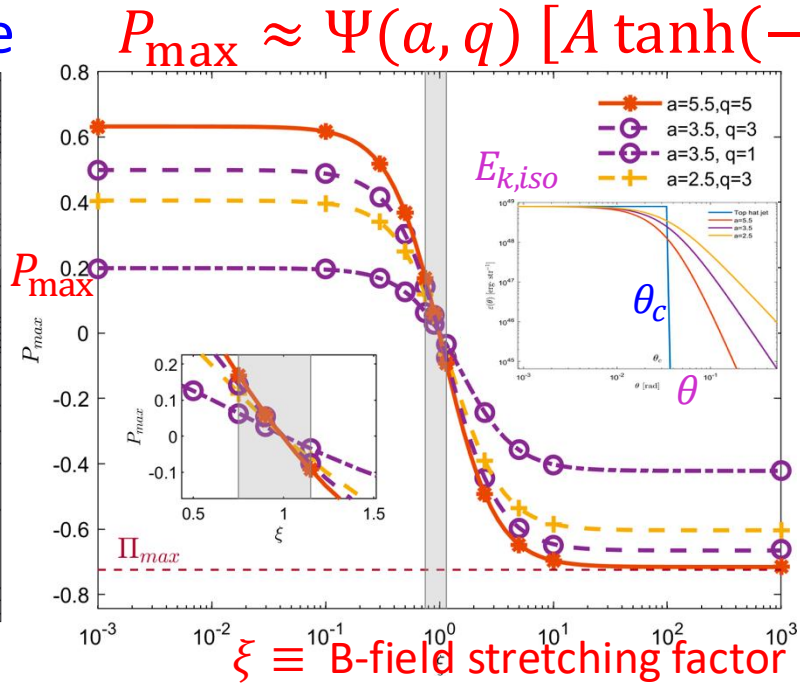
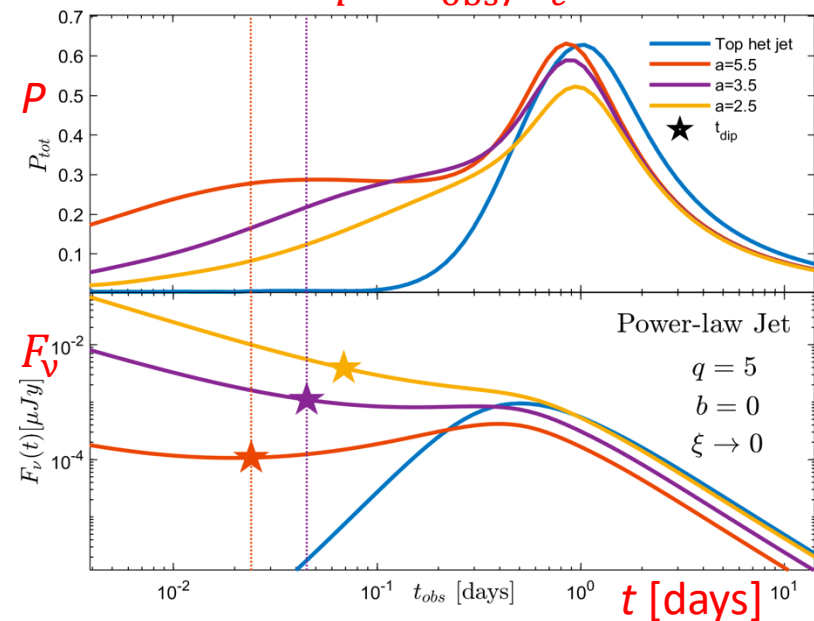
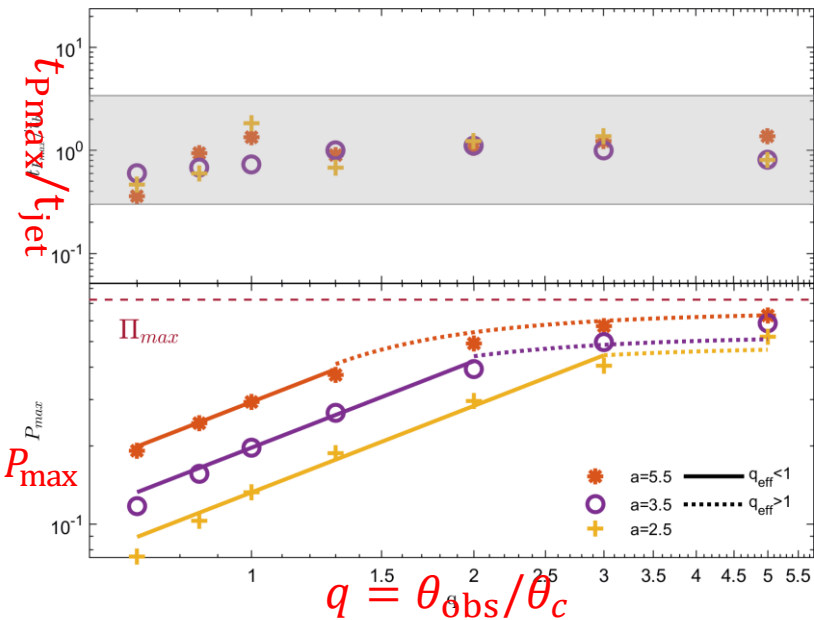


- Particularly energetic GRBs seem to have **shallow Jets**
- Earlier polarimetry of the B.O.A.T could have helped constrain its jet structure & post-shock B-field structure



Afterglow Polarization from Steep Jets ($a \geq 2$; Birenbaum et al. 2026)

P peaks near the jet break time



ULGRB250702B: Prompt Emission

$T_{\text{GRB}} \approx 25 \text{ ks}$ (longest ever)

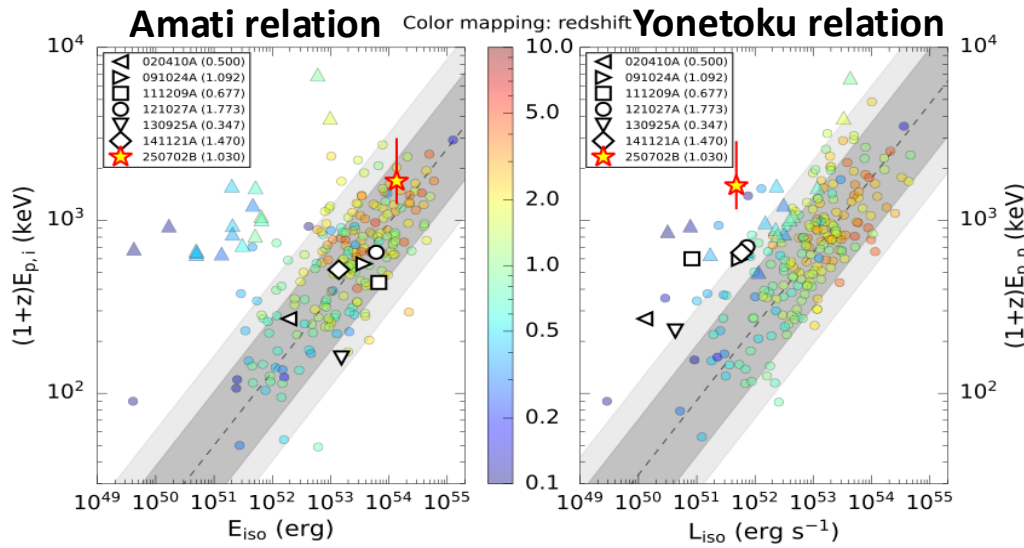
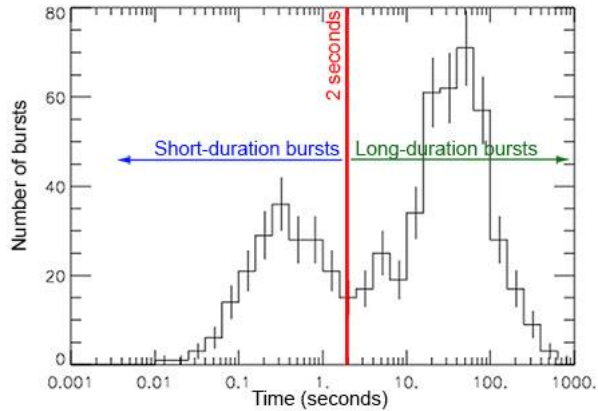
$E_{\gamma,\text{iso}} \geq 1.4 \times 10^{54} \text{ erg}$

$L_{\gamma,\text{iso}} \approx 5 \times 10^{51} \text{ erg/s}$

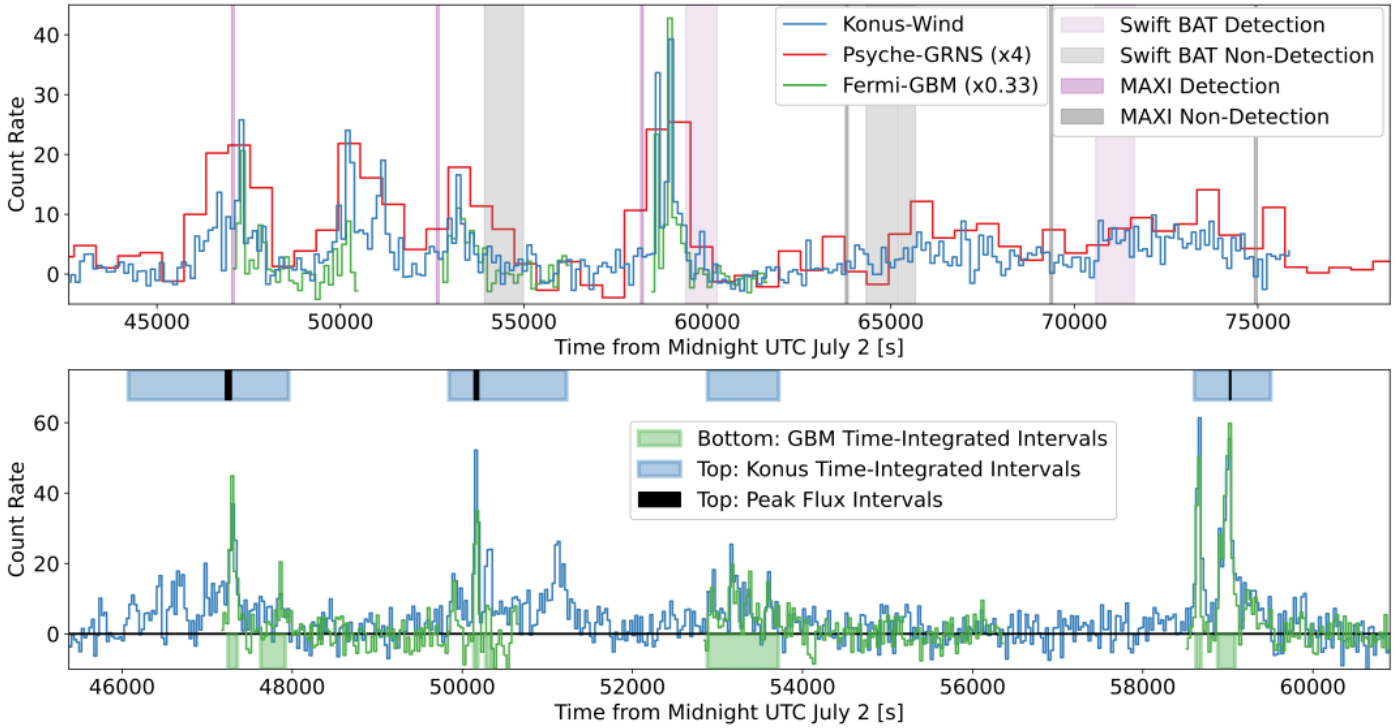
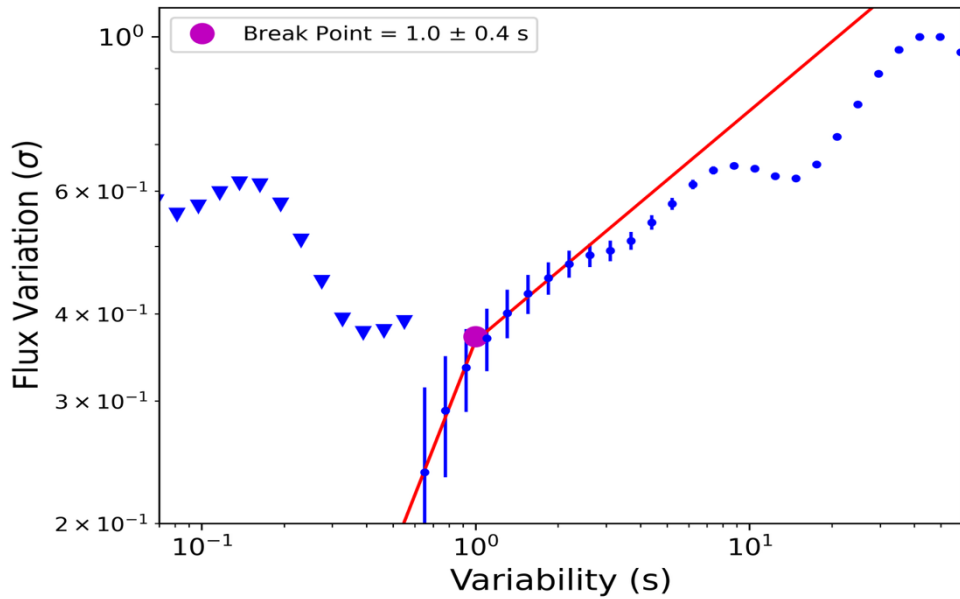
$E_{p,z} \sim 1.5 - 6 \text{ MeV}$

$t_{\text{MV}} \approx 1\text{s}, E_{\text{max}} \approx 5 \text{ MeV} \Rightarrow \Gamma > 56$

$E_{\text{cut}} \approx 4 \text{ MeV} \Rightarrow \Gamma \approx 81$



(Neights et al. 2025)



Ultra-Long GRBs:

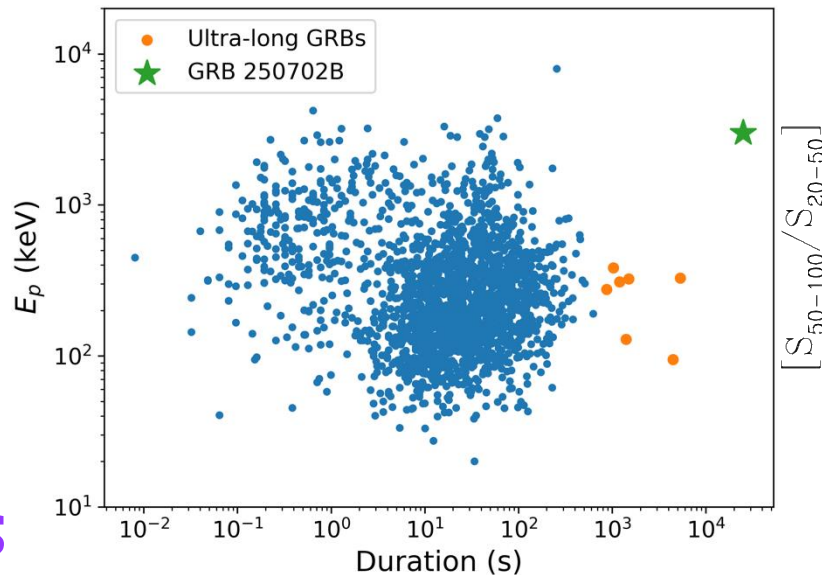
$$T_{\text{GRB}} \approx 25 \text{ ks}$$

$$E_{\gamma,\text{iso}} \geq 1.4 \times 10^{54} \text{ erg}$$

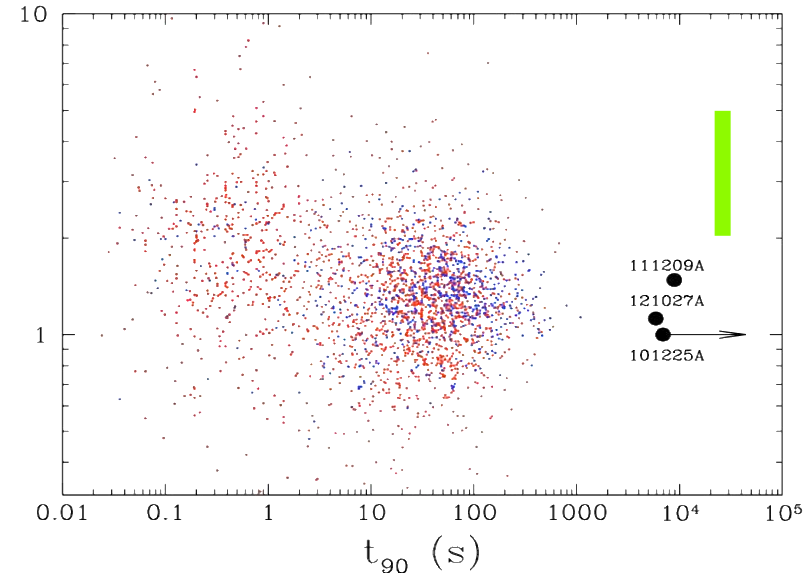
$$L_{\gamma,\text{iso}} \approx 5 \times 10^{51} \text{ erg/s}$$

$$\bar{L}_{\gamma,\text{iso}} \geq 1.1 \times 10^{50} \text{ erg/s}$$

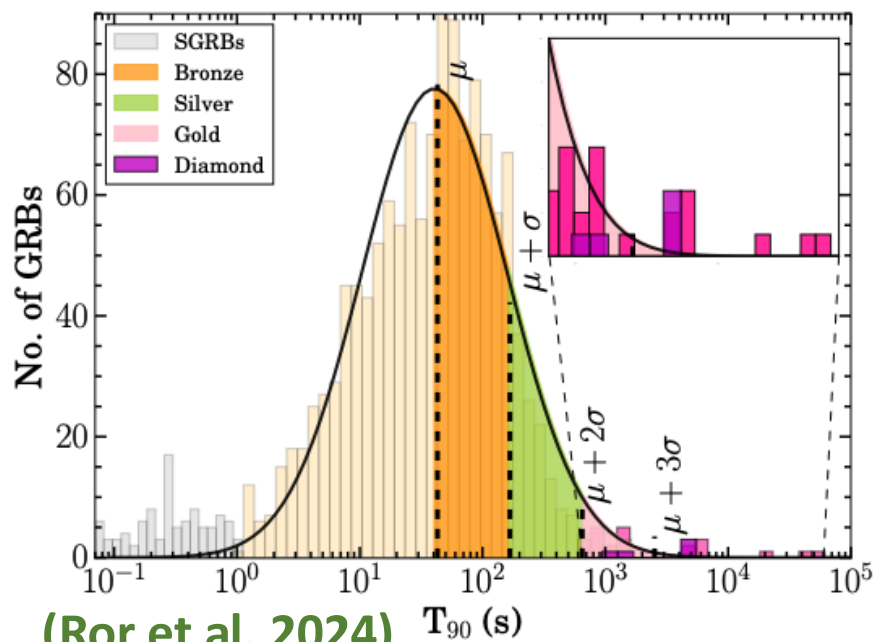
$$E_{p,z} \sim 1.5 - 6 \text{ MeV}$$



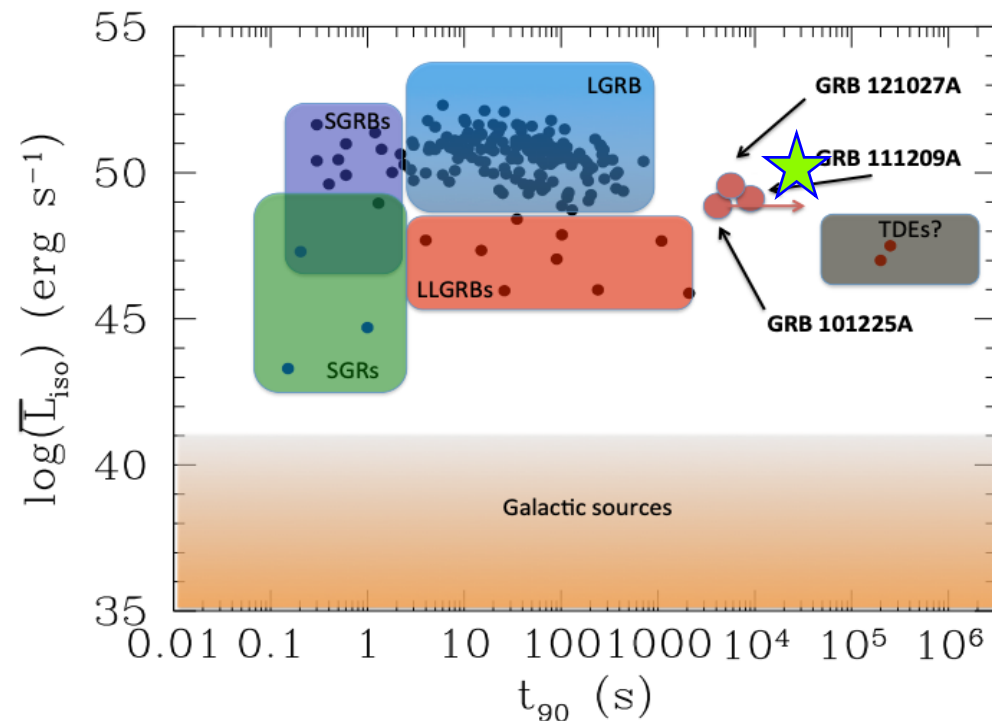
(Neights et al. 2025)



(Levan et al. 2014)



(Ror et al. 2024)

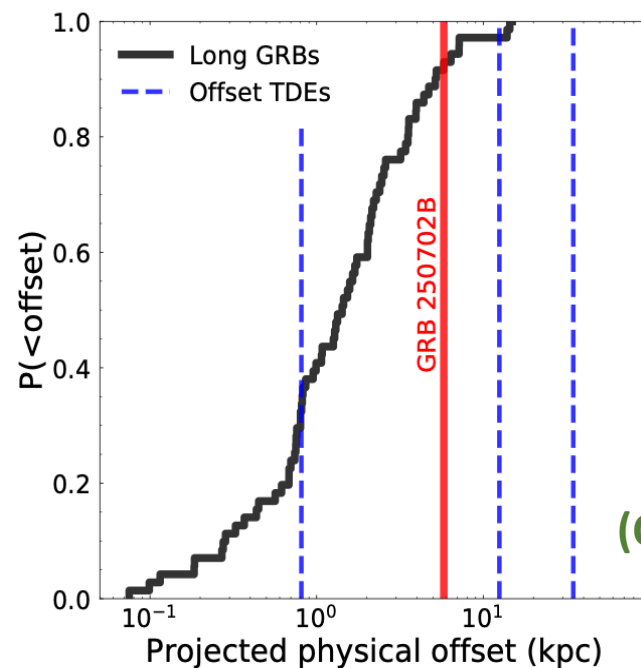


ULGRB250702B

Redshift &
Host galaxy:

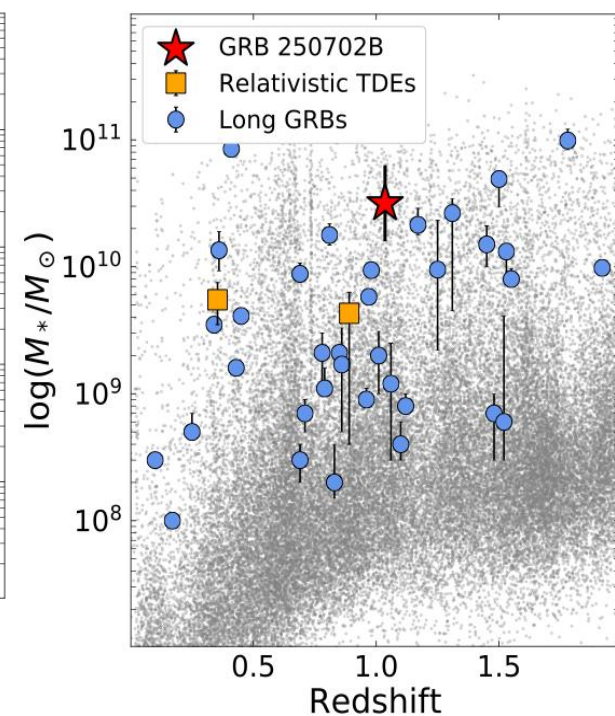
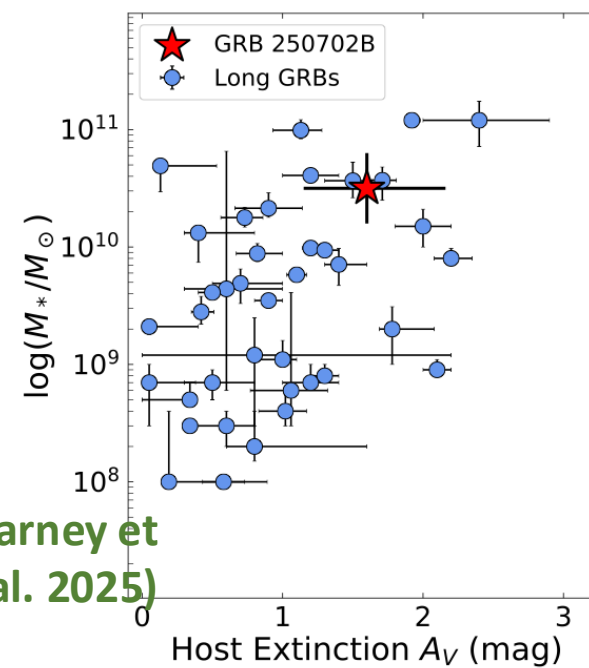
JWST: $z = 1.036$

Massive spiral
galaxy (with dust
lane?) or merger

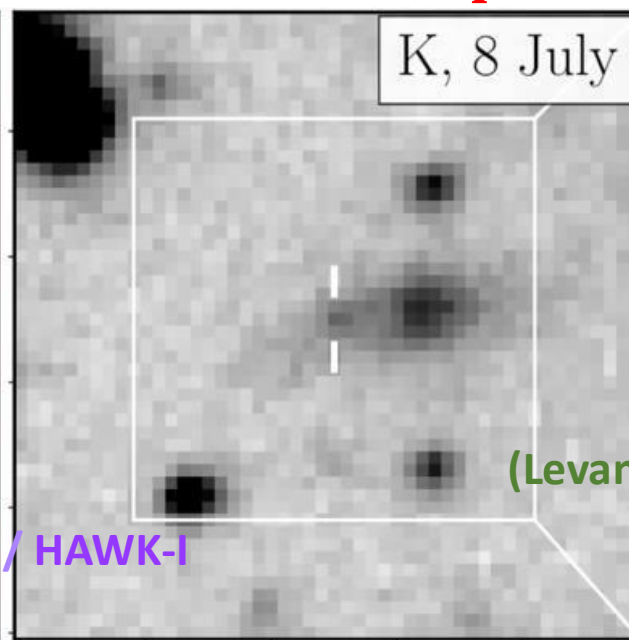
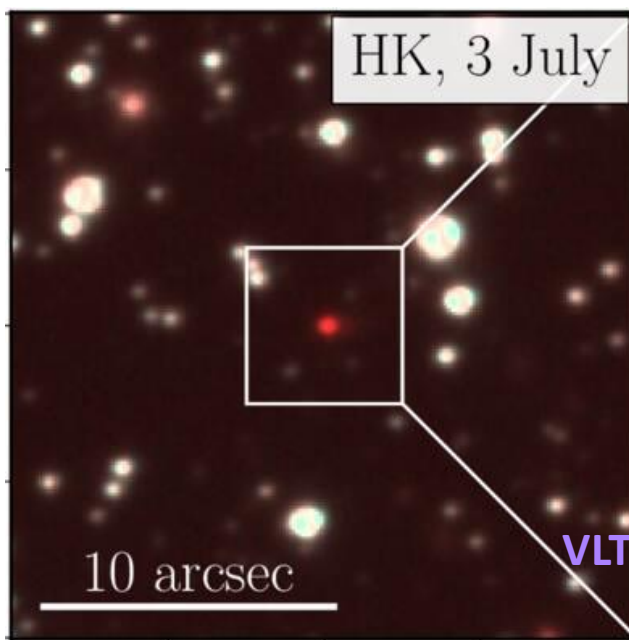


(Carney et al. 2025)

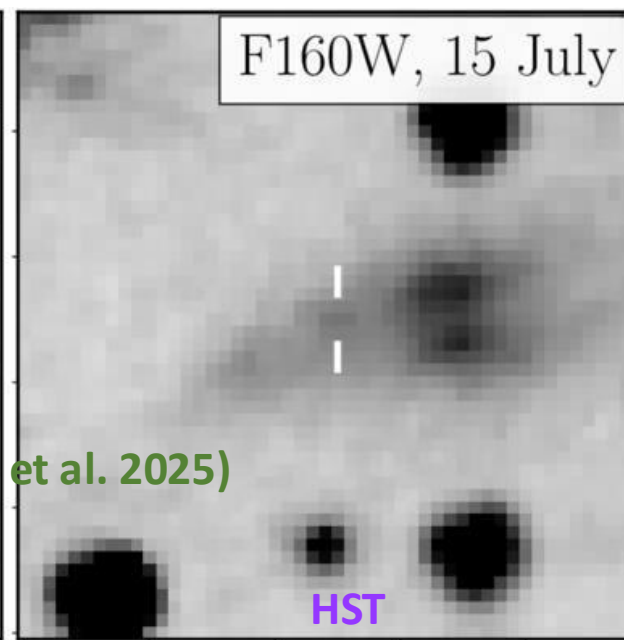
offset = 5.7 kpc



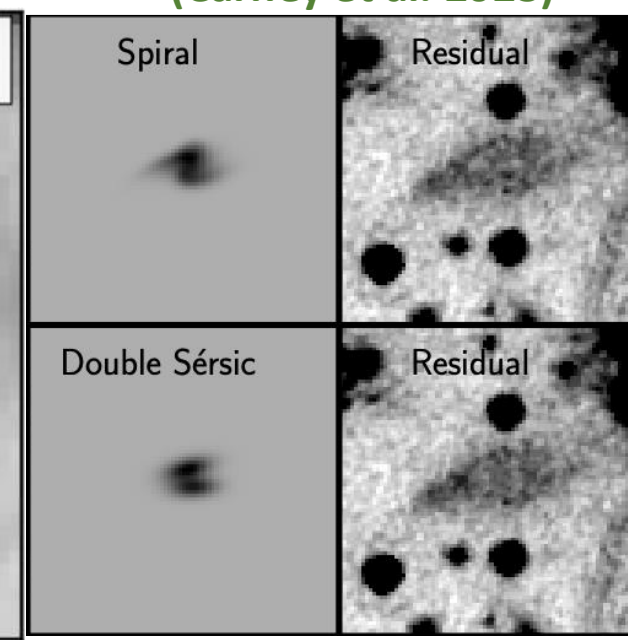
(Carney et al. 2025)



(Levan et al. 2025)

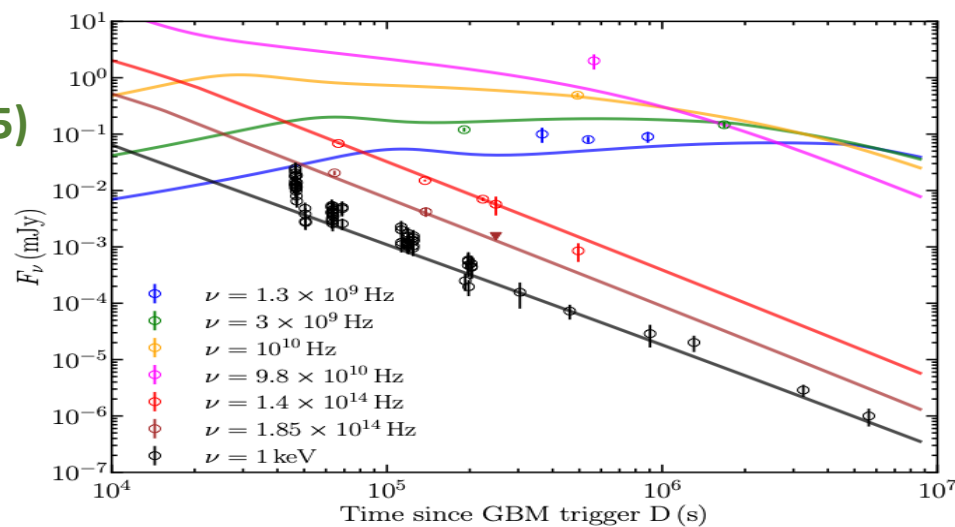
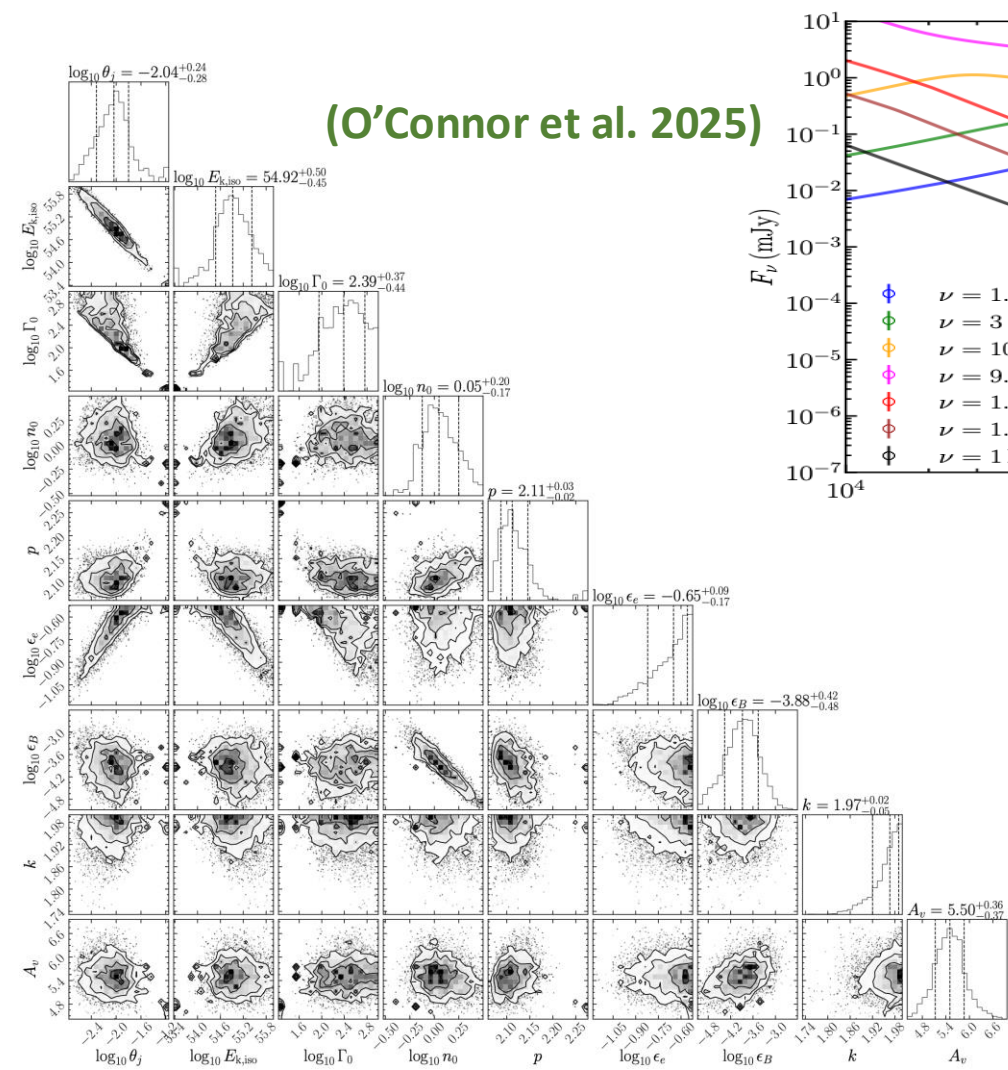


HST

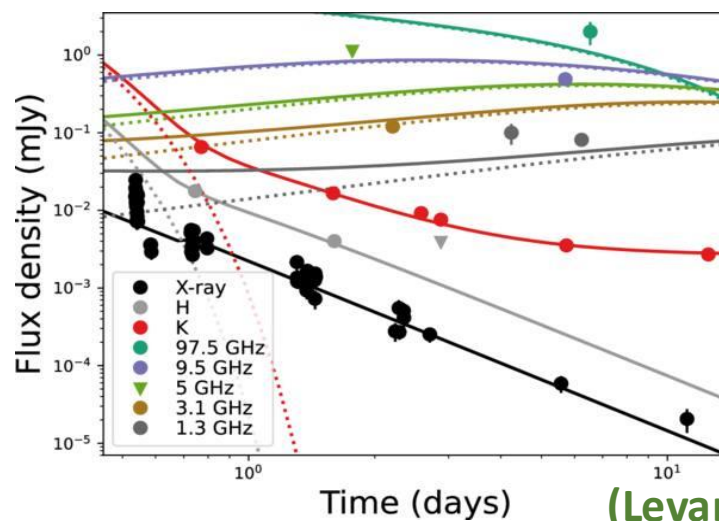


ULGRB250702B Afterglow: X-ray, IR, radio

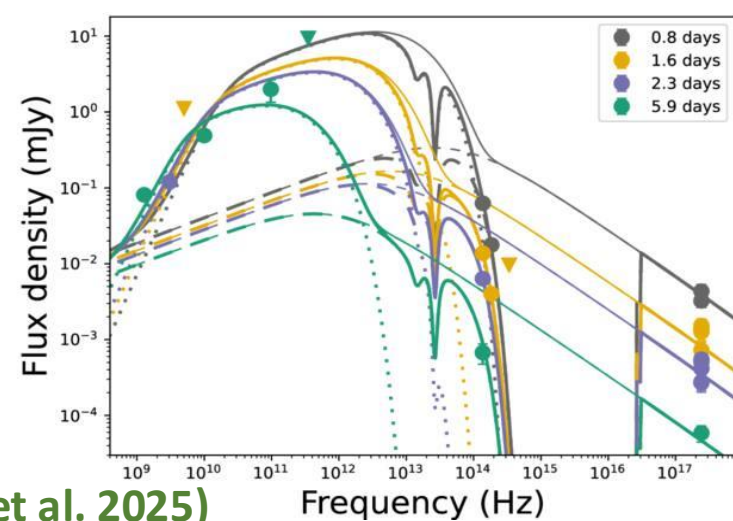
- **Extinction:** $A_V = 0.847$ mag (MW) + 2 – 9 mag (host); $N_H \approx (3 - 5) \times 10^{22} \text{ cm}^{-2}$



- **Stratified external medium**
- $E_{k,iso} \sim \text{few} \times E_{\gamma,iso}$
- **Narrow jet:** $\theta_{jet} \sim 10^{-2}$ rad
- **Still many degeneracies; more data expected**



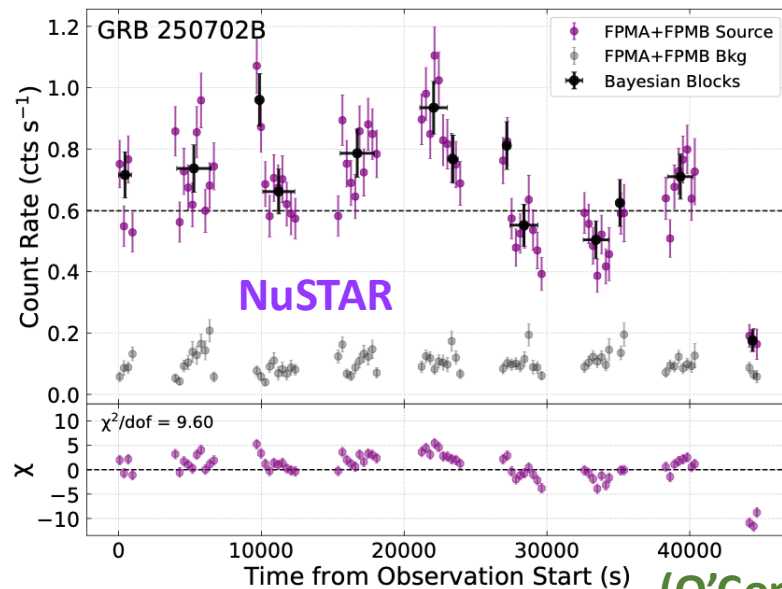
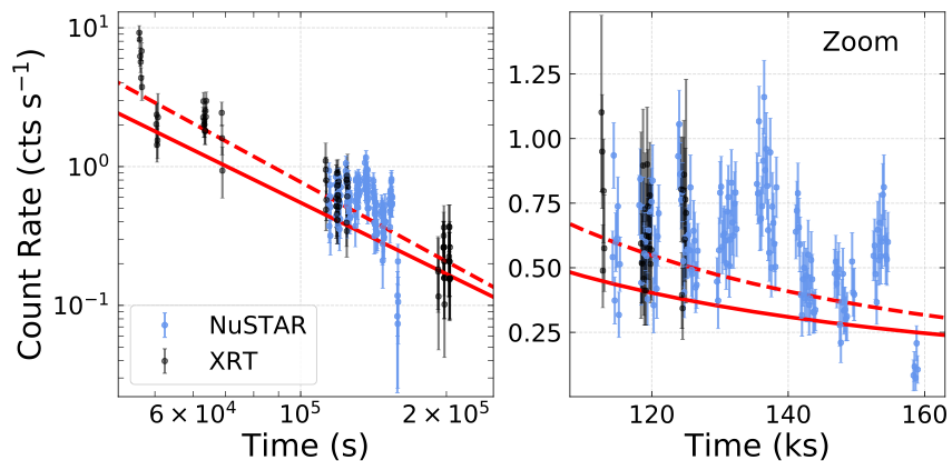
(Levan et al. 2025)



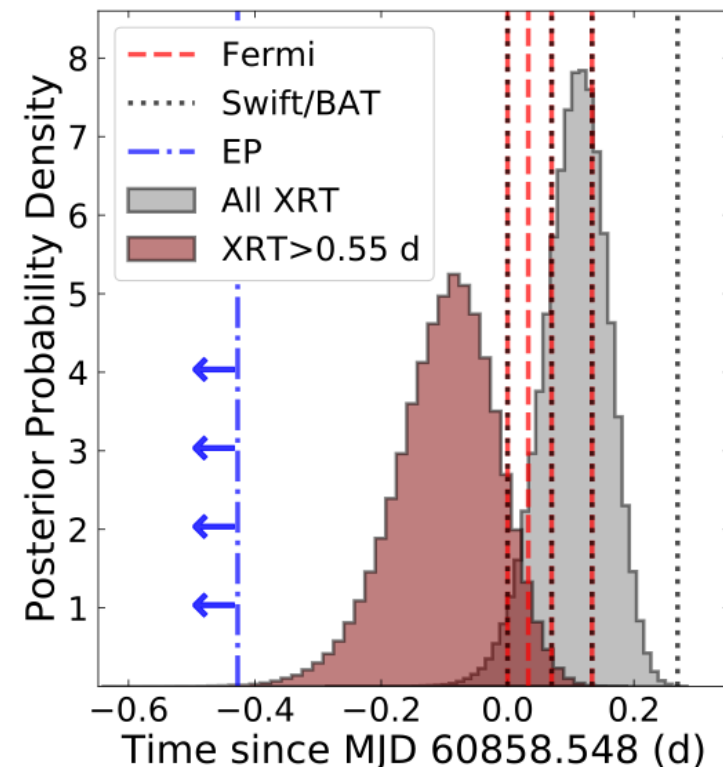
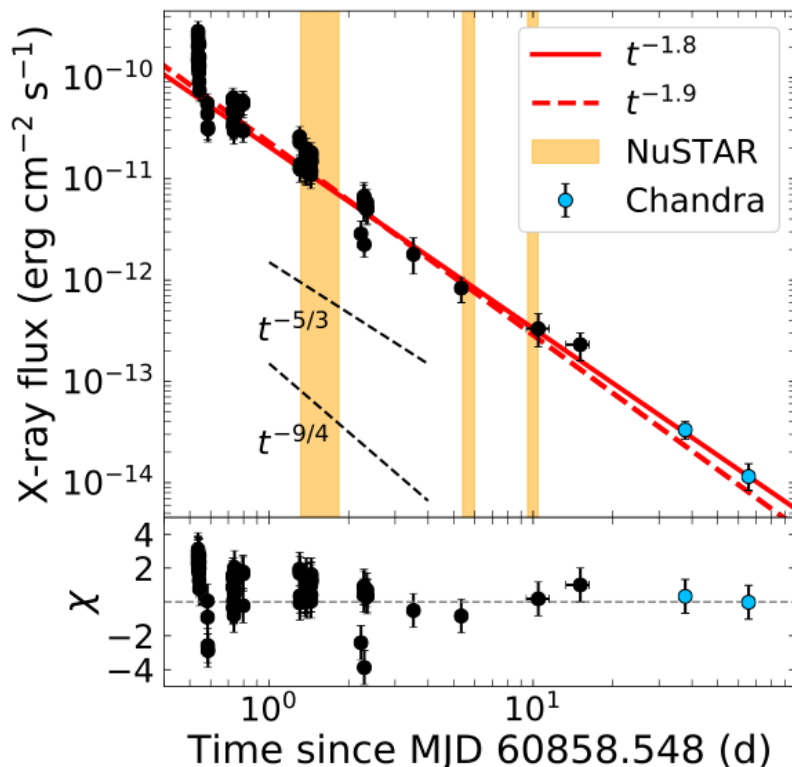
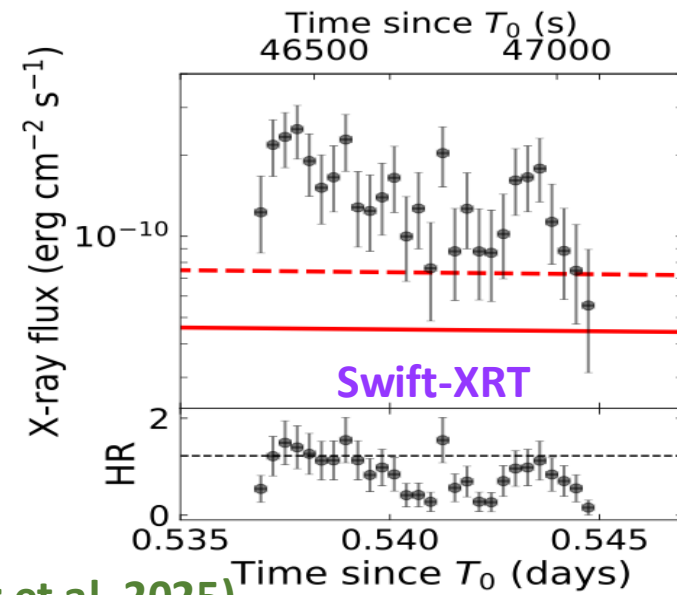
ULGRB250702B

X-ray Afterglow:

- Very rapid variability in first few days: $\Delta t/t \sim 10^{-2.5}$
⇒ rules out afterglow origin
- Likely late time residual source activity (accretion powered jet)
- Decay slope broadly consistent with TDE expectations

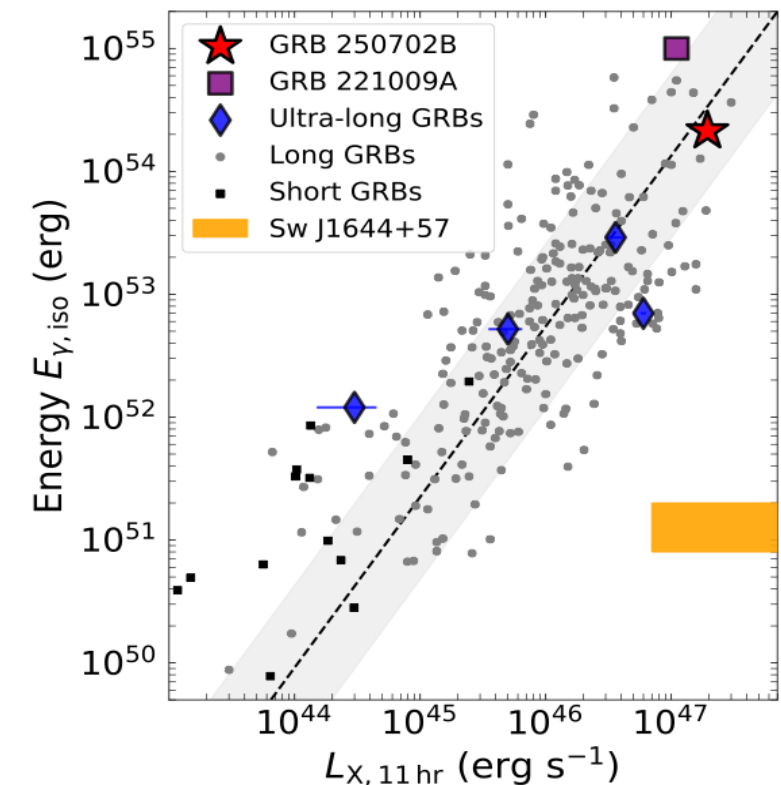
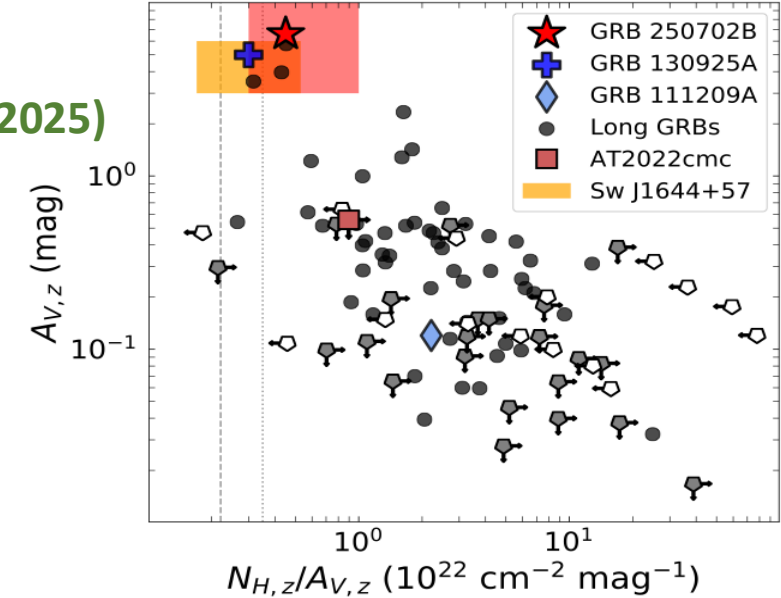
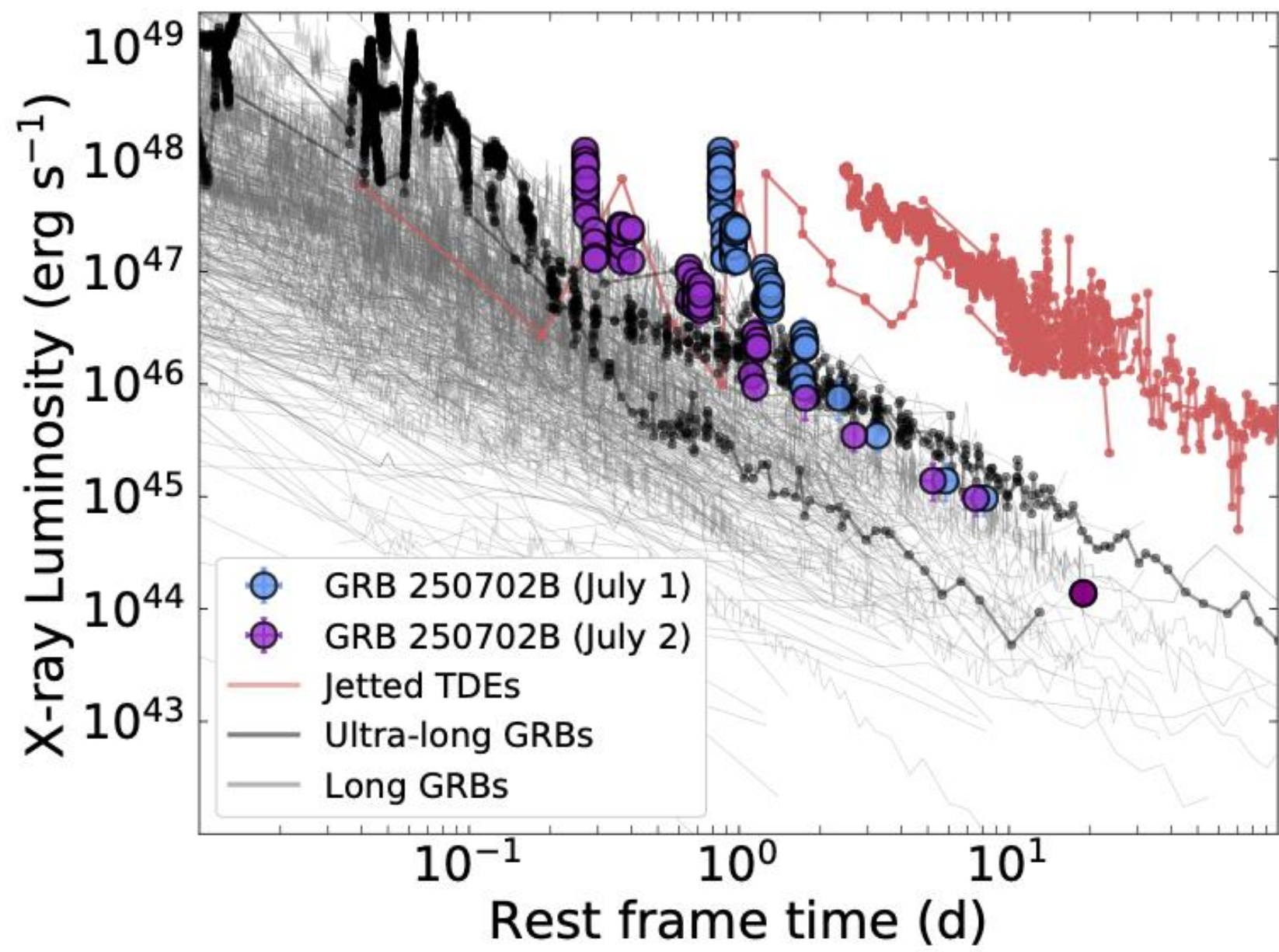


(O'Connor et al. 2025)



ULGRB250702B: GRB or TDE?

(O'Connor et al. 2025)



ULGRB250702B: Einstein Probe: X-ray starts 1 day earlier

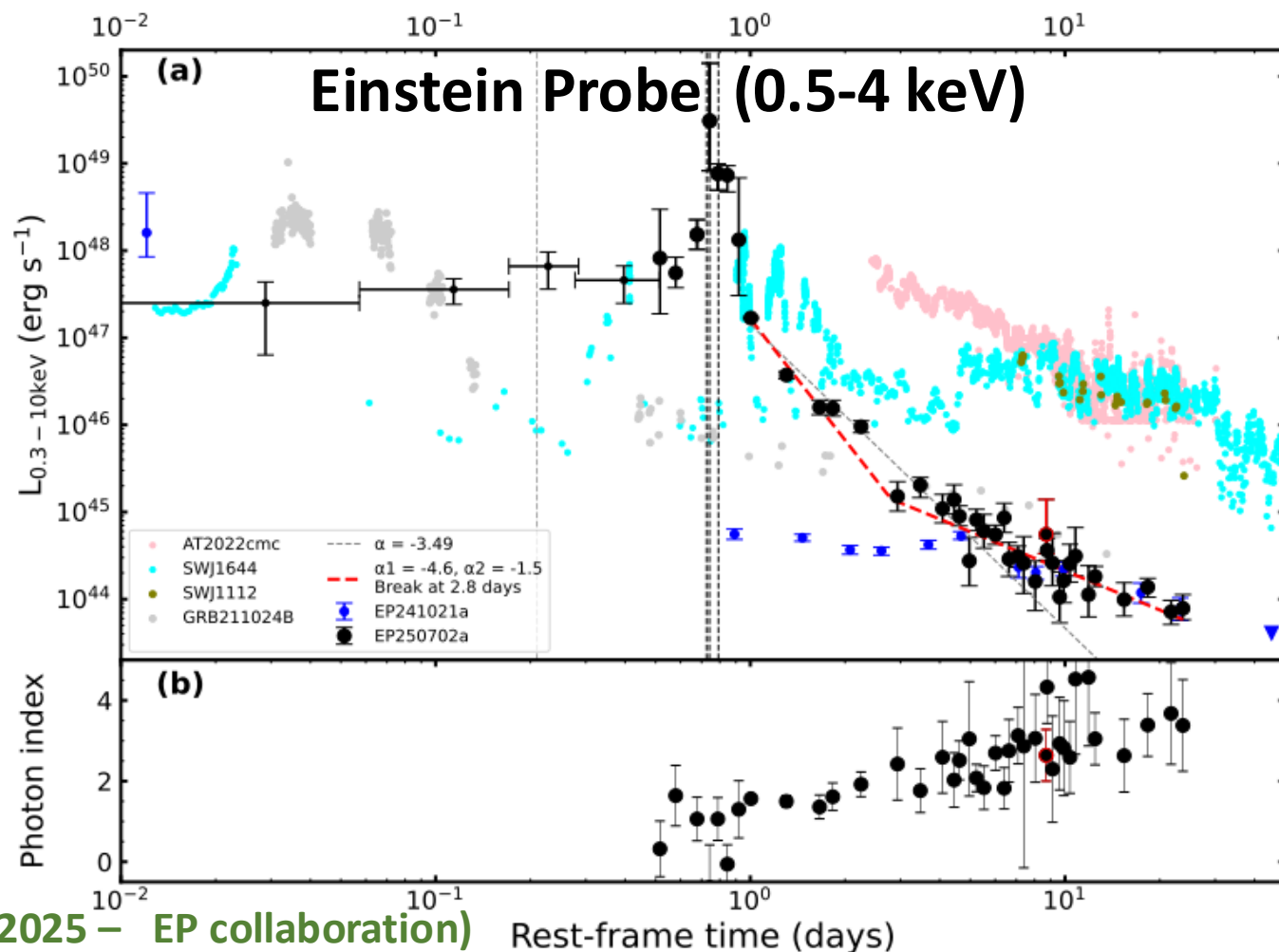
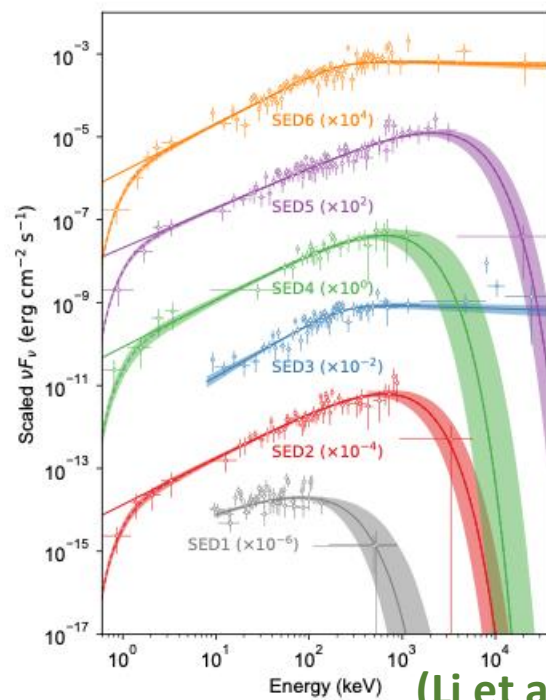
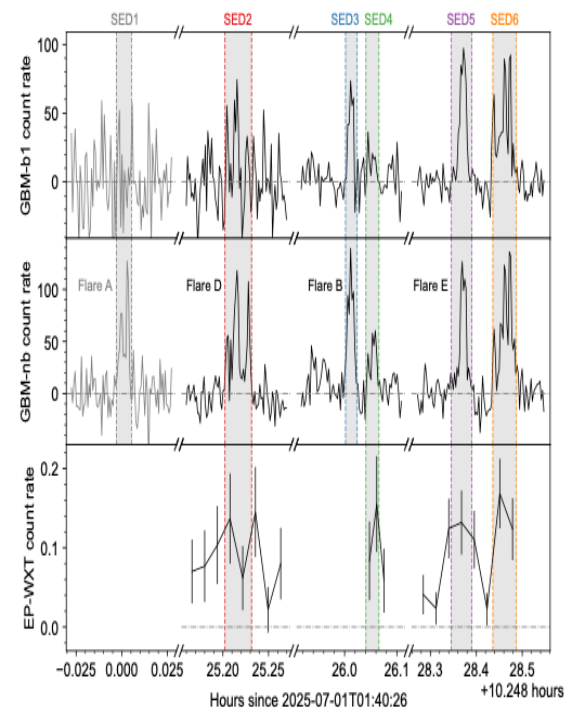
$$E_{X,\text{iso}}(\text{pre-peak}) \sim 10^{52.5} \text{ erg}$$

$$E_{X,\text{iso}}(\text{peak}) \sim 10^{53} \text{ erg}$$

$$E_{\gamma,\text{iso}}(\text{peak}) \geq 1.4 \times 10^{54} \text{ erg}$$

⇒ likely beamed relativistic jet

- Insight-HXMT/HE + Fermi /GBM: ~ 50 s precursor found ~ 25 hr before the main peak (Zhang+26)

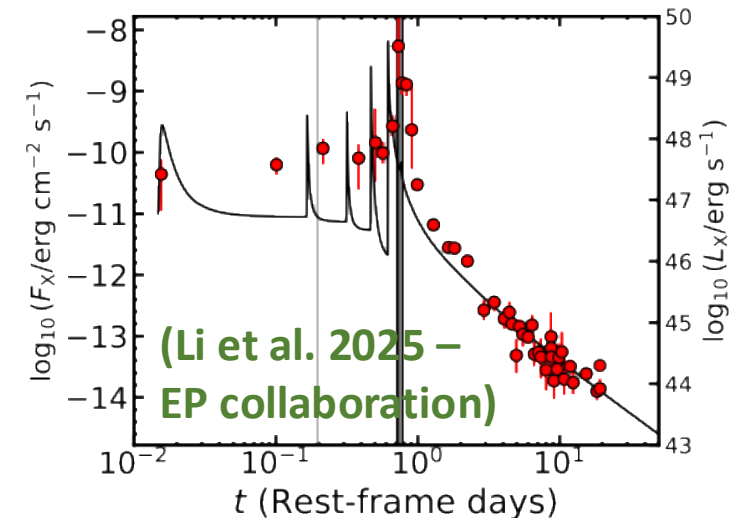
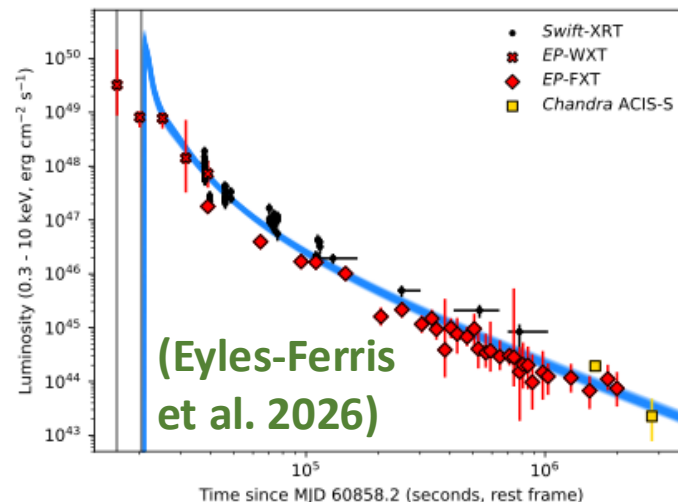


What can it be?

- Unusual type of collapsar (variant of long GRBs)? He-core + NS/BH? (JWST rules out a typical broad-line SN Ic, associated with LGRSs)
- Extreme Ultra-Long GRB? (maybe, but unclear what ULGRBs are...)
- TDE-SMBH: offset from host + $t_{\text{MV},z} \approx 0.5 \text{ s} > \frac{r_g}{c} \Rightarrow M_{\text{BH}} < 5 \times 10^4 M_{\odot}$
- TDE-stellar-mass-BH (micro-TDE): possible (Beniamini, Perets & JG 2025)
- TDE-IMBH: possible (MS or WD?); $t_{\text{MV},z} \approx 0.5 \text{ s}$ & $\Gamma > 56$ still favor a stellar-mass BH engine

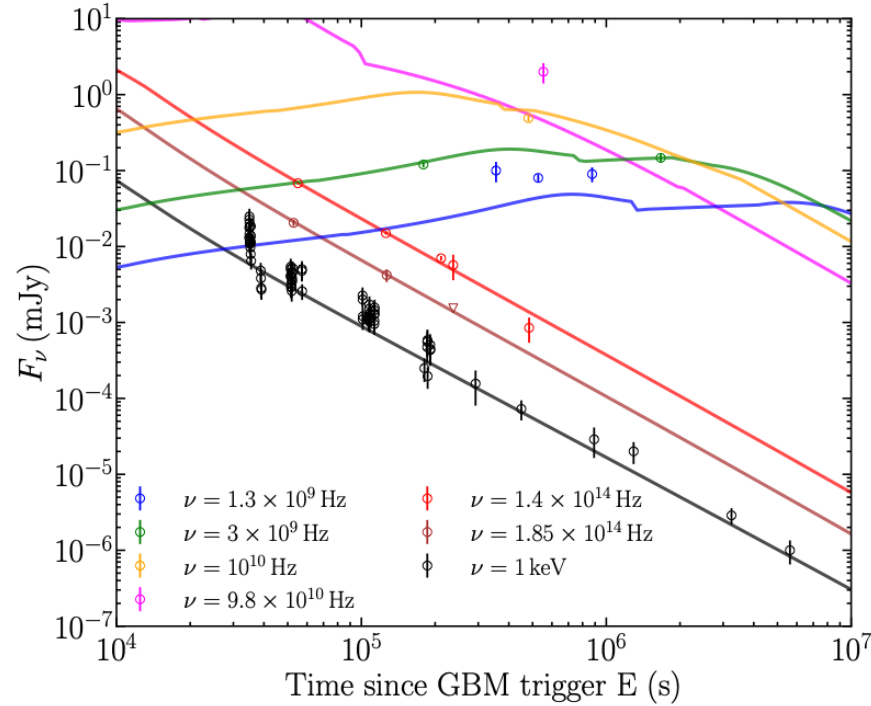
TDE-IMBH of WD:

Timescales & lightcurve shape don't quite match



A milli-TDE Model for GRB250702B: MS Star Disrupted by IMBH (JG, Peters, Gill, Beniamini, O'Connor 2025)

GRB250702B Afterglow Fit:



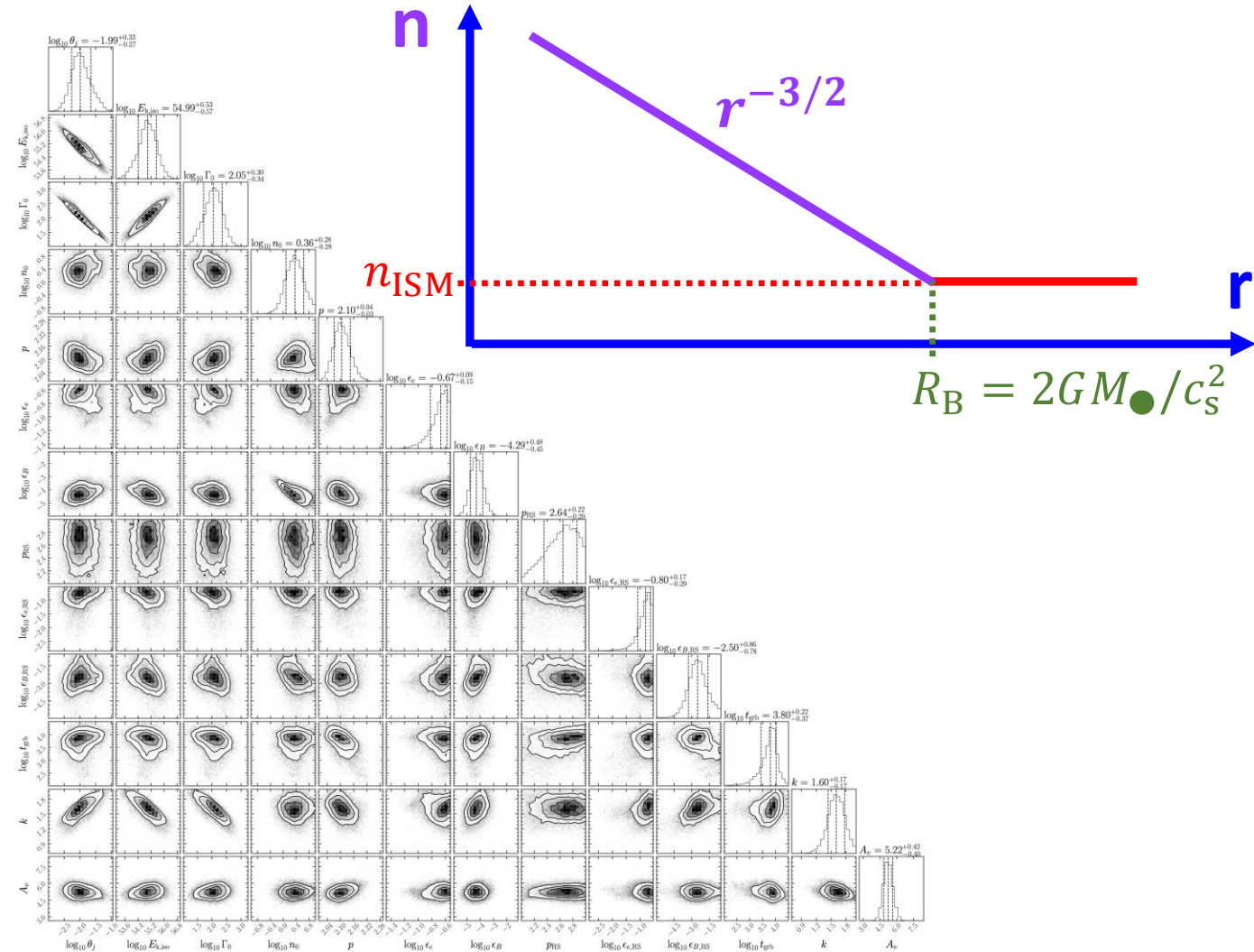
$k = -1.60 \pm 0.17$ (consistent with -3/2)

$$n(r) = n_d(r/R_d)^{-k}, R_d = 10^{18} \text{ cm}$$

$$\log_{10} n_{d,0} = 0.36 \pm 0.28$$

$$(\log_{10} \theta_j = -1.99^{+0.33}_{-0.27}; \log_{10} E_{\text{k,iso}} = 54.99^{+0.53}_{-0.57} \\ 50 \lesssim \log_{10} E_{\text{k}} \lesssim 51)$$

Bondi Accretion:



A milli-TDE Model for GRB250702B: MS Star Disrupted by IMBH (JG, Peters, Gill, Beniamini, O'Connor 2025)

Inferred Bondi Radius:

$$\begin{aligned} R_B &= R_d (n_d/n_{\text{ISM}})^{2/3} \\ &= 0.56^{+0.31}_{-0.19} n_0^{-2/3} \text{ pc} \end{aligned}$$

Inferred IMBH Mass:

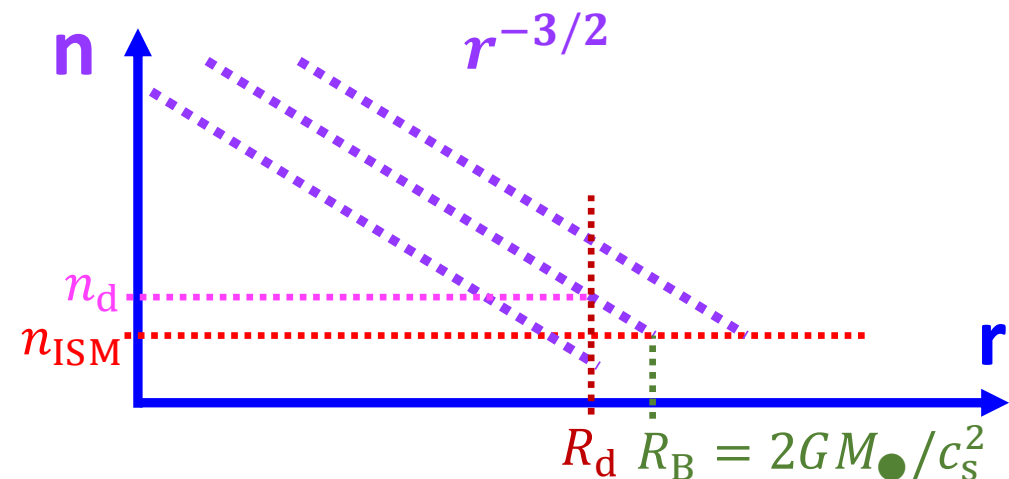
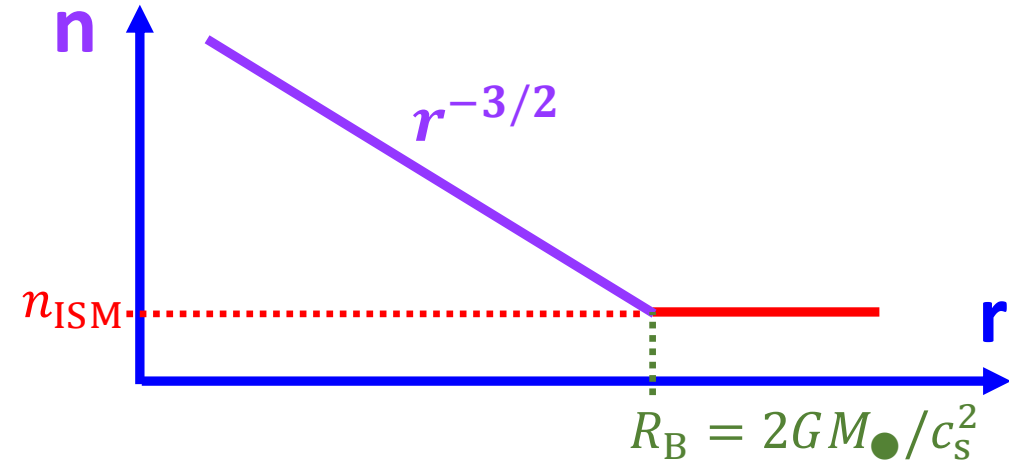
$$\begin{aligned} M_\bullet &= 6.55^{+3.51}_{-2.29} \times 10^3 n_0^{-2/3} c_{s,6}^2 M_\odot \end{aligned}$$

$$k = -1.60 \pm 0.17 \text{ (consistent with } -3/2 \text{)}$$

$$n(r) = n_d (r/R_d)^{-k}, R_d = 10^{18} \text{ cm}$$

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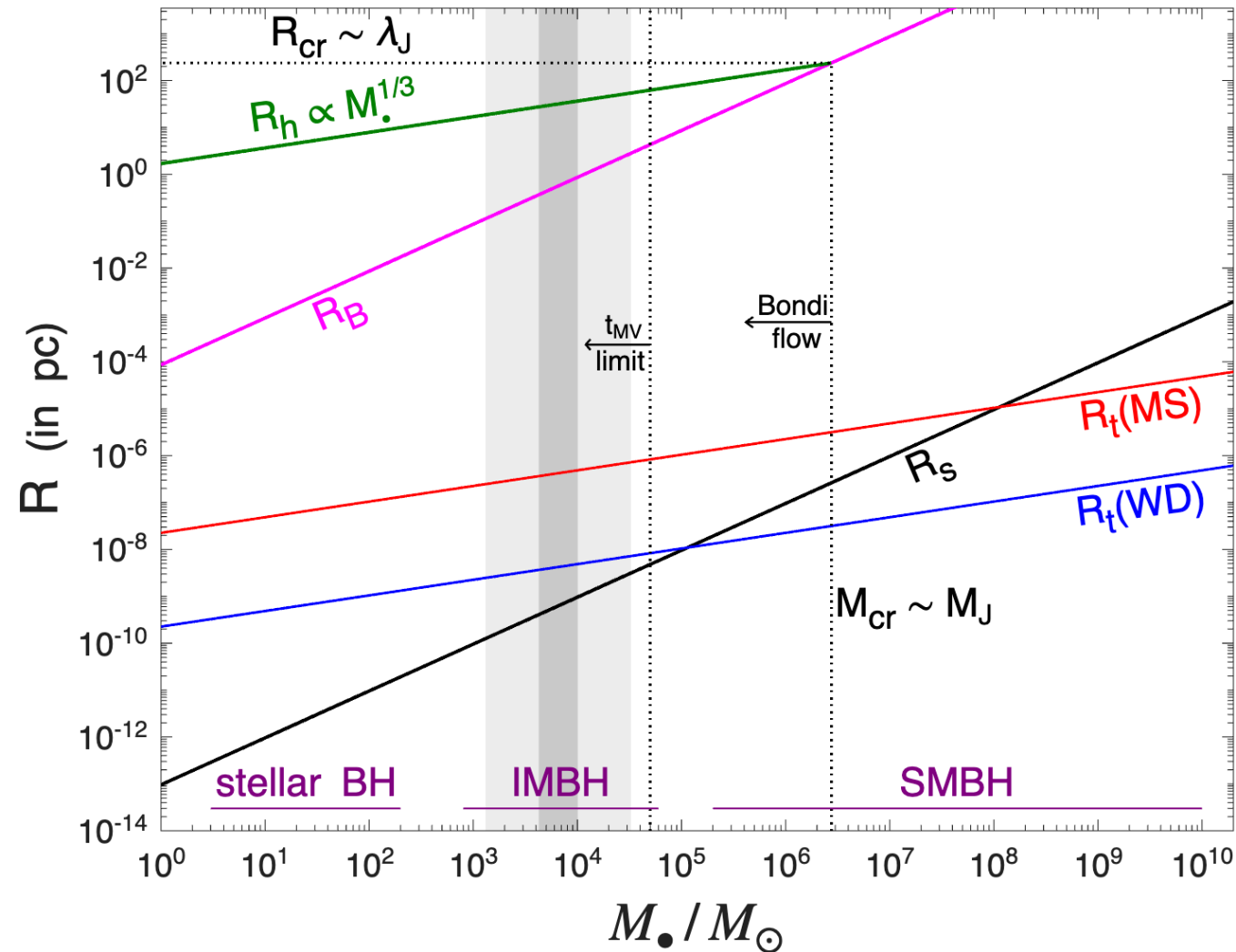
$$M_{\bullet} \\ = 6.55^{+3.51}_{-2.29} \times 10^3 n_0^{-2/3} c_{s,6}^2 M_{\odot}$$

- Also consistent with the observed afterglow emission coming from $r < R_B$

$$k = -1.60 \pm 0.17 \text{ (consistent with } -3/2)$$

$$n(r) = n_d (r/R_d)^{-k}, R_d = 10^{18} \text{ cm}$$

$$\log_{10} n_{d,0} = 0.36 \pm 0.28$$



Bondi Hoyle Lyttleton Accretion:

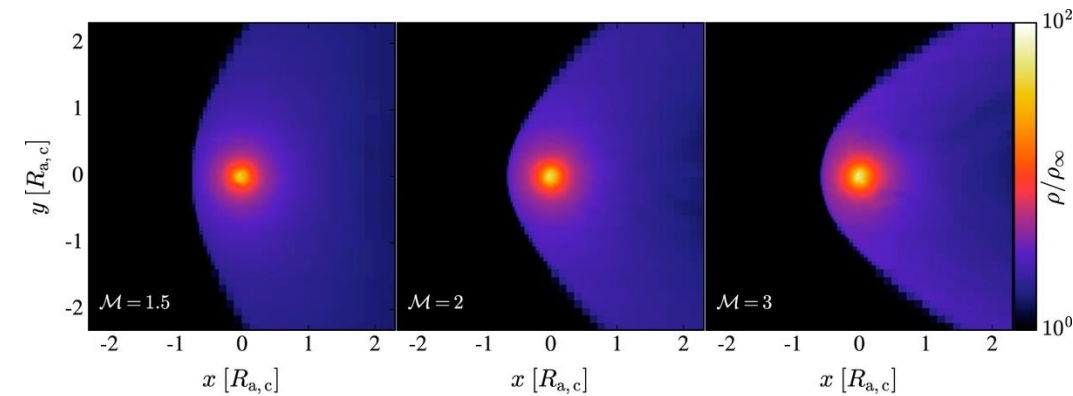
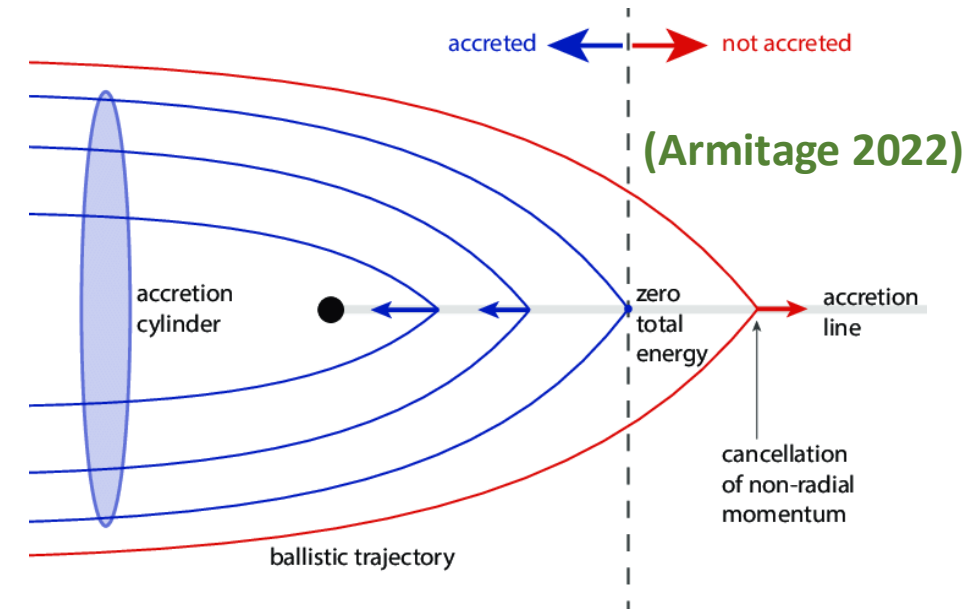
Modified Bondi Radius (v_{BH} added):

$$R_{\text{BHL}} \approx \frac{R_{\text{B}}}{1 + \mathcal{M}^2} = \frac{2GM_{\bullet}}{c_s^2(1 + \mathcal{M}^2)}; \quad \mathcal{M} \equiv \frac{v_{\text{BH}}}{c_s}$$

Modified inferred IMBH Mass:

$$M_{\bullet} = 6.55^{+3.51}_{-2.29} \times 10^3 n_0^{-2/3} c_{s,6}^2 (1 + \mathcal{M}^2) M_{\odot}$$

$$t_{\text{MV},z} > \frac{r_g}{c} \Rightarrow \frac{M_{\text{BH}}}{M_{\odot}} < 5 \times 10^4 \Rightarrow v_{\text{BH}} \lesssim 28 n_0^{1/3} \text{ km/s}$$



(Kaaz, Antoni & Ramirez-Ruiz 2019)

Relevant Timescales: Main Sequence vs. White Dwarf

Disruption to 1st periastron passage ($r_t \rightarrow r_p$):

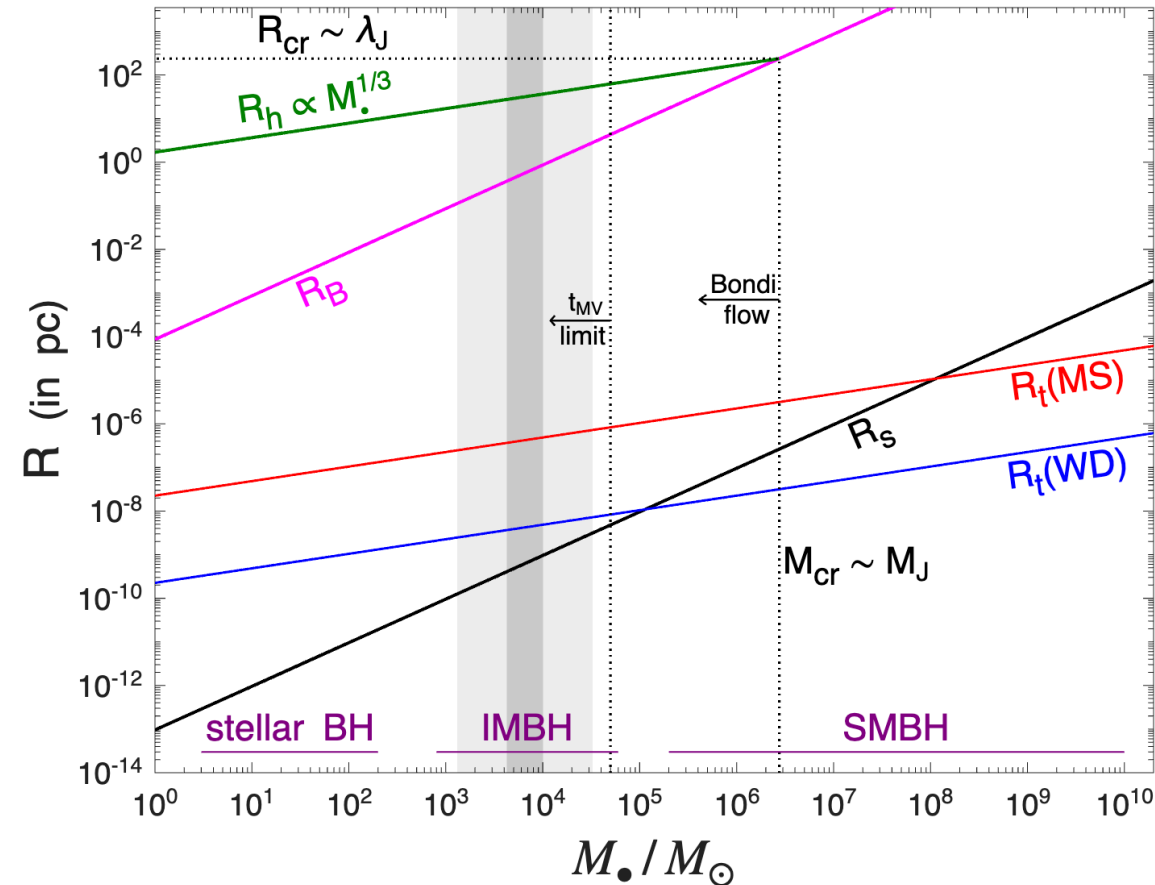
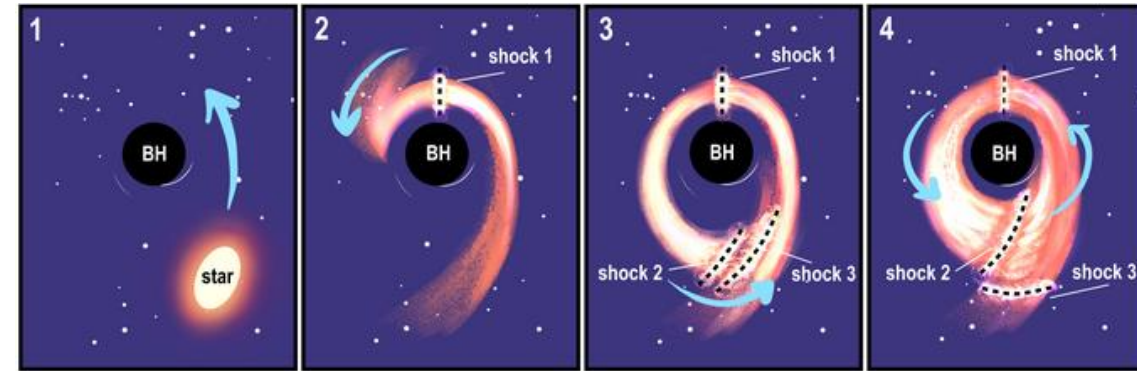
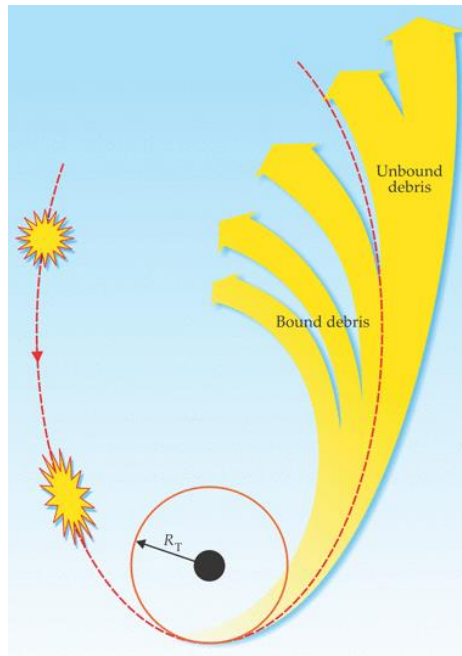
$$t_1 = \sqrt{\frac{\pi^2 R_*^3}{2GM_*}} = \begin{cases} 3.54 \times 10^3 R_{*,0}^{3/2} M_{*,0}^{-1/2} \text{ s (MS)} , \\ 3.54 R_{*,-2}^{3/2} M_{*,0}^{-1/2} \text{ s (WD)} . \end{cases}$$

Orbital time of the most bound debris:

$$t_{\min} \approx \begin{cases} 3.5 \times 10^4 A_{\beta,-1} \eta^2 M_{\bullet,4}^{1/2} R_{*,0}^{3/2} M_{*,0}^{-1} \text{ s (MS)} , \\ 3.5 \times 10^1 A_{\beta,-1} \eta^2 M_{\bullet,4}^{1/2} R_{*,0}^{3/2} M_{*,0}^{-1} \text{ s (WD)} . \end{cases}$$

Where $A_\beta \sim 1$ for frozen in approx, or $\sim \beta^{-3}$ for efficient dissipation near r_p
($\beta \equiv r_t/r_p =$ penetration factor)

- The observed timescales greatly prefer a MS star over a WD
- For a WD $r_p \lesssim r_{\text{ISCO}}$ is possible, in which case no accretion disk or jet would form



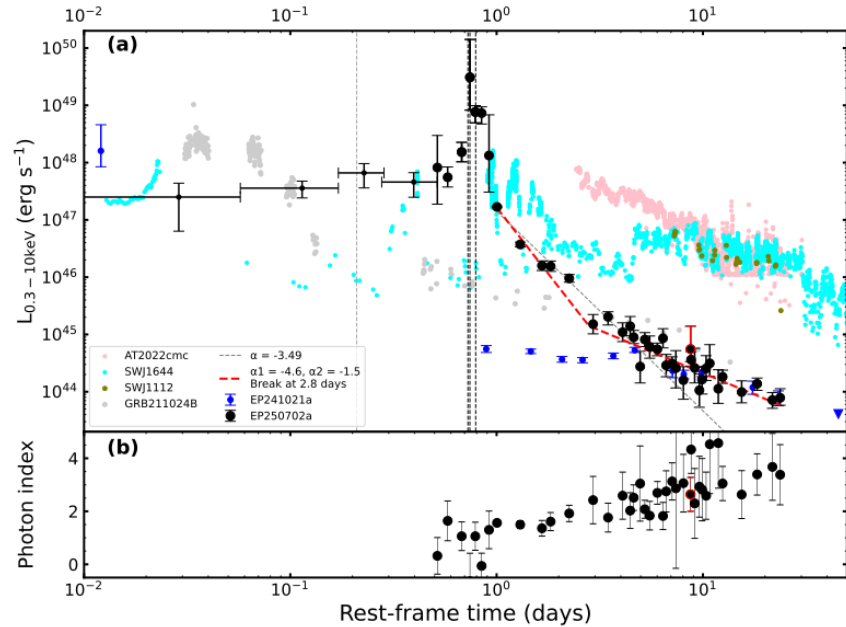
Relevant Timescales: Main Sequence vs. White Dwarf

Disruption to 1st periastron passage ($r_t \rightarrow r_p$):

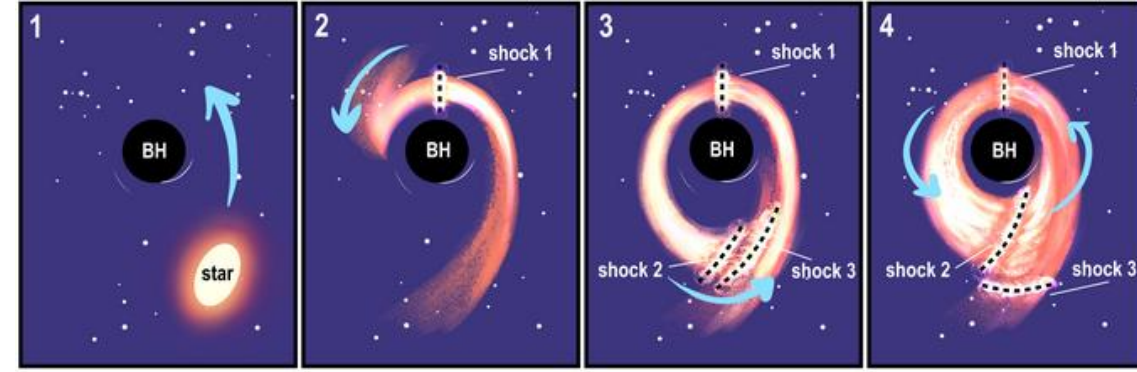
$$t_1 = \sqrt{\frac{\pi^2 R_*^3}{2GM_*}} = \begin{cases} 3.54 \times 10^3 R_{*,0}^{3/2} M_{*,0}^{-1/2} \text{ s} & (\text{MS}) , \\ 3.54 R_{*, -2}^{3/2} M_{*,0}^{-1/2} \text{ s} & (\text{WD}) . \end{cases}$$

Orbital time of the most bound debris:

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(Li et al. 2025 EP collaboration)



For $R_{\text{circ}} \approx 2r_p = 2r_t/\beta$, $t_{\text{acc}} \approx t_{\text{vis}}(R_{\text{circ}})$:

$$t_{\text{acc}} \approx 4.5 \times 10^4 \beta^{-3/2} \alpha_{-1}^{-1} h^{-2} M_{*, 0}^{-1/2} R_{*, 0}^{3/2} \text{ s}$$

Approximate rise and peak timescales:

$$t_{\text{X, rise}} \sim t_{\text{circ}} + t_{\text{acc}} , \quad t_{\text{main}} \sim t_{\min} + t_{\text{acc}} ,$$

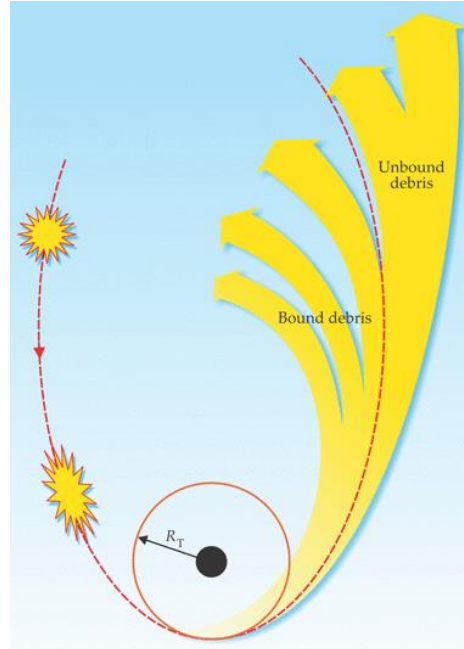
$$t_{\text{circ}} \approx 3.5 \times 10^4 f_{0.5} \beta_{0.5}^{-3} \eta^2 M_{\bullet, 4}^{1/2} R_{*, 0}^{3/2} M_{*, 0}^{-1} \text{ s} ,$$

$$t_{\min} \approx 1.1 \times 10^4 \beta_{0.5}^{-3} \eta^2 M_{\bullet, 4}^{1/2} R_{*, 0}^{3/2} M_{*, 0}^{-1} \text{ s} ,$$

$$t_{\text{acc}} \approx 3.2 \times 10^4 \beta_{0.5}^{-3/2} \alpha_{-1}^{-1} h_{-0.3}^{-2} M_{*, 0}^{-1/2} R_{*, 0}^{3/2} \text{ s} ,$$

Energetics limit on the jet beamin factor:

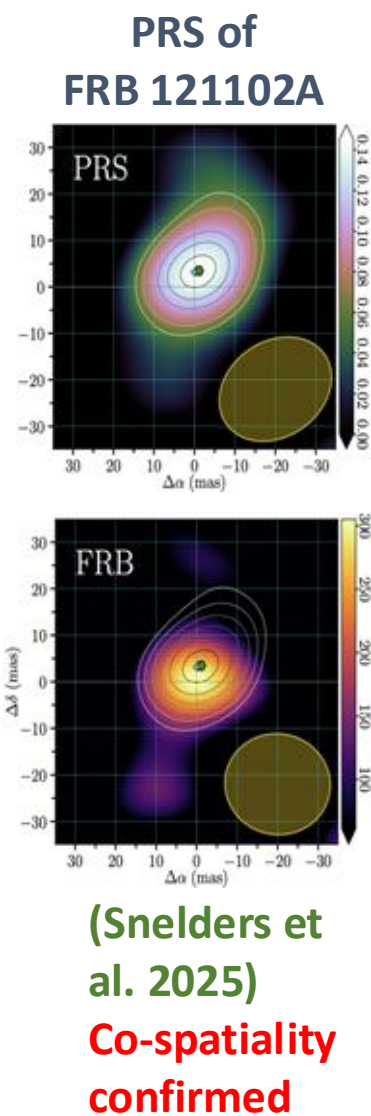
$$f_b \leq 6.4 \times 10^{-4} \eta_{j, -2} \eta_{\gamma, -1} M_{*, 0} \left(\frac{f_{\text{fb}}}{0.5} \right) \left(\frac{1.4 \times 10^{54} \text{ erg}}{E_{\gamma, \text{iso}}} \right)$$



Persistent Radio Sources (PRSs) of Repeating FRBs: Implications for Magnetar Progenitors (Rahaman, Acharia, Beniamini & JG 2025)

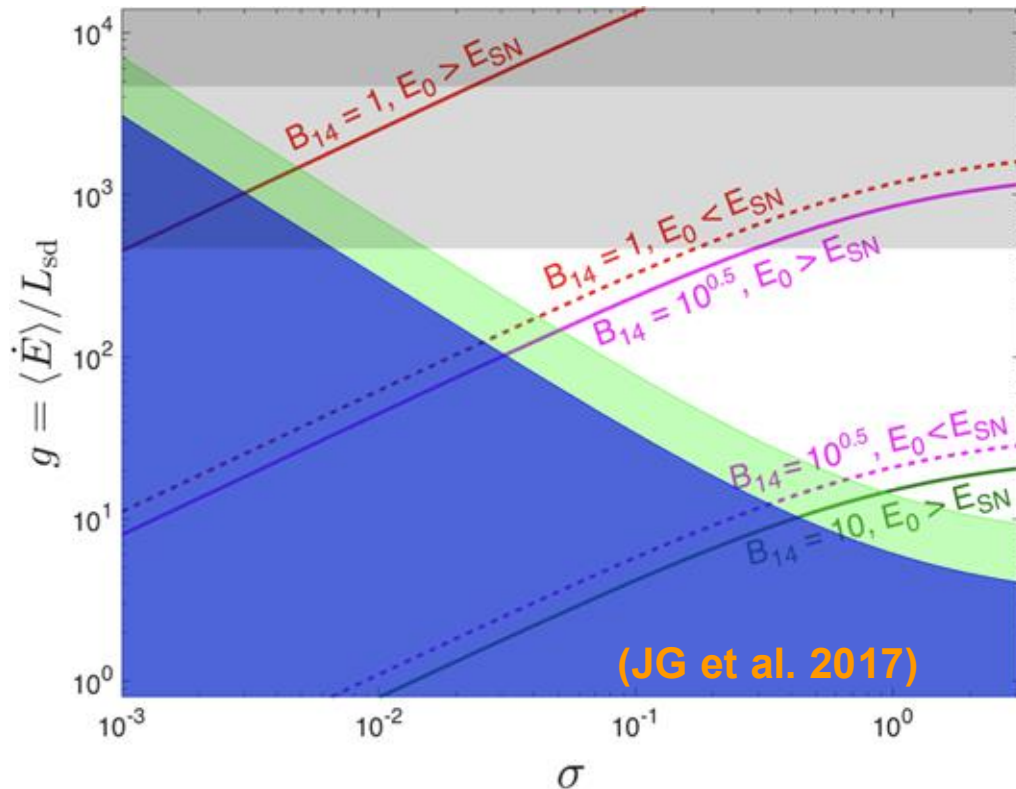
- A handful of confirmed/candidate PRSs associated with repeating FRBs
- Appear to have large DM_{host} , large + variable R, low-Z high-SFR hosts
- Model: synchrotron emission from a compact Magnetar Wind Nebula (MWN)

Property	FRB 20121102A	FRB 20190417A	FRB 20190520B	FRB 20201124A	FRB 20240114A
z (redshift)	0.193	0.128	0.241	0.098	0.130
Host galaxy	Dwarf	Dwarf	Dwarf	Spiral	Dwarf
DM [pc cm ⁻³] DM _{host,rest}	558 ≲ 203	1379 > 1228	1204 137 – 707	413 150 – 220	528 142 ± 107
RM _{rest} [rad m ⁻²]	(0.44 – 1.5) · 10 ⁵	(5.04 – 6.44) · 10 ³	(–3.6 – 2.0) · 10 ⁴	–661 ± 42	449 ± 13
offset _{PRS-FRB} [pc]	< 40	< 26	< 80	< 188	≲ 28
R _{proj} [pc]	< 0.7 (at 5 GHz)	< 23	< 9	< 700	< 0.4
ν _{obs} [GHz]	1 – 26	1.4	1.5, 3, 5.5	6, 15, 22	0.65, 1.3, 5
F _ν [μJy]	180 (at 3 GHz)	190 (at 1.4 GHz)	202 (at 3 GHz)	8,20,30 (6,15,22)	66,72,46 (.65,1.3,5)
α (spectral index)	–0.2 – –1	–1.2 ± 0.4	–0.41 ± 0.04	1.00 ± 0.43	–0.34 ± 0.21
L _ν [erg s ⁻¹ Hz ⁻¹]	2 × 10 ²⁹ (1.4 GHz)	8 × 10 ²⁸ (1.4 GHz)	3 × 10 ²⁹ (1.7 GHz)	2 × 10 ²⁷ (6 GHz)	2 × 10 ²⁸ (5 GHz)

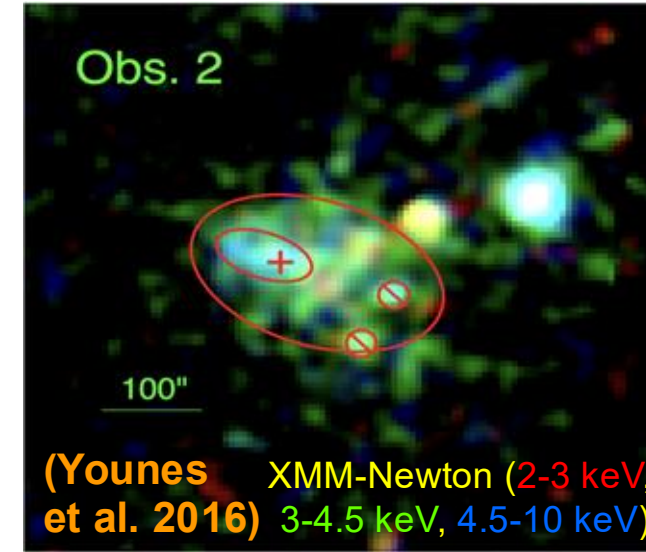


One Known Galactic Magnetar Wind Nebula: **Swift J1834–0846**

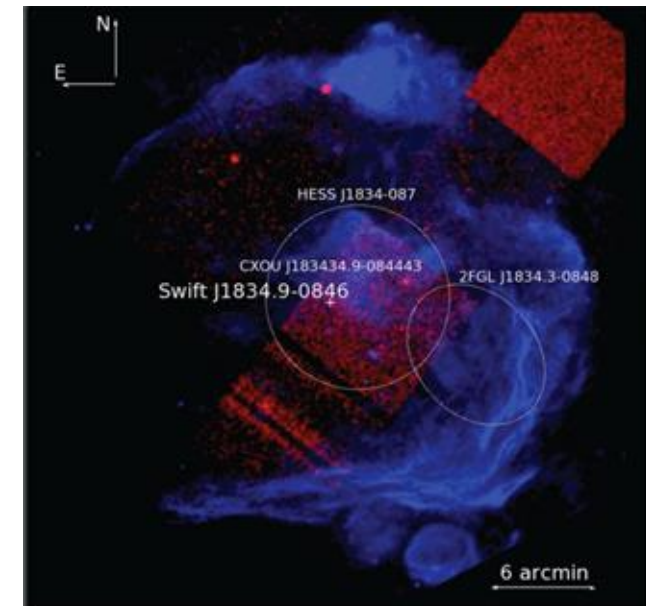
- Quite rare (1 out of over 30 Galactic magnetars)
- MWN size: diffusion-dominated cooling length of X-ray emitting e^\pm
- Spindown power or B_{dipole} decay cannot power the MWN
- B_{int} decay can power the MWN for a current $B_{\text{int}} \gtrsim 10^{15.5}$ G through outflows associated with bursting activity (e.g. giant flares)



$$\begin{aligned}
 P &= 2.48 \text{ s} \\
 \dot{P} &= 7.96 \times 10^{-12} \text{ s s}^{-1} \\
 \tau_c &= 4.9 \text{ kyr} \quad (n = 3) \\
 t &= \tau_c - 10^5 \text{ yr} \\
 B_d &= 10^{14} \text{ G} \\
 L_{\text{sd}} &= 2 \times 10^{34} \text{ erg s}^{-1}
 \end{aligned}$$

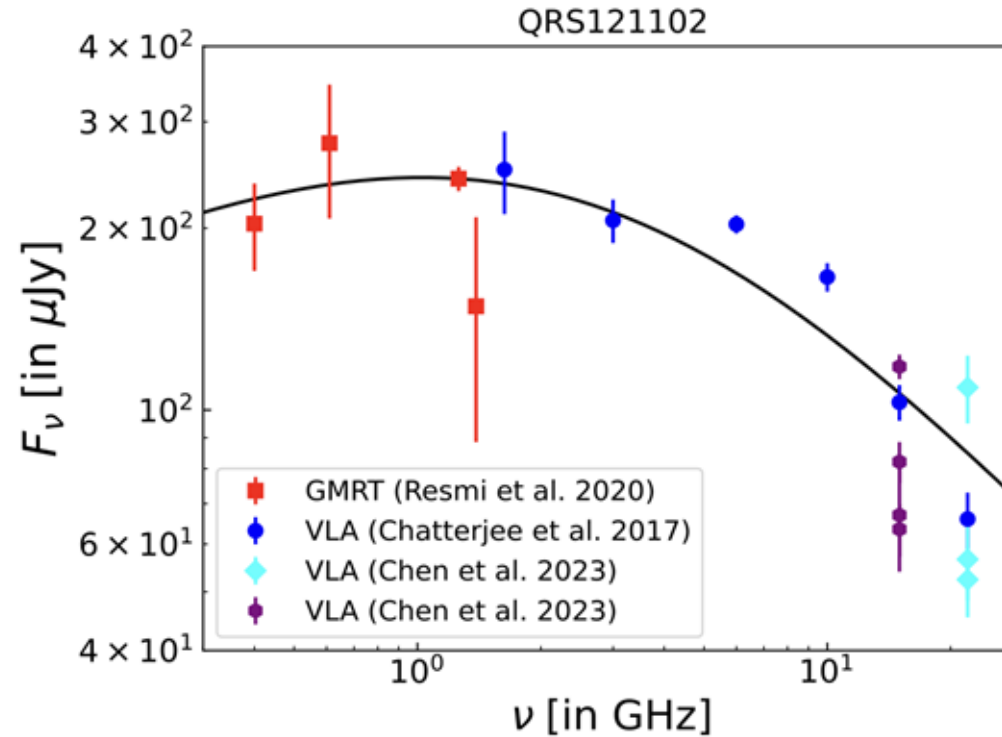
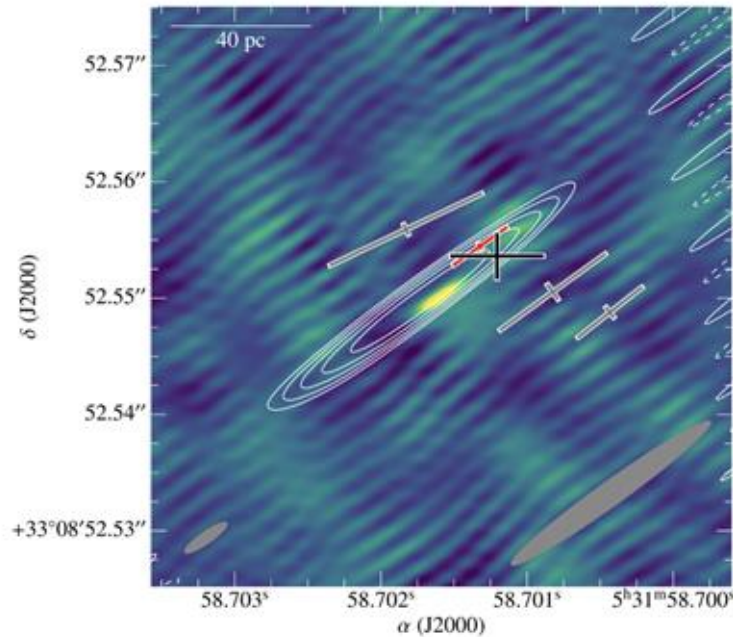


Associated with SNR W41



Persistent Radio Sources (PRSs) of Repeating FRBs: Implications for Magnetar Progenitors (Rahaman, Acharia, Beniamini & JG 2025)

PRS of FRB 121102A: the best constrained source



$$L_\nu \sim 2$$

$$\times 10^{29} \frac{\text{erg}}{\text{s Hz}}$$

$$\nu_{\text{sa}} < 0.5 \text{ GHz}$$

$$\nu_{\text{m}} < 3 \text{ GHz}$$

$$\nu_{\text{c}} > 22 \text{ GHz}$$

$$R_{\text{max}} = 0.7 \text{ pc (imaging)}$$

(Marcote et al. 2017)

$$R_{\text{eq}} \sim 0.1 \text{ pc (equipartition)}$$

$$R_{\text{min}} = 0.03 \text{ pc (scintillation)}$$

(Chen et al. 2023)

$$R_{\text{min}} < R_{\text{eq}} < R_{\text{max}}$$

PRS is very compact

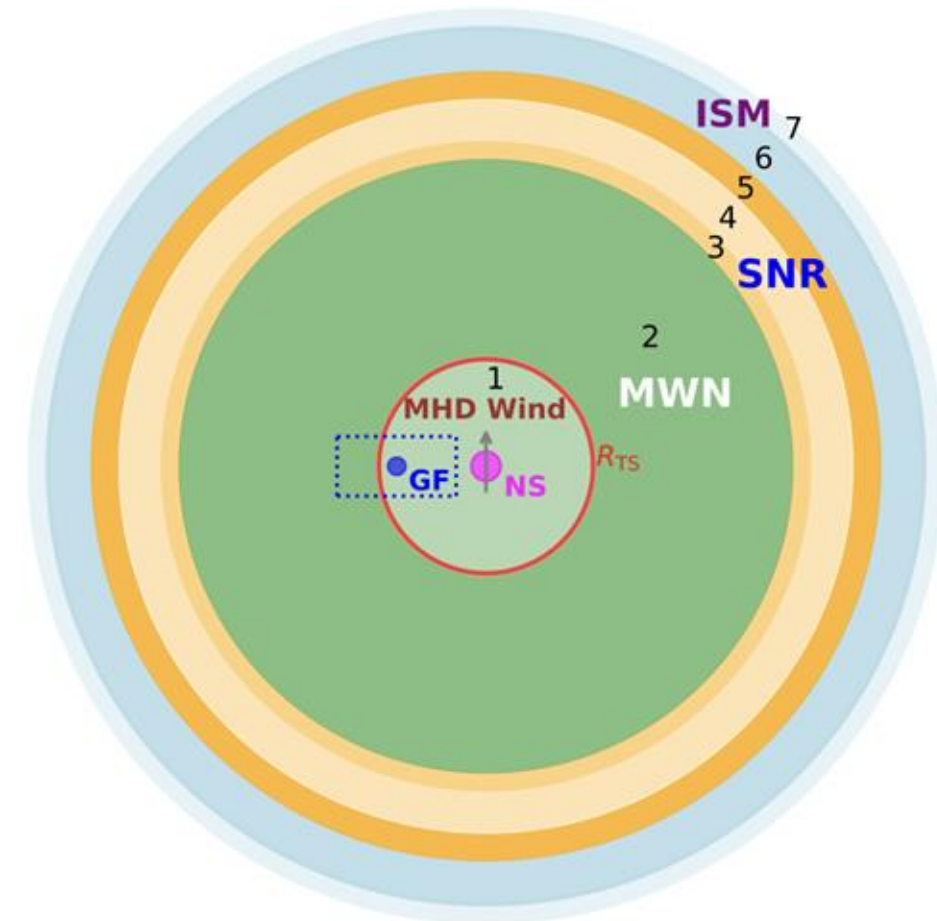
Persistent Radio Sources (PRSs) of Repeating FRBs: Implications for Magnetar Progenitors (Rahaman, Acharia, Beniamini & JG 2025)

- Magnetar Wind Nebula (**MWN** = PRS candidate) is **confined** in a SuperNova Remnant (**SNR**)
- **Can a millisecond-magnetar work?** (Murase et al. 2016; Metzgar et al. 2017; Margalit & Metzgar 2018; Omand et al. 2018; Murase et al. 2021; Bhattacharya et al. 2024)
- **No** – the **compact size** and **minimal age** exclude this!!!

$$t = \frac{R_{\text{SNR}}}{v_{\text{SNR}}} = R_{\text{SNR}} \sqrt{\frac{M_{\text{SNR}}}{2(E_{\text{rot}} + E_{\text{SN}})}}$$
$$\approx \begin{cases} 2.5 R_{17.3} M_3^{\frac{1}{2}} P_{i,-3} \text{ yr} & \text{for } E_{\text{rot}} > E_{\text{SN}}, \\ 63.4 R_{17.3} M_{10}^{\frac{1}{2}} E_{50}^{-\frac{1}{2}} \text{ yr} & \text{for } E_{\text{rot}} < E_{\text{SN}}. \end{cases}$$

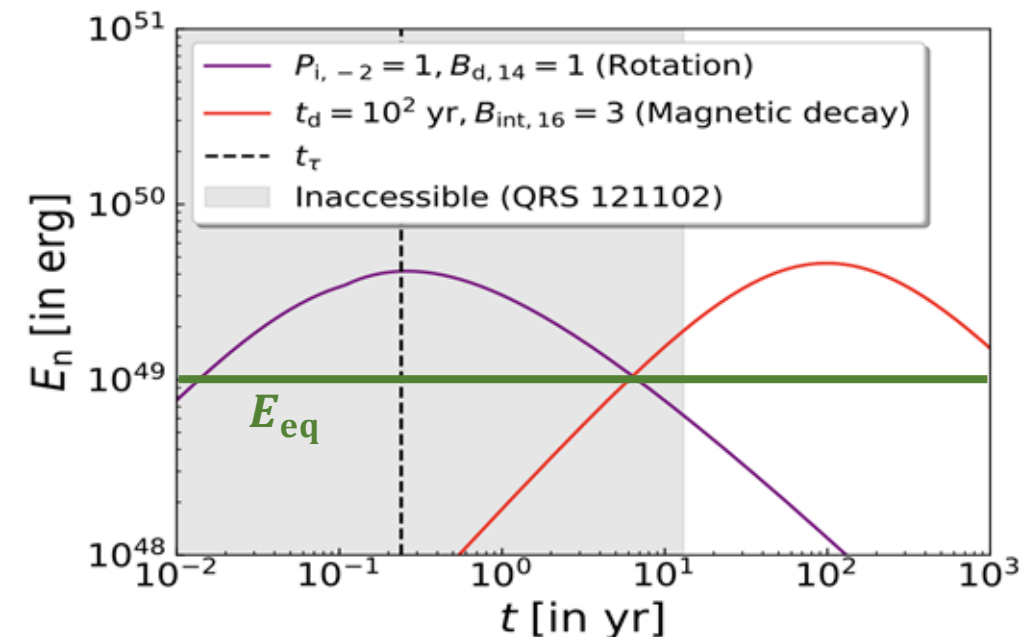
$$E = E_{\text{SN}} + E_{\text{rot}} \quad (\text{Total SNR Energy})$$

This source is already observed for over 13 years

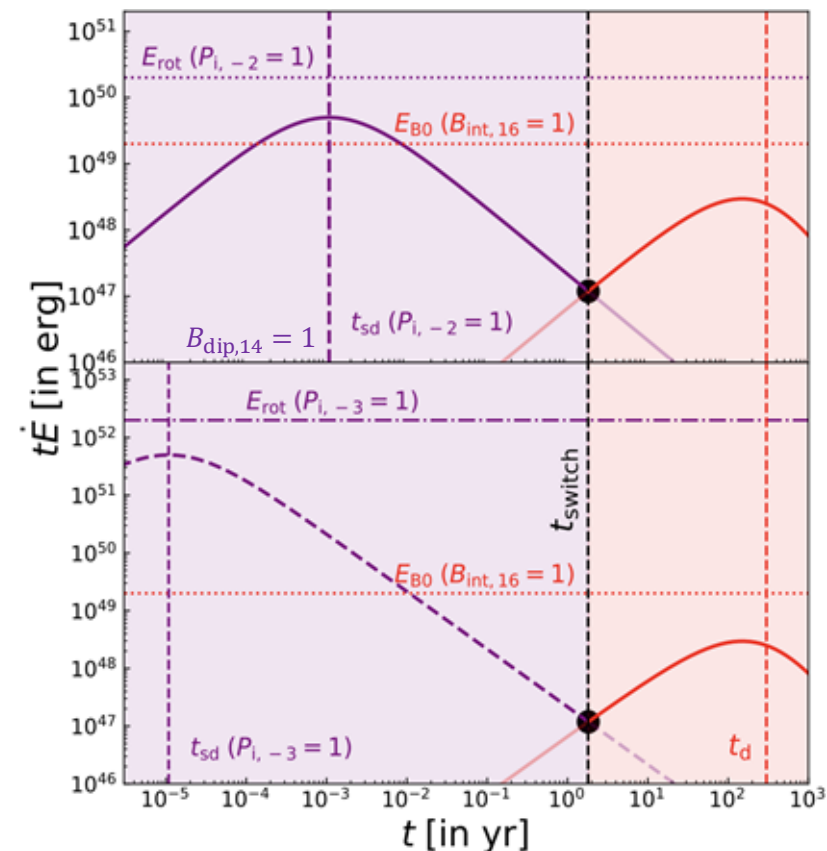
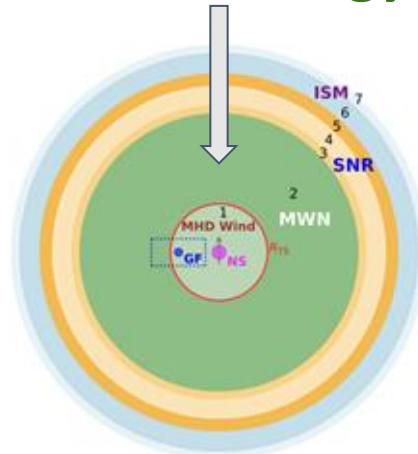


Persistent Radio Sources (PRSs) of Repeating FRBs: Implications for Magnetar Progenitors (Rahaman, Acharia, Beniamini & JG 2025)

- Equipartition/minimum nebular energy: $E_{\text{eq}} \sim 10^{49}$ erg
- To power the FRB: $B_{\text{dip}} \geq 10^{14}$ G (Lu & Kumar 2018, Beniamini & Kumar 2025)
- For an age of $t > 13$ years spindown cannot power the MWN
- After a characteristic time, t_{switch} , the nebular energy input is dominated by B_{int} decay
- We require B_{int} decay time: $t_d \sim 10^{2.5}$ yr $\gg t_{\text{sd}}$ (Other B-powered models need months: Murase+16, Metzger+17, Margalit & Metzgar 18, Omand+18, Murase+21, Bhattacharya+24)



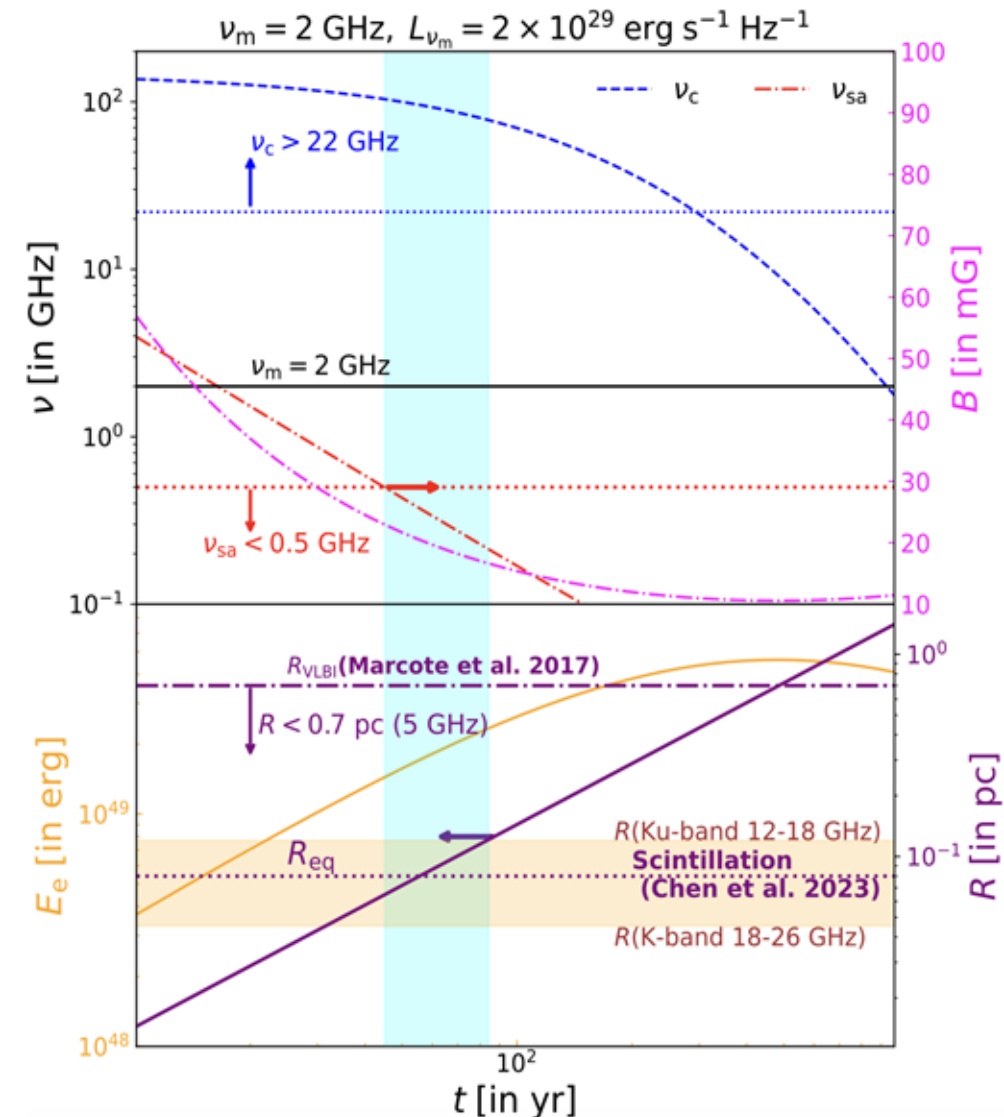
Nebular energy



Persistent Radio Sources (PRSs) of Repeating FRBs: Implications for Magnetar Progenitors (Rahaman, Acharia, Beniamini & JG 2025)

Model favored by all of the observations:
Extreme magnetar & weak SN explosion

- Extreme initial internal B-field:
 $B_{\text{int}} \sim (1 - 3) \times 10^{16} \text{ G}$ with a rather small decay time $t_d \sim 10^2 - 10^{2.5} \text{ yr}$
- Weak SN explosion: $E_{\text{SN}} \sim 10^{50} - 10^{51} \text{ erg}$ given to an ejected mass of $M_{\text{ej}} \sim (3 - 10) M_{\odot}$
- Age of PRS/FRB source: $13 \text{ yr} < t \lesssim 100 \text{ yr}$
- The slowest allowed B_{int} decay time $t_{d,\text{max}} \sim 500 \text{ yr}$ favors a sub-energetic SN explosion $E_{\text{SN}} \sim 10^{50} \text{ erg}$ with $M_{\text{ej}} \gtrsim 10 M_{\odot}$ & a low-ionization fraction ($\sim 3\%$)
- Similar results hold for the PRS of FRB 20190520B
- PRS of FRB 20201124A is rather poorly constrained



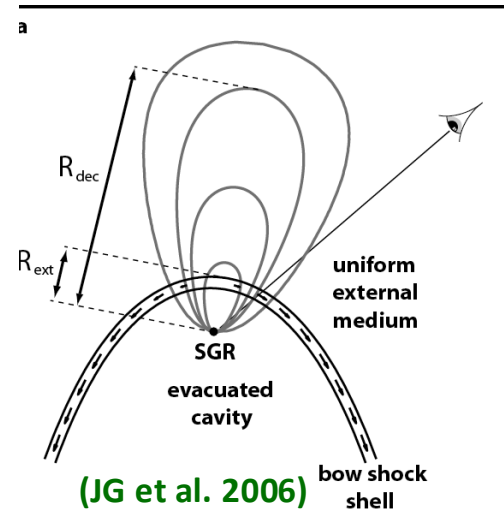
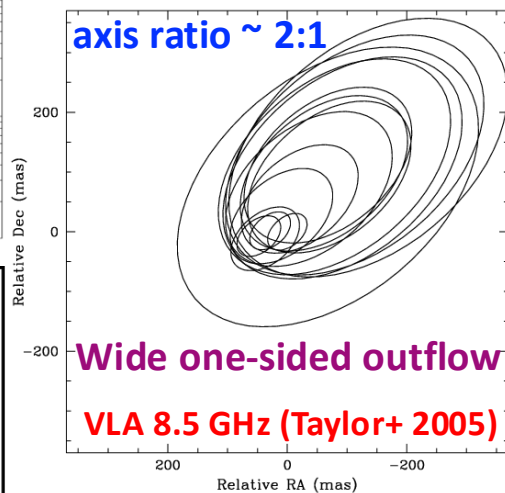
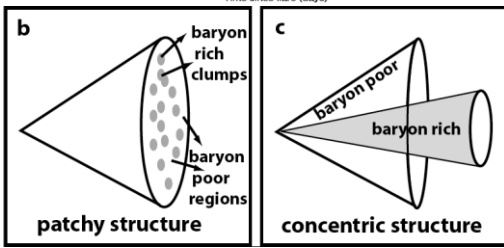
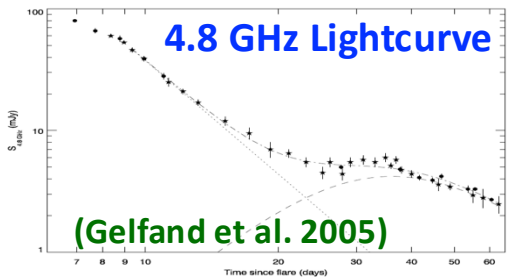
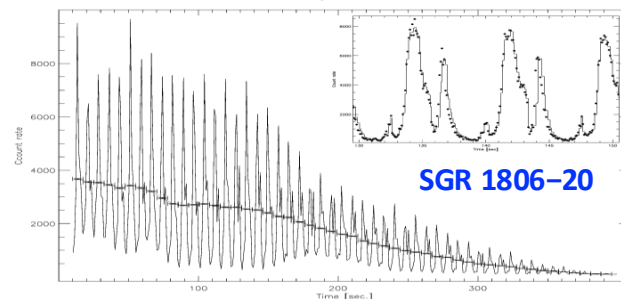
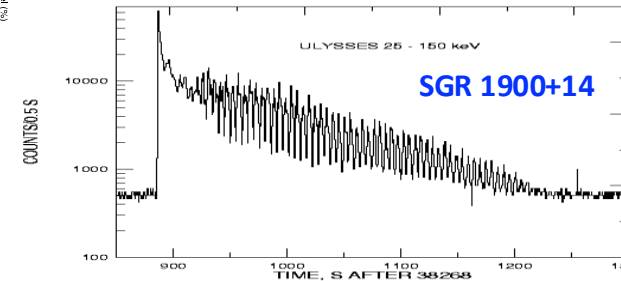
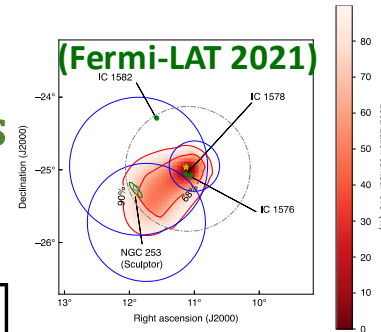
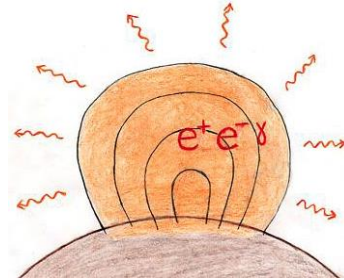
Extragalactic Magnetar Giant Flares (GFs)

- There are **3 known Galactic GFs** (including the 5.3.1979 GF from SGR 0526-66 in the **LMC**)
- **Similar energy in pulsating tail** ($\sim 10^{44}$ erg): $e^+e^-\gamma$ that is trapped on closed field lines
- The **initial spike energy varies greatly** ($E_{\text{spike}}/E_{\text{tail}} \sim 1 - 10^{2.5}$): what cannot be trapped
- 2 of the 3 Galactic GFs created a **radio nebula**, implying $u = \Gamma\beta \sim 1$ outflow: $E_k \sim E_{\text{spike}}$
- Recently: 10 good candidates for **extragalactic GFs** (only the initial spike is detectable)
- **Sculptor galaxy** (3.5 Mpc): $\Gamma \sim 100$, $E_k \sim E_{\text{spike}} \sim 10^{46.5}$ erg (similar to SGR 1806-20)

Observed delay of emission from outflow collision with an

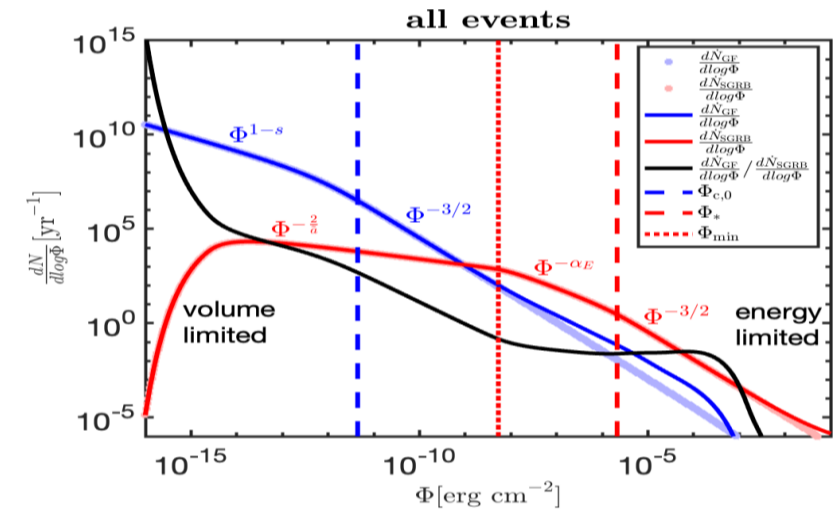
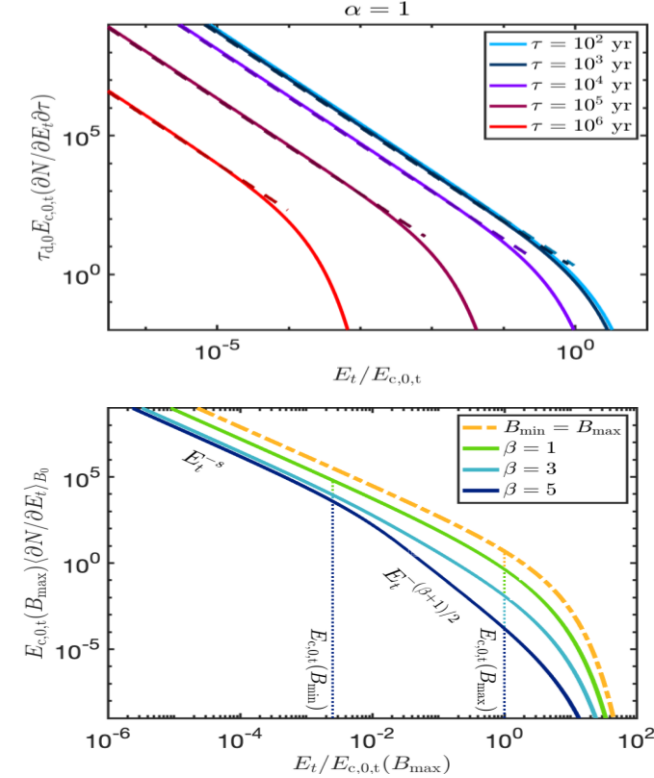
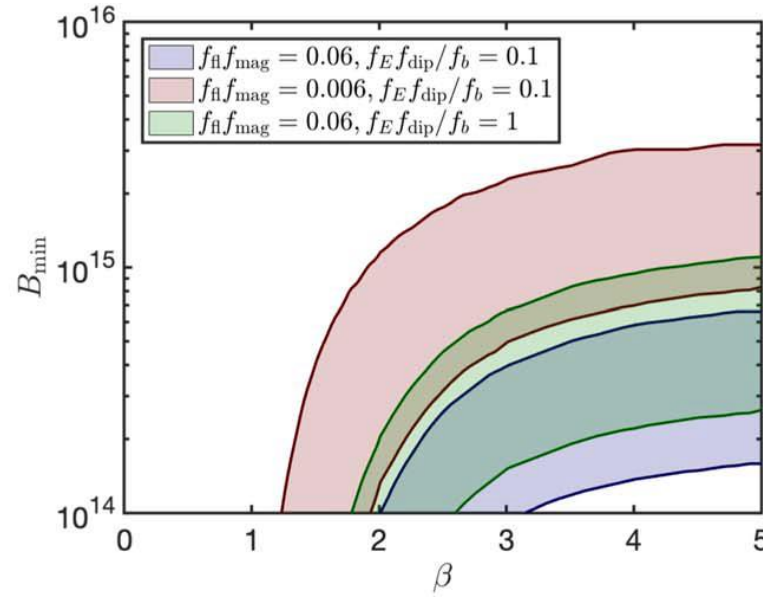
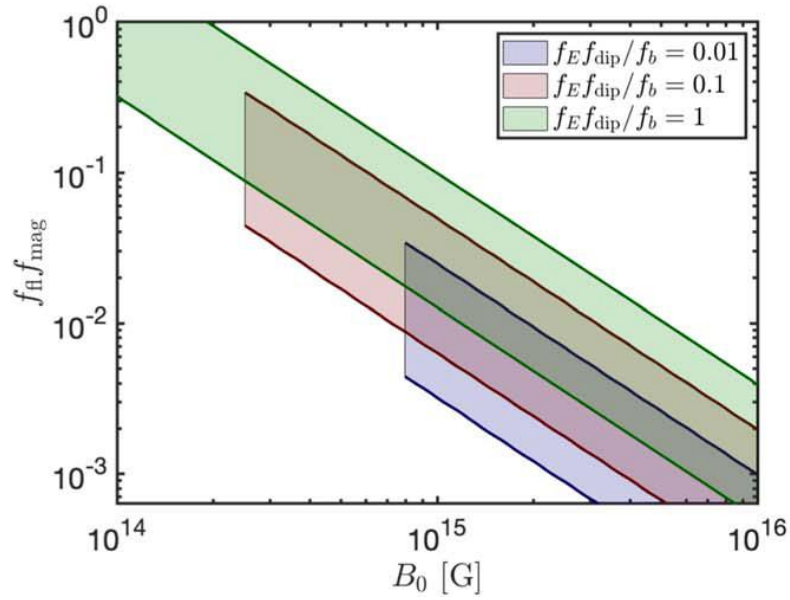
external bow-shock sell: $\Delta t_{\text{obs,r}} = \frac{R_{\text{ext}}(1-\beta)}{\beta c} \sim \frac{R_{\text{ext}}}{2c\Gamma^2} \sim 17R_{16}\Gamma_2^{-2} \text{ s}$

(LAT \sim GeV photons at 19, 180, 284 s)



Extragalactic Magnetar Giant Flares (GFs) (Beniamini et al. 2025)

- Each magnetar is born with **initial B-field** B_0 of energy $E_{B,0}$, which **decays** on timescale $\tau_{d,0}$ and powers the GFs: $\int dE_t E_t \frac{\partial^2 N}{\partial E_t \partial \tau} = f_{fl} |\dot{E}_B|$
- Power law energy distribution**: $\frac{\partial^2 N}{\partial E_t \partial \tau} \propto E_t^{-s} e^{-E_t/E_{c,t}}$ ($s \approx 1.7$ from obs.)
with **cutoff energy**: $E_{c,t} = f_E f_{dip} E_B(\tau)$, $f_{dip} = E_{dip}/E_B$, $f_E = E_{c,t}/E_{dip}$
- Allow **beaming**: observed isotropic energy $E = E_t/f_b$ (E_t = true energy)
- Allow for a **distribution in B_0** : $P(B_0) \propto B_0^{-\beta}$ $B_{min} < B_0 < B_{max}$
- Detailed predictions**: can constrain magnetar properties by fits to data



The End