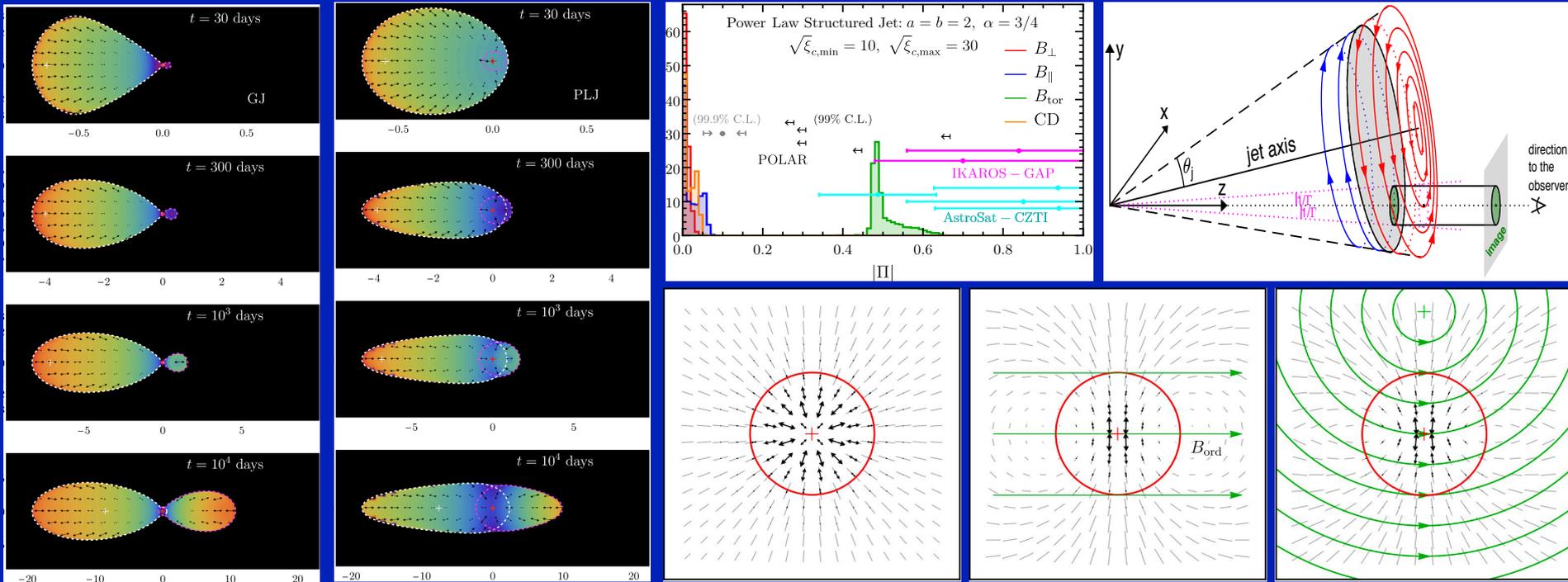


Polarization in Gamma-Ray Bursts

Jonathan Granot

Open University of Israel & George Washington University

Collaborators: R. Gill, A. Königl, F. De Colle, E. Ramirez-Ruiz, T. Piran



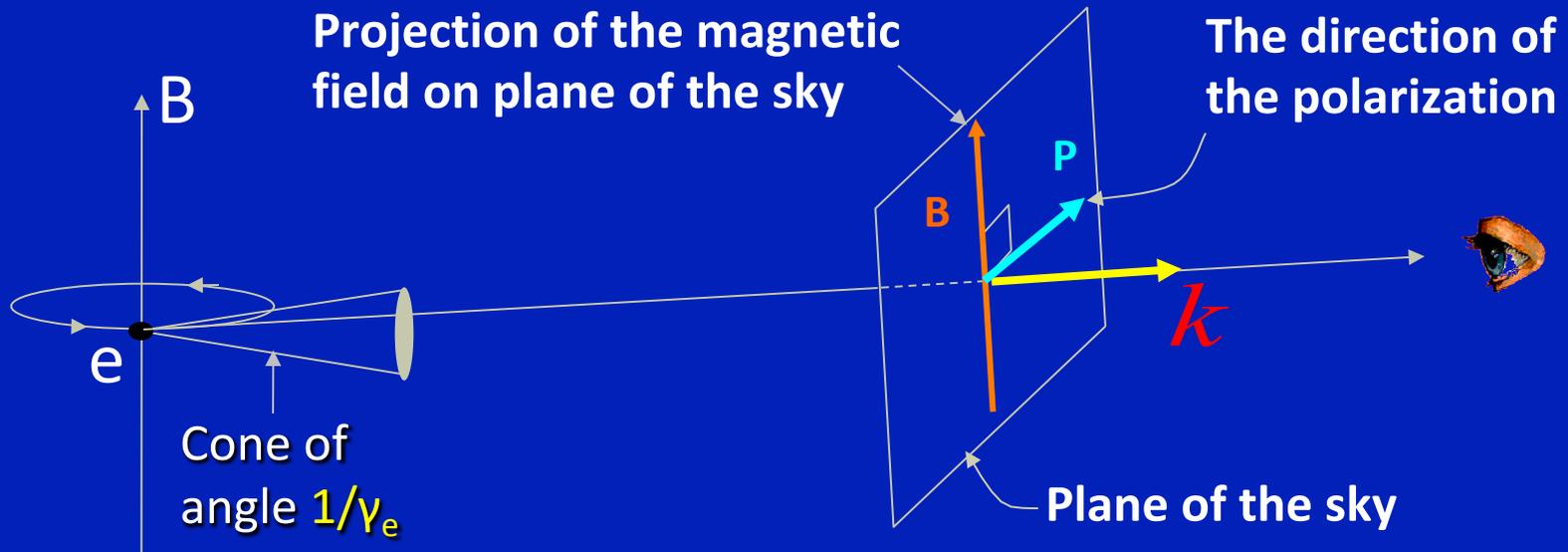
Gamma-Ray Bursts & Related Astrophysics in Multi-Messenger Era

Nanjing, China, 13 May 2019

Outline of the Talk:

- Polarization of synchrotron rad. from a relativistic source
- **Afterglow: Jet structure & dynamics, B-field structure (ES)**
 - ◆ Top hat vs. structured jet
 - ◆ **Shock-produced vs. ordered B-field**, or combining the two
 - ◆ Shock-produced B-field's degree of anisotropy
- Reverse shock: optical flash & **radio flare** (ejecta B-field)
- Prompt GRB: emission mechanism, Jet structure, ejecta B
 - ◆ **High P: Syn. + ordered B vs. sharp jet + special viewing angle**
 - ◆ **Different emission mechanisms**
 - ◆ **What can be learned from single GRBs or a large sample**
- **Conclusions**

Polarization of Synchrotron Emission

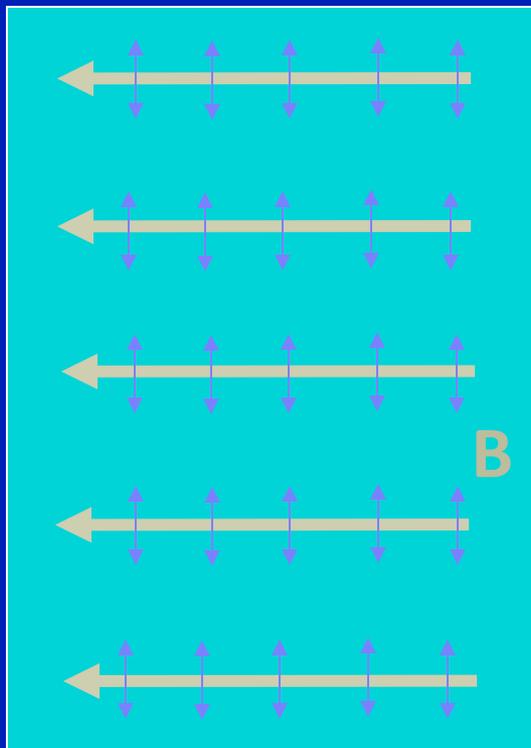


- **linear polarization** is perpendicular to the projection of **B** on the plane of the sky (normal to the wave vector)
- The maximal polarization is for the local emission from an ordered **B**-field: $P_{\max} = (\alpha+1)/(\alpha+5/3)$ where $F_\nu \propto \nu^{-\alpha}$, $-1/3 \leq \alpha \lesssim 1.5 \Rightarrow 50\% \leq P_{\max} \lesssim 80\%$
(Rybicki & Lightman 1979; Granot 2003)

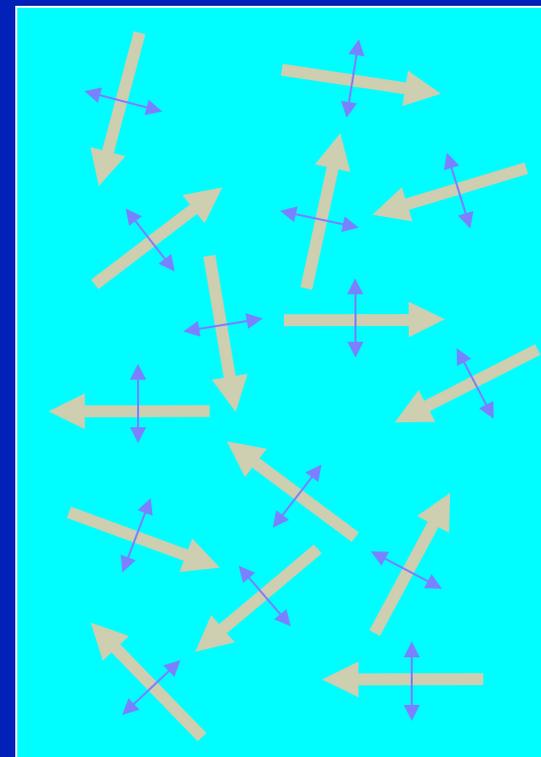
In the source rest frame:

- A uniform field produces $\mathbf{P} = \mathbf{P}_{\max}$
- For a field random when projected on the plane of the sky: $\mathbf{P} = \mathbf{0}$
- In particular, for a field isotropically tangled in 3D: $\mathbf{P} = \mathbf{0}$

Uniform B

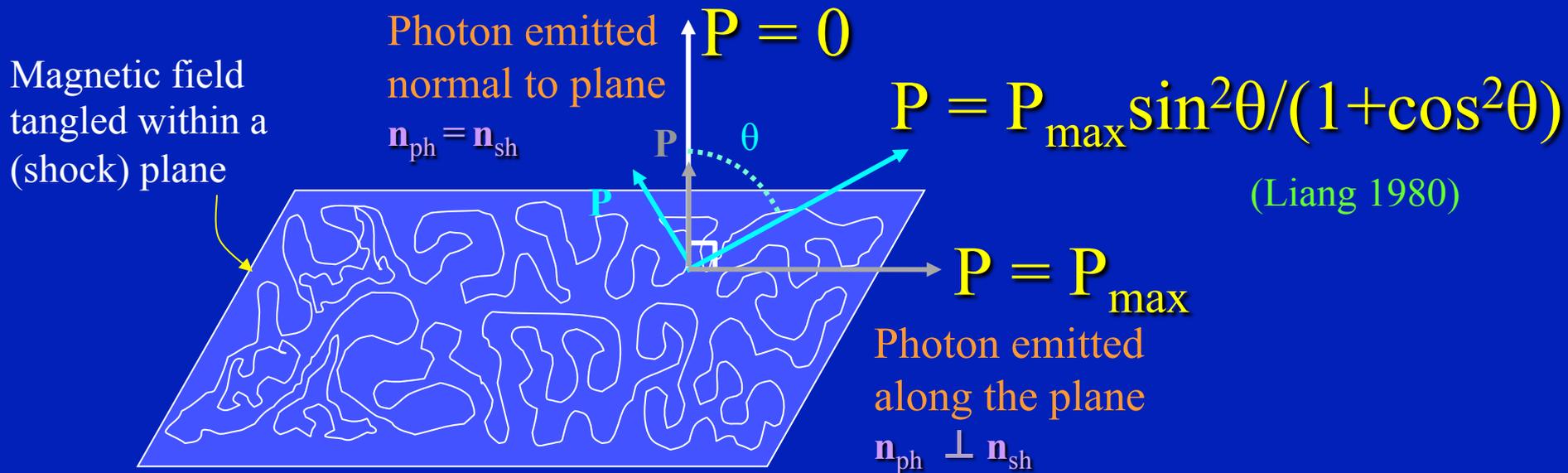


Random B



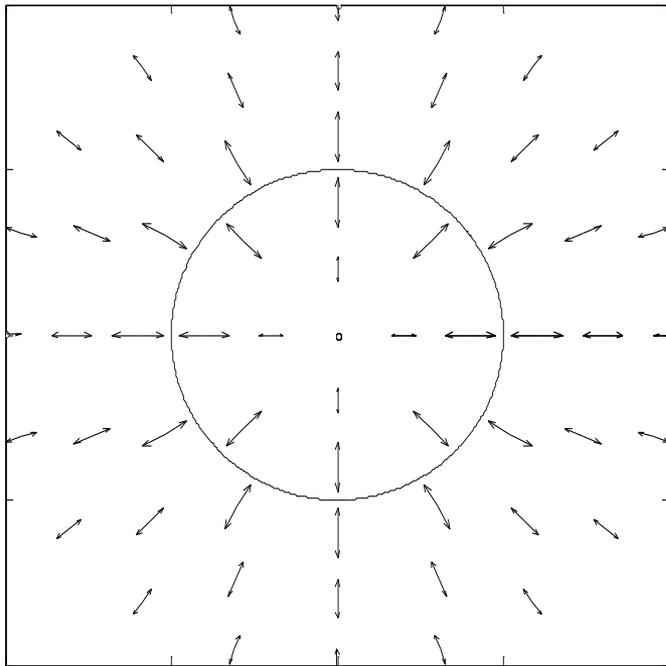
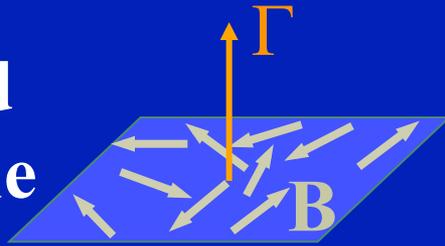
Shock Produced Magnetic Field:

- A magnetic field that is produced at a relativistic collisionless shock, due to the two-stream instability, is expected to be **tangled within the plane of the shock** (Medvedev & Loeb 1999)



Relativistic Source:

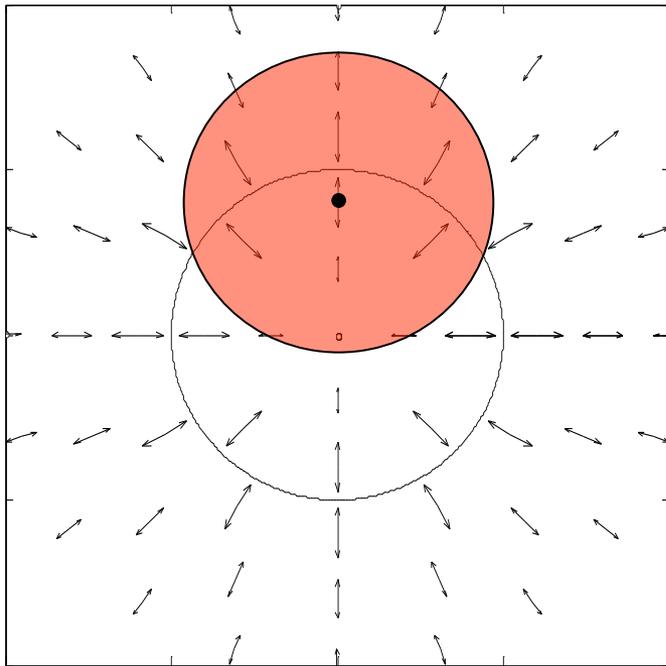
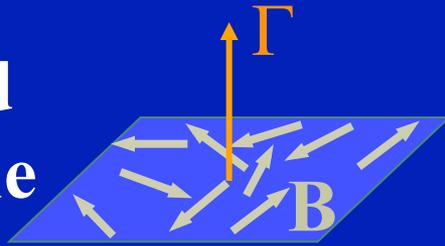
Random field
in shock plane



Sari 99; Ghisellini & Lazzati 99

Relativistic Source:

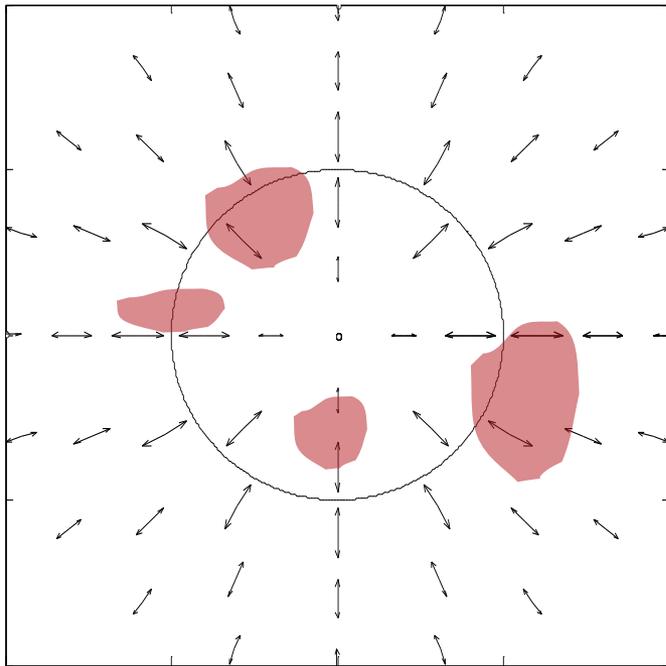
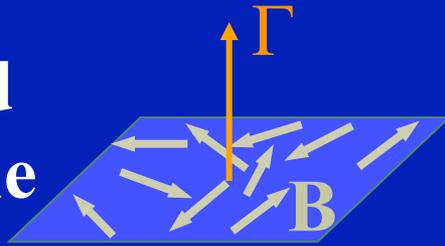
Random field
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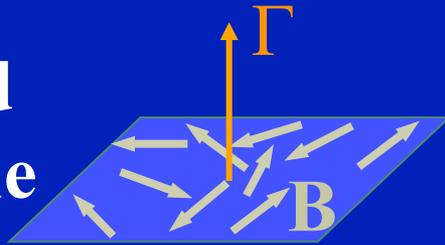
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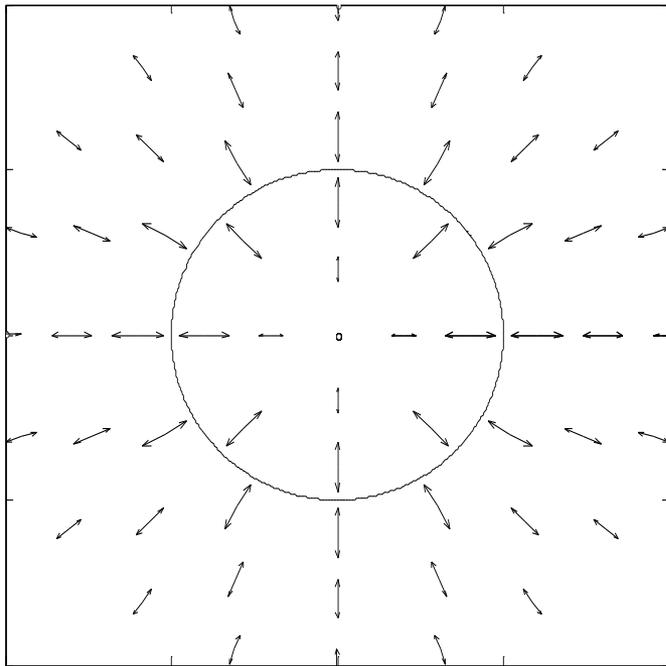
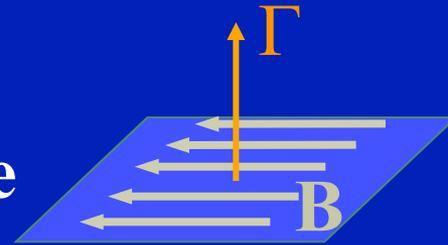
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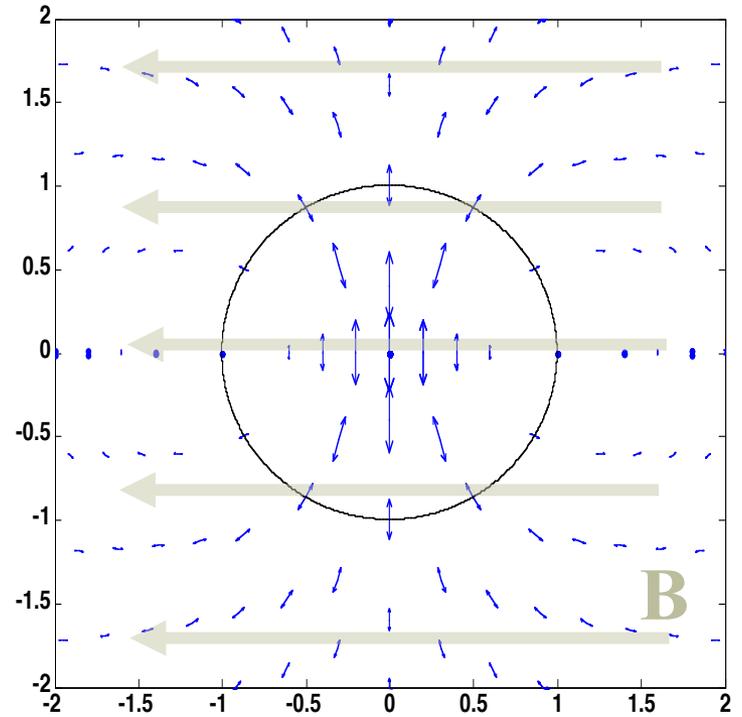
Random field
in shock plane



Ordered field
in shock plane



Sari 99; Ghisellini & Lazzati 99



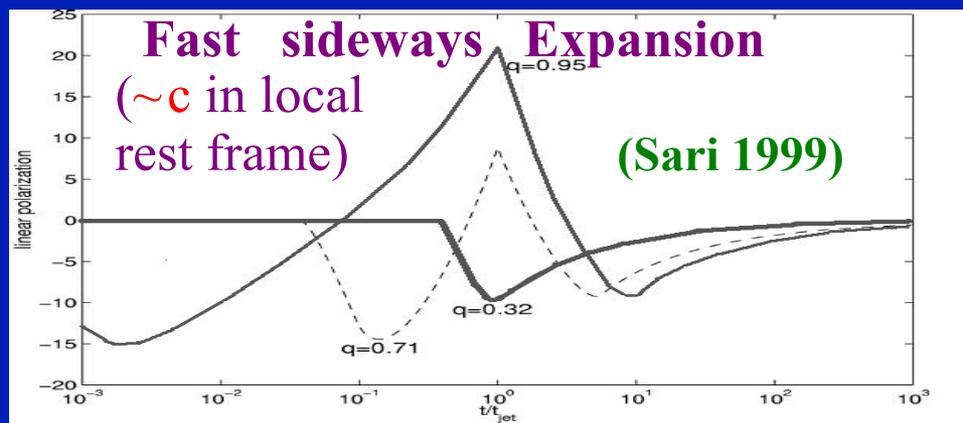
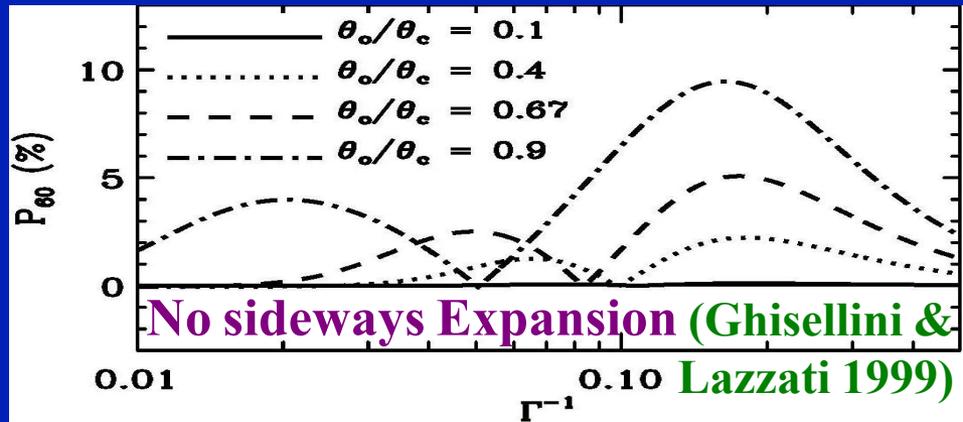
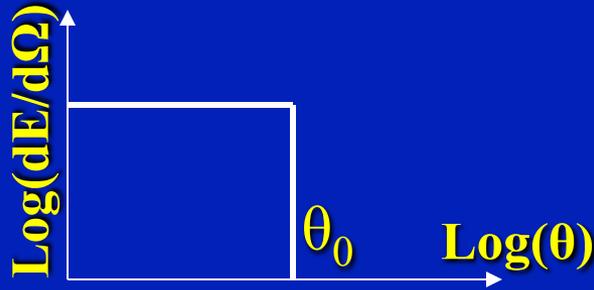
Granot & Königl 03

$$P \sim P_{\max}$$

Afterglow: Two “Traditional” Jet Structures

Uniform (top hat) jet:

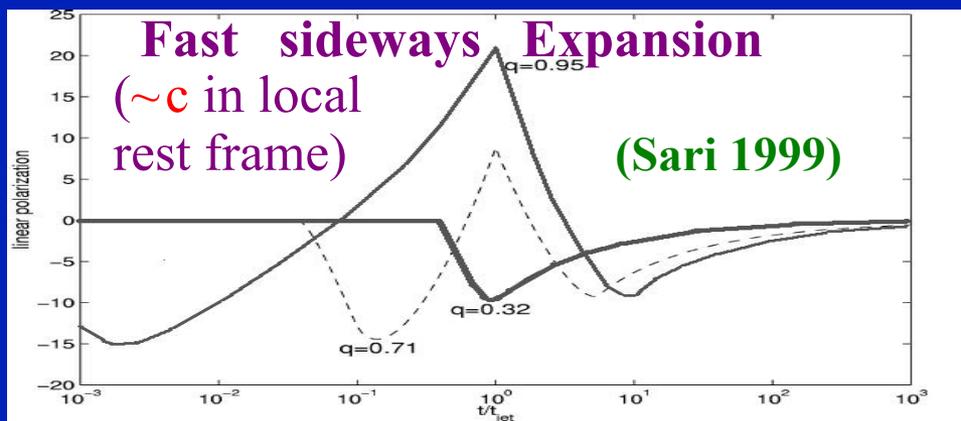
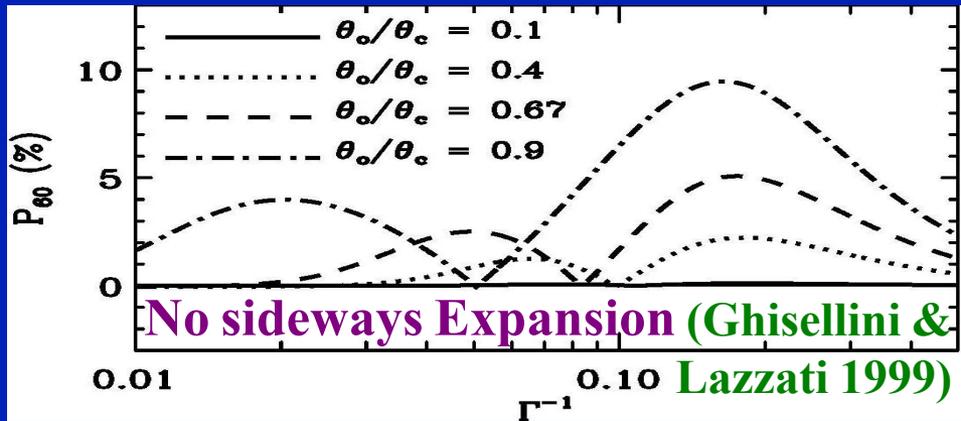
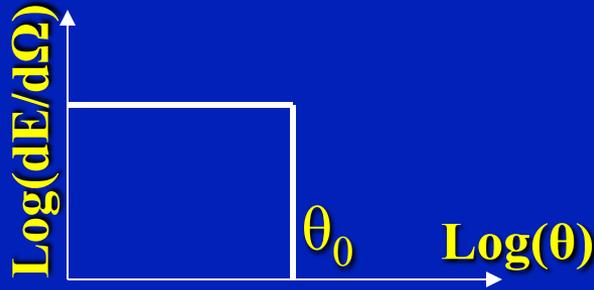
(Rhoads 97,99;
Sari+99, ...)



Afterglow: Two “Traditional” Jet Structures

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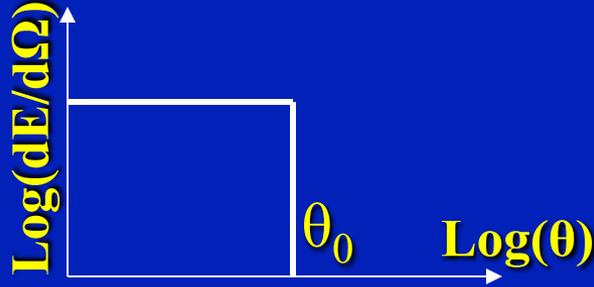
Main Prediction:
P vanishes & reappears
 with θ_p rotated by 90°
Is not clearly observed

Also: $P \lesssim 10\%-20\%$
 While $P_{obs} \sim 1-3\%$

Afterglow: Two “Traditional” Jet Structures

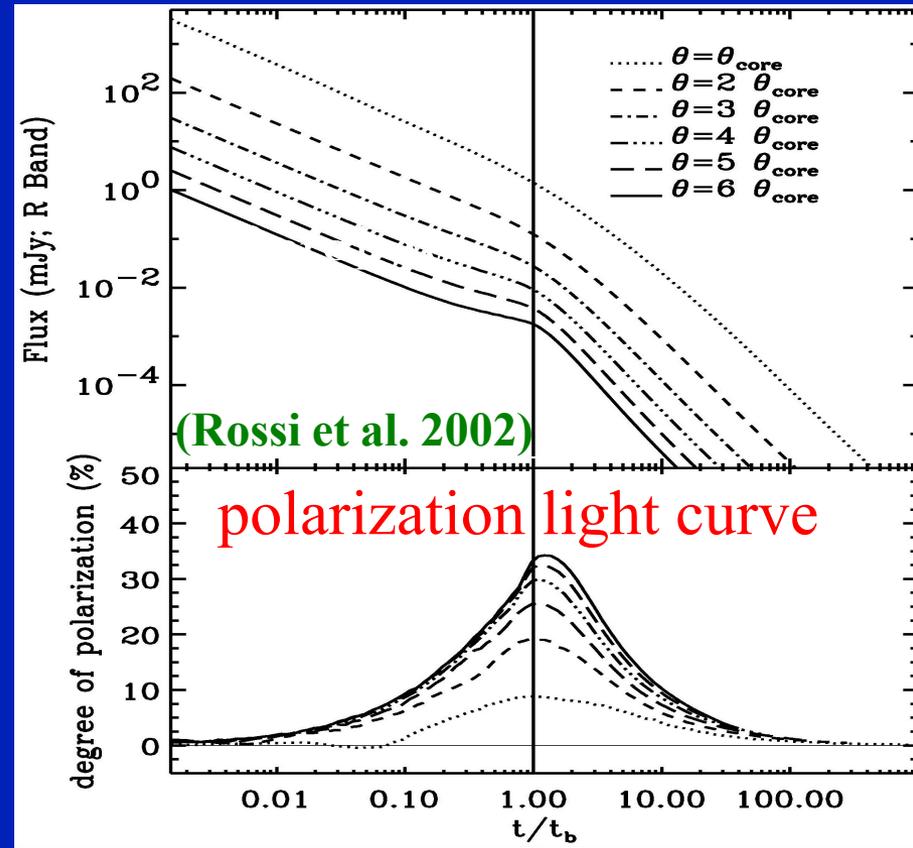
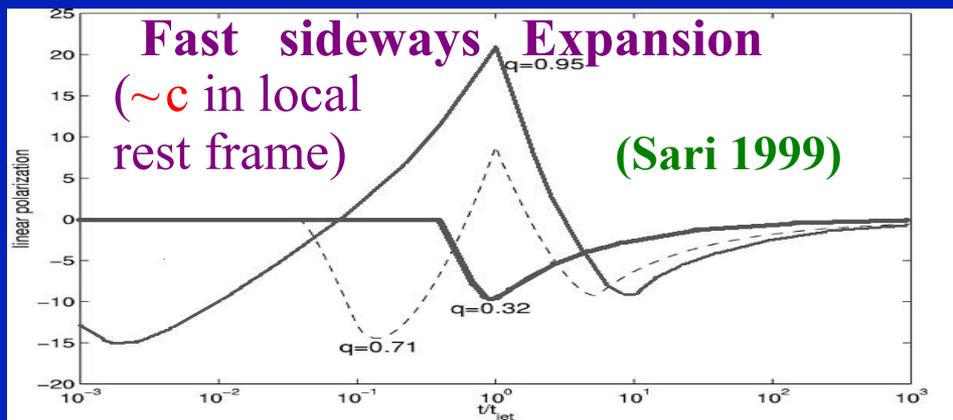
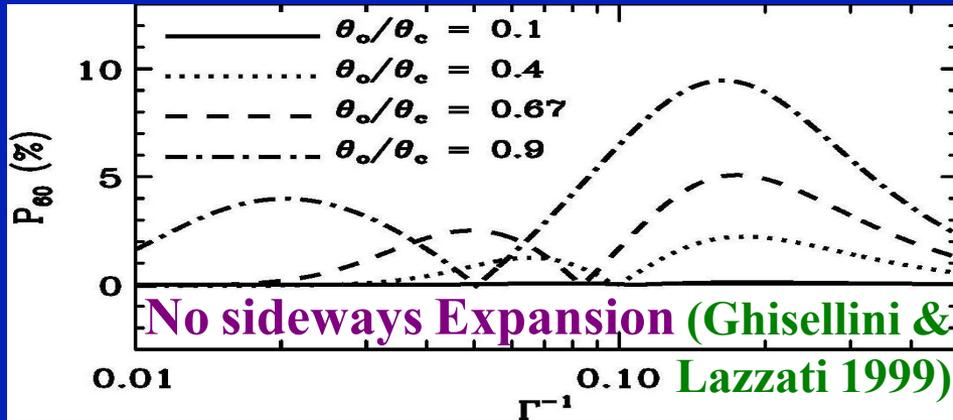
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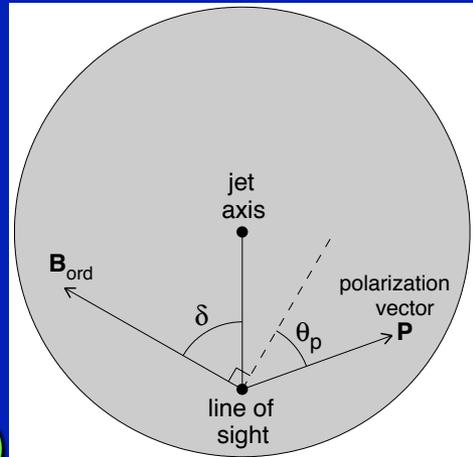
Structured jet:

(Postnov+01; Rossi+02;
Zhang & Meszaros 02)

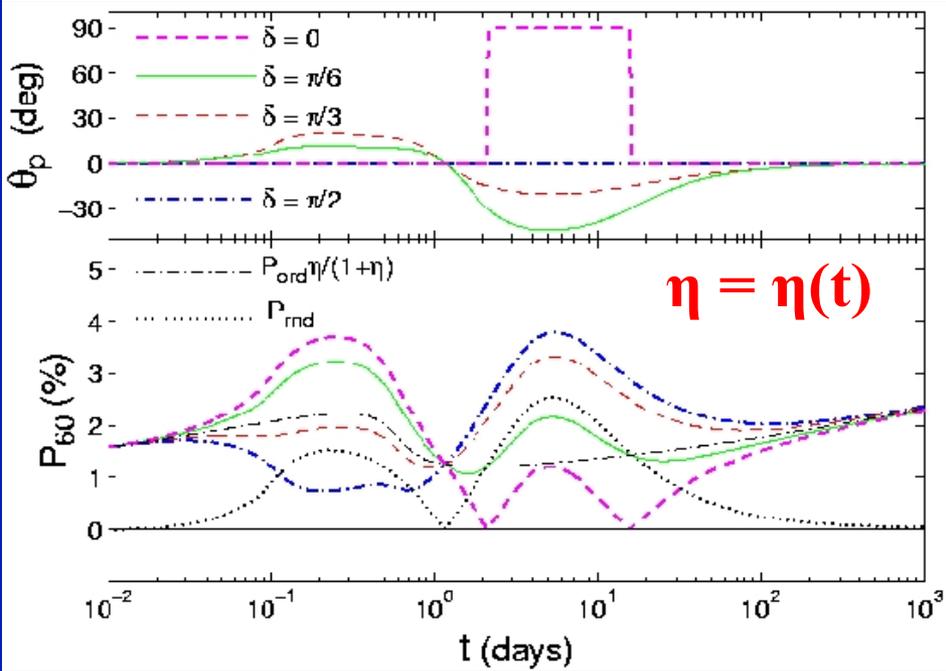
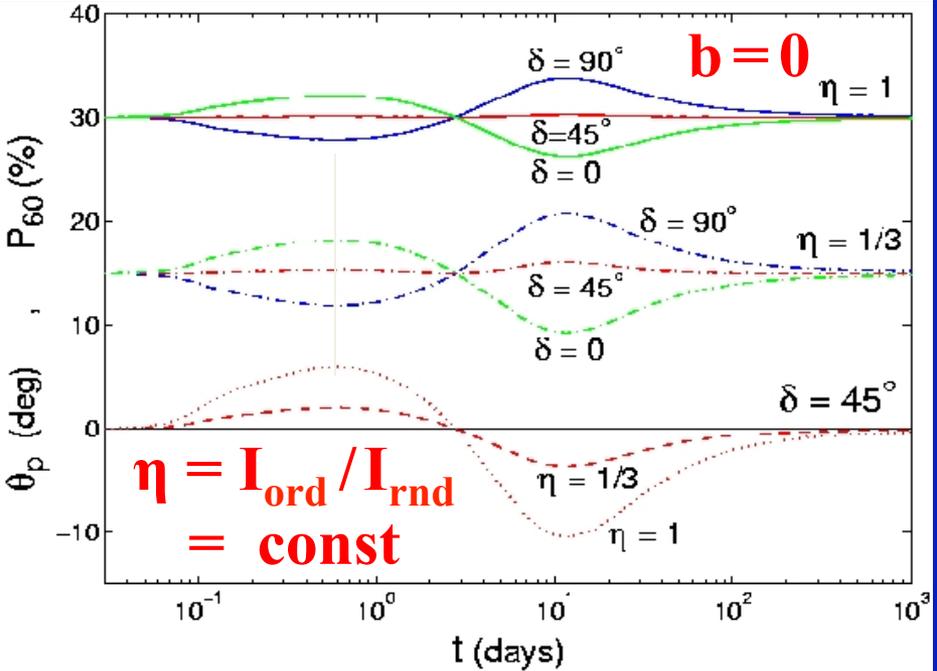


Combining Ordered \mathbf{B}_{ord} & Random \mathbf{B}_{rnd} Fields

- $P_{\text{ord}} \sim P_{\text{max}} \sim 60\%$ & $\theta_p = 90^\circ$ w.r.t. the direction of \mathbf{B}_{ord}
- In the afterglow $P \lesssim 3\% \Rightarrow I_{\text{ord}} \ll I_{\text{rnd}}$
but we can still have $I_{\text{ord}} P_{\text{ord}} \gtrsim I_{\text{rnd}} P_{\text{rnd}}$
- $\Rightarrow \mathbf{B}_{\text{rnd}}$ dominates I_{total}
but \mathbf{B}_{ord} dominates IP & P_{total}

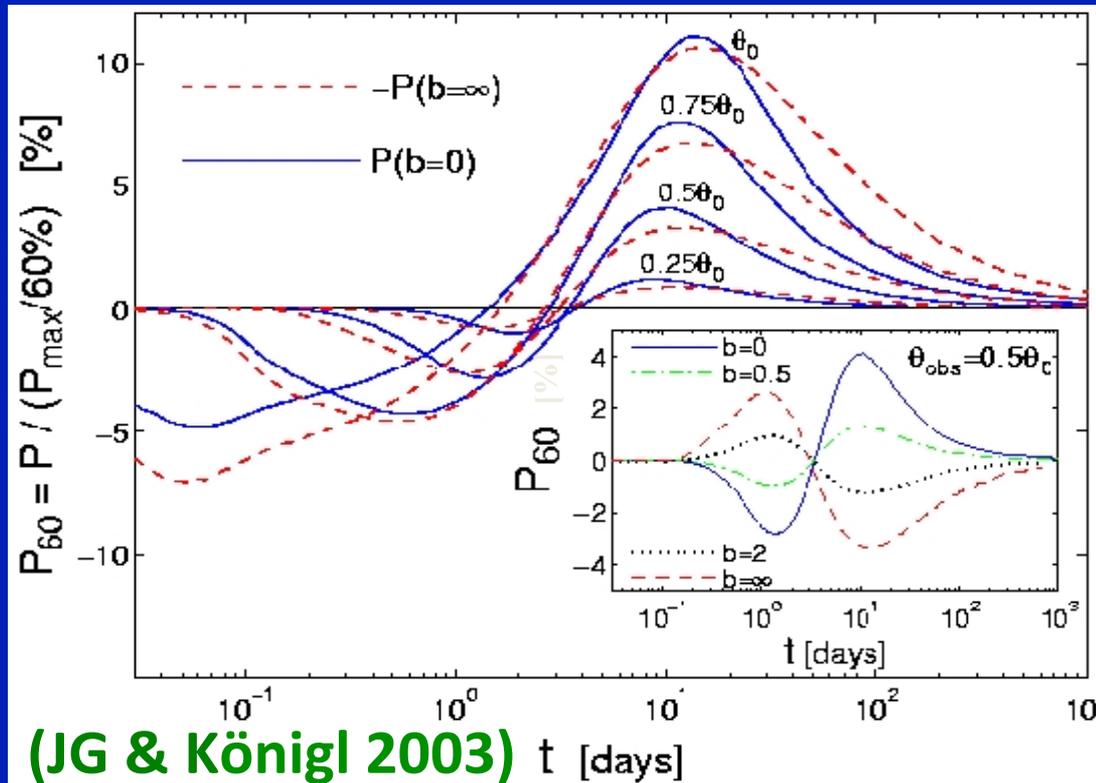


JG & Königl (2003)



The Random B-field's Degree of Anisotropy:

- $b = 2\langle B_{\parallel}^2 \rangle / \langle B_{\text{perp}}^2 \rangle$ parameterizes the asymmetry of \mathbf{B}_{rnd}
- $\text{Sign}(b-1)$ determines θ_p ($P > 0$ is along the direction from the line of sight to the jet axis & $P < 0$ is rotated by 90°)
- For $b \approx 1$ the polarization is very low (field is almost isotropic)
- $P \lesssim 3\%$ in afterglows observations $\Rightarrow 0.5 \lesssim b \lesssim 2$



$$P = P_{\text{max}} / [1 + 2 / (b-1) \sin^2 \theta']$$

(valid for $j'_v \propto [B' \sin \chi']^2$)

$$\theta_0 = 5^\circ$$

$$E_{\text{jet}} = 3 \times 10^{51} \text{ erg}$$

$$n = 1 \text{ cm}^{-3}$$

$$z = 1$$

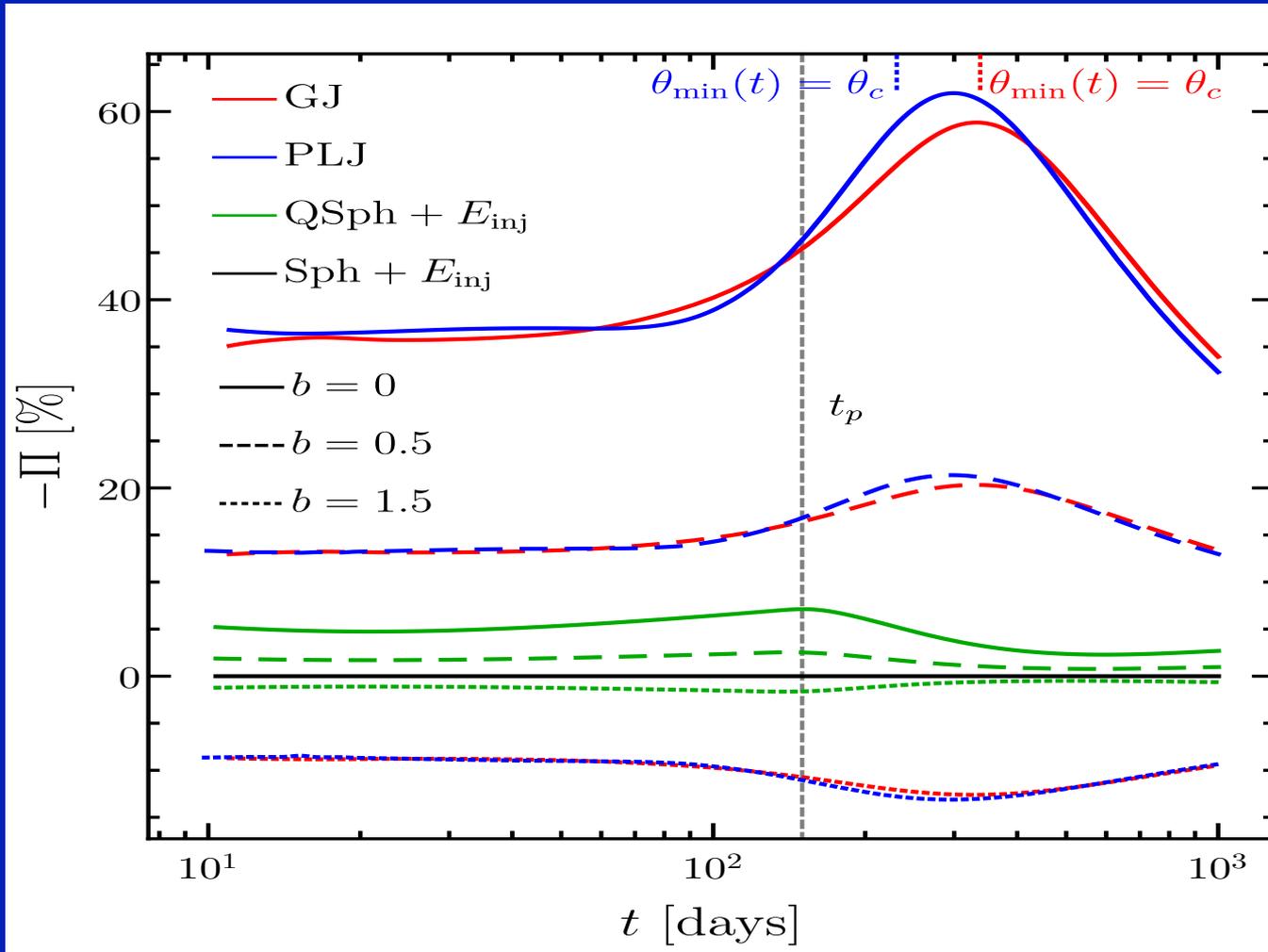
$$p = 2.5$$

$$\epsilon_e = 0.1$$

$$\epsilon_B = 0.01$$

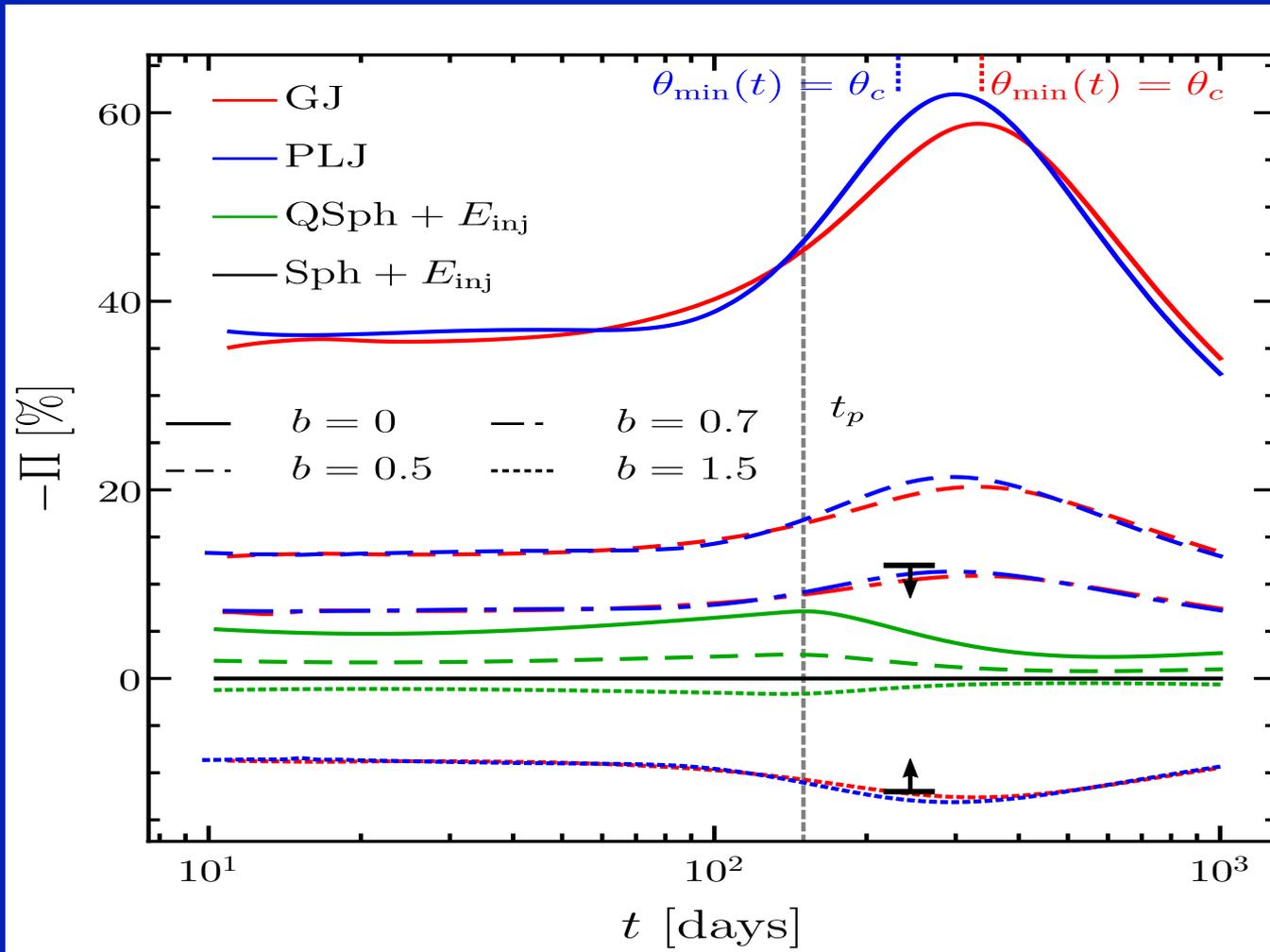
GW170817/GRB170817A Afterglow (Gill & JG 18)

- Assuming a shock-produced B-field with $b \equiv 2\langle B_{\parallel}^2 \rangle / \langle B_{\perp}^2 \rangle$
- Data favor two core-dominated jet models with similar $P(t)$



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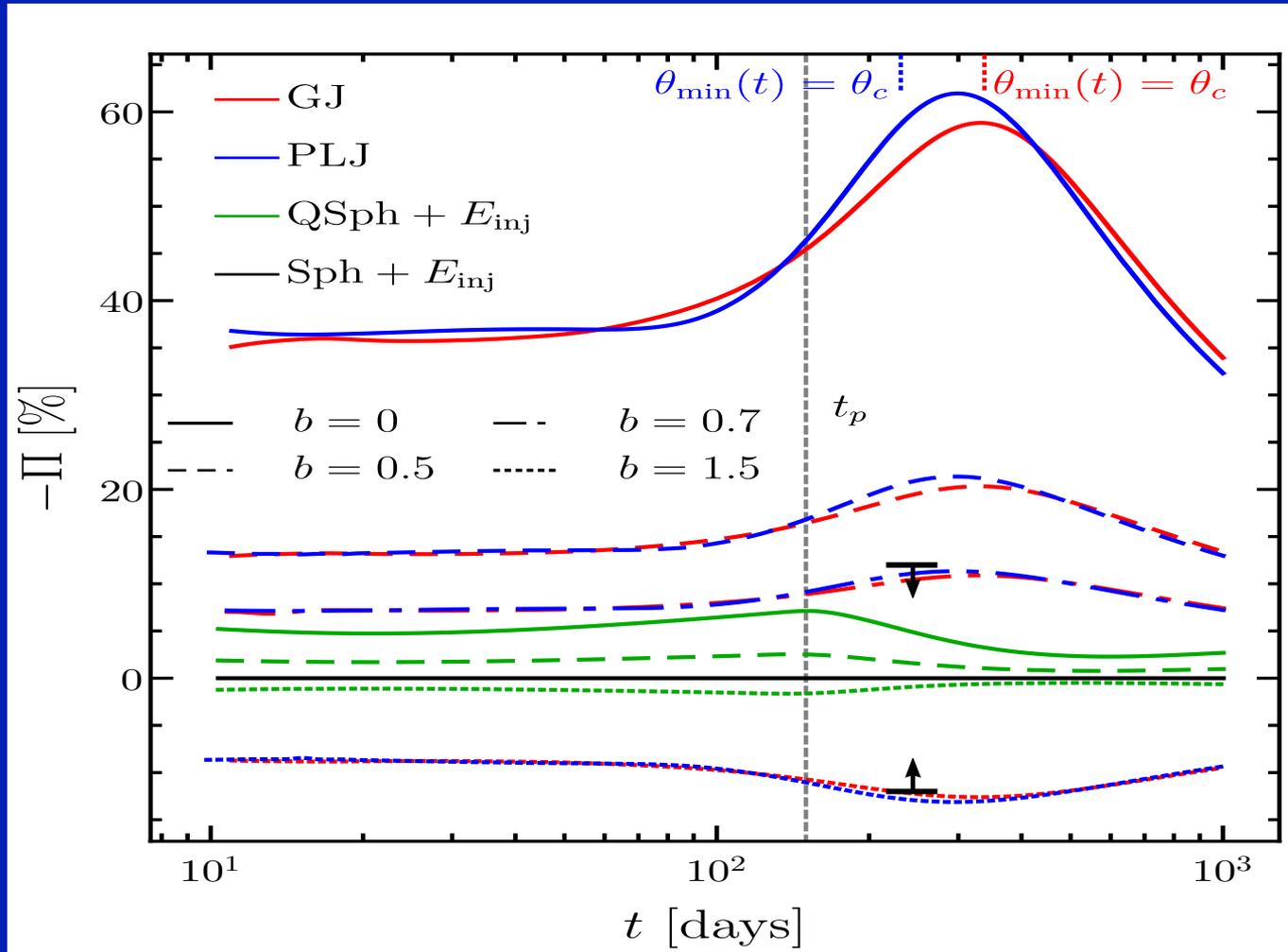
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New: upper limit
 $P_{lin} < 12\%$ @
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(Corsi + 2018)

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$0.7 \lesssim b \lesssim 1.5$
for jet models



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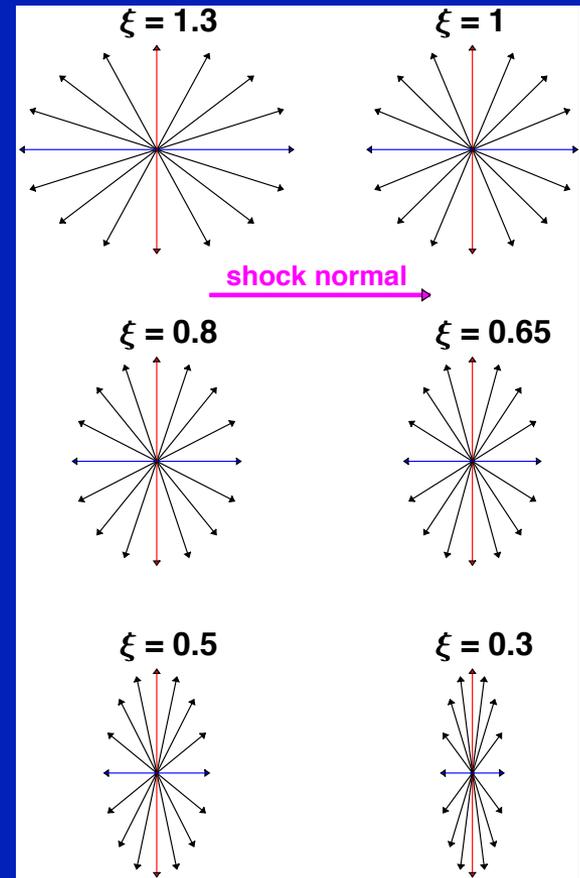
More realistic assumptions \Rightarrow B-field in collisionless shocks:

- 2D emitting shell \rightarrow 3D emitting volume (local BM76 radial profile)

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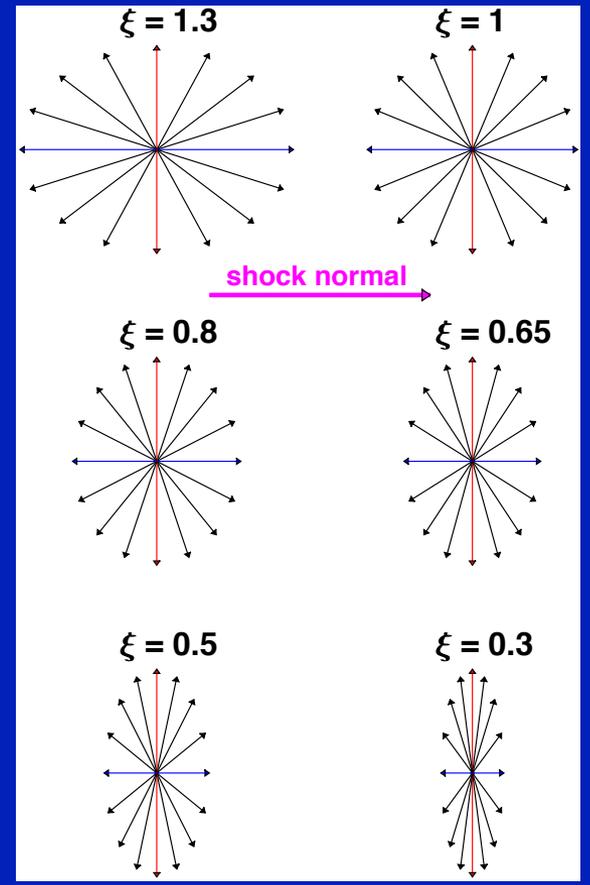
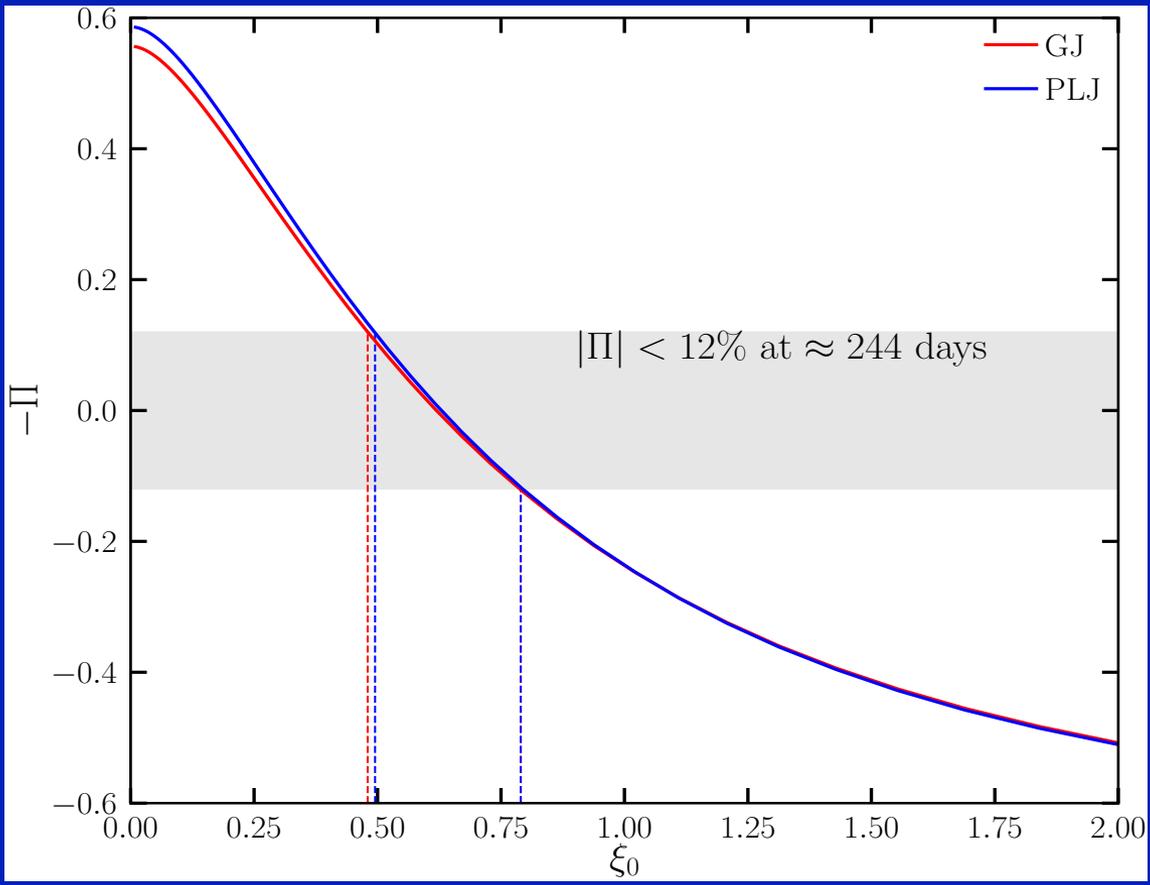
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- B-field evolution by faster radial expansion: $L'_r/L'_{\theta,\phi} \propto \chi^{(7-2k)/(8-2k)}$
B-field isotropic in 3D with $B'_r \rightarrow \xi B'_r$ (Sari 1999); $\xi = \xi_0 \chi^{(7-2k)/(8-2k)}$



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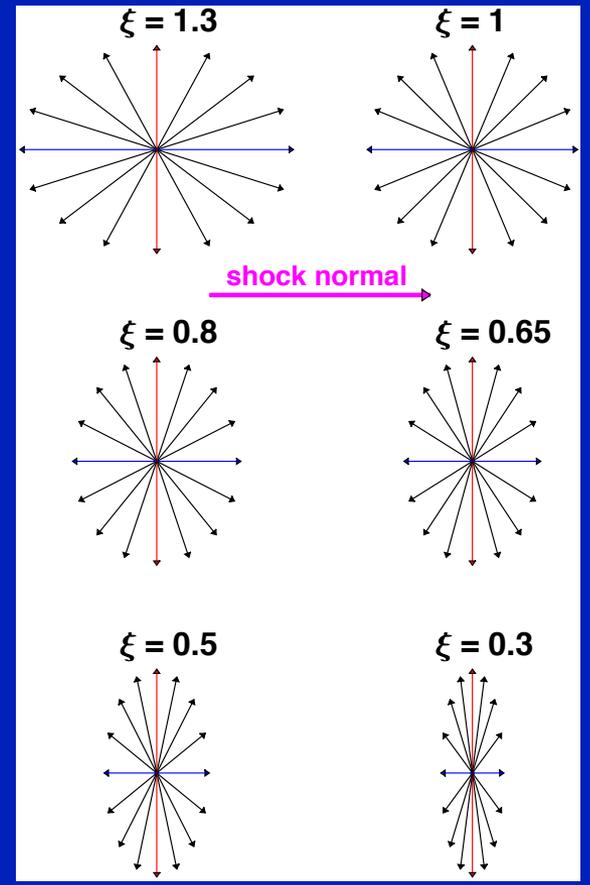
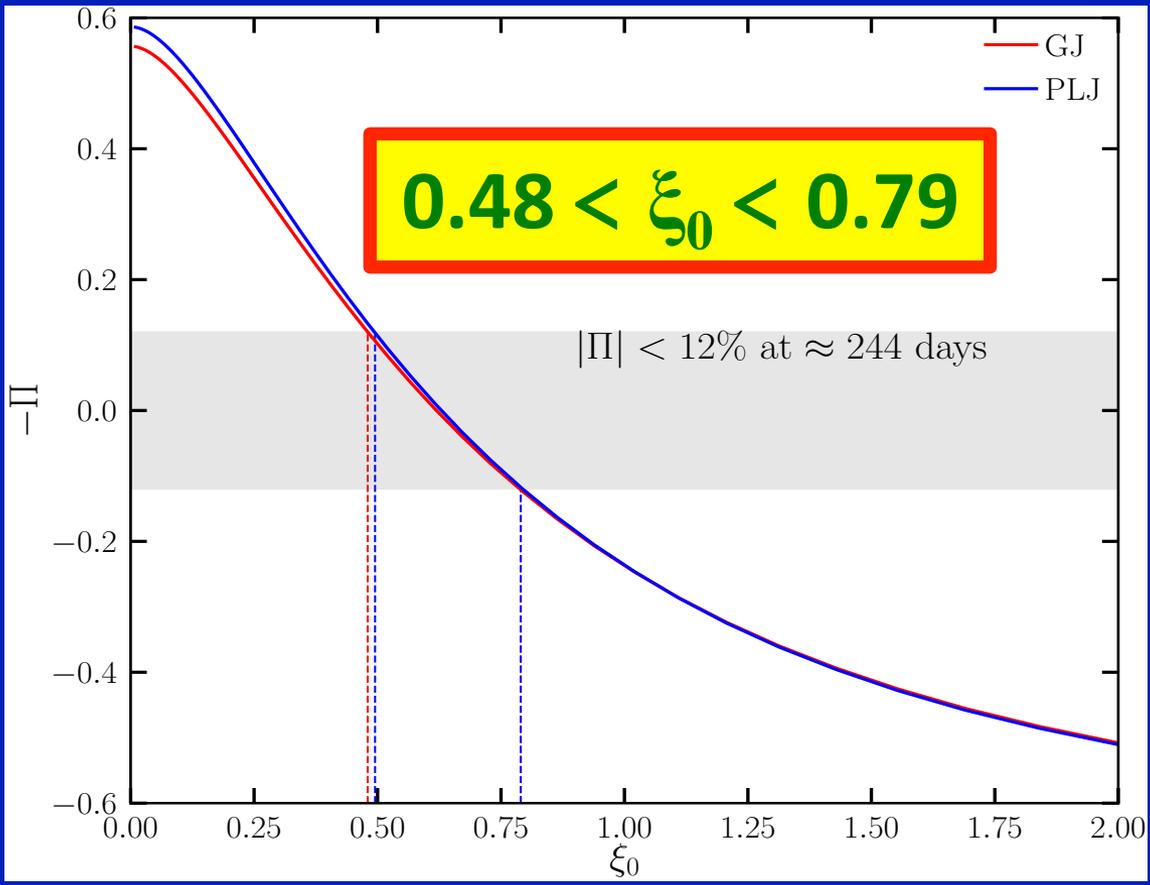
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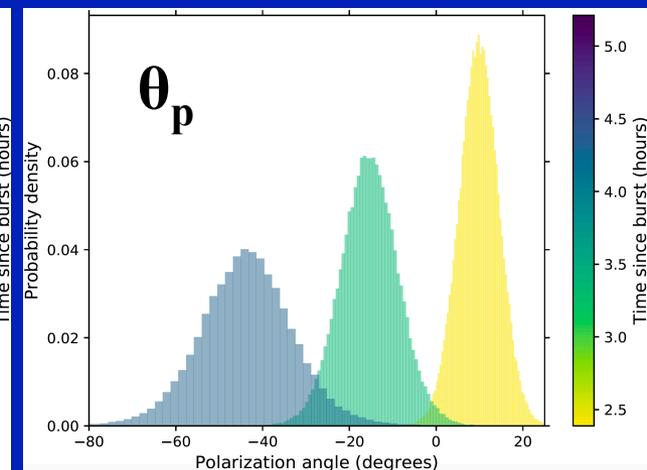
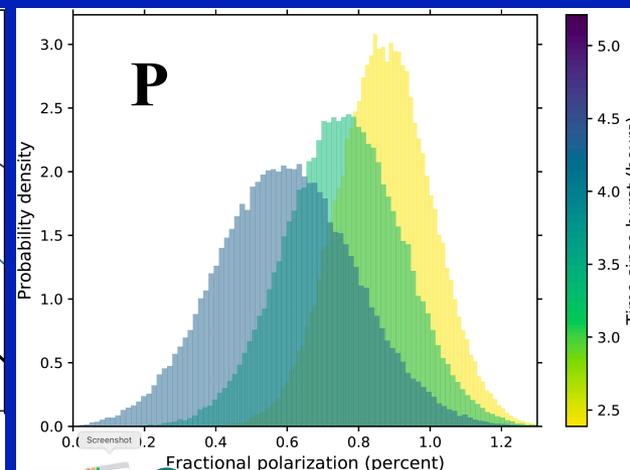
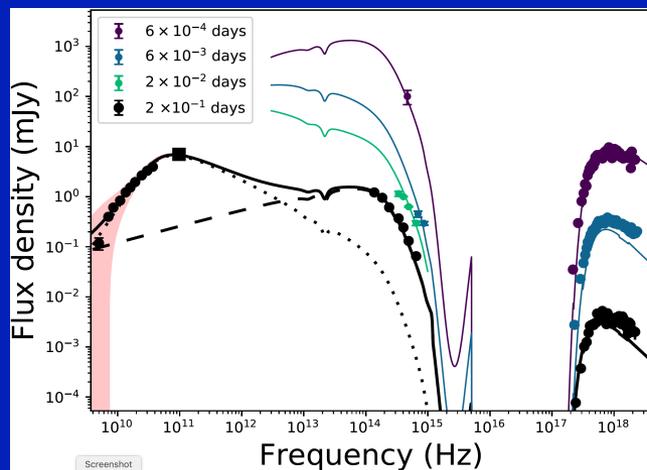
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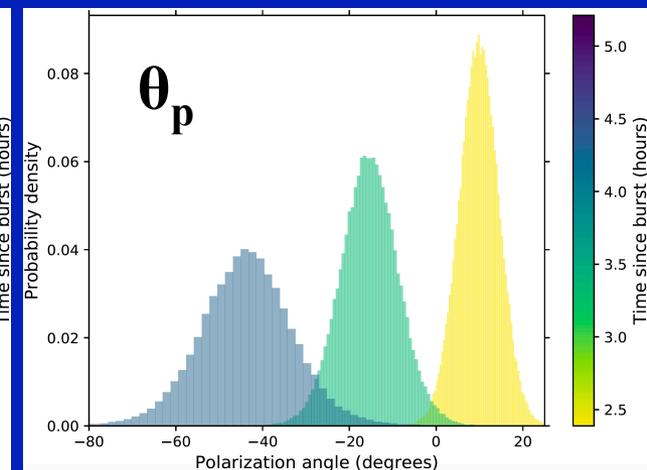
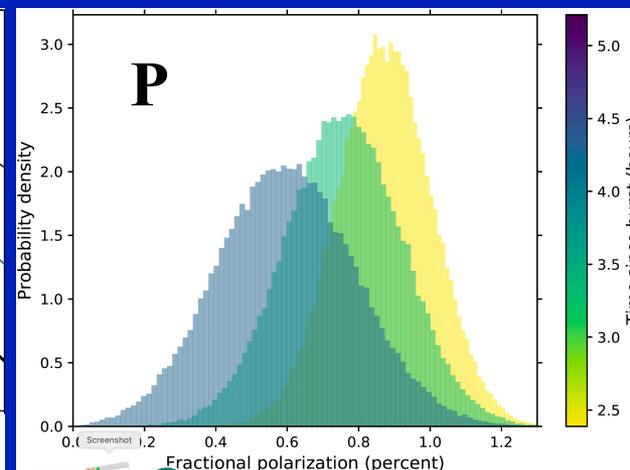
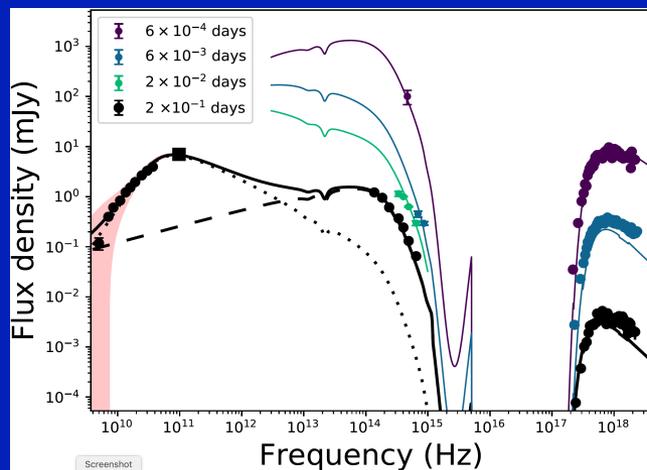
Reverse shock Pol.: Ejecta B-field (Laskar + 2019)

- ALMA observed GRB190114C reverse shock @ 97.5 GHz:
 $P \approx 0.9 \rightarrow 0.6\%$, $\Delta\theta_p \approx 54^\circ$ (2.2 \rightarrow 5.2 hr); 1st GRB radio pol.



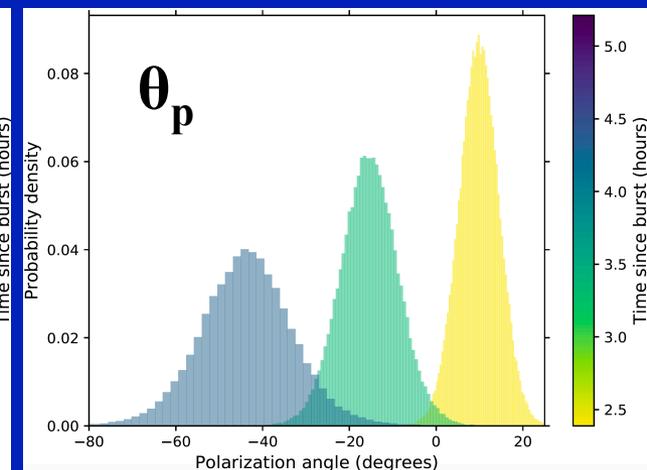
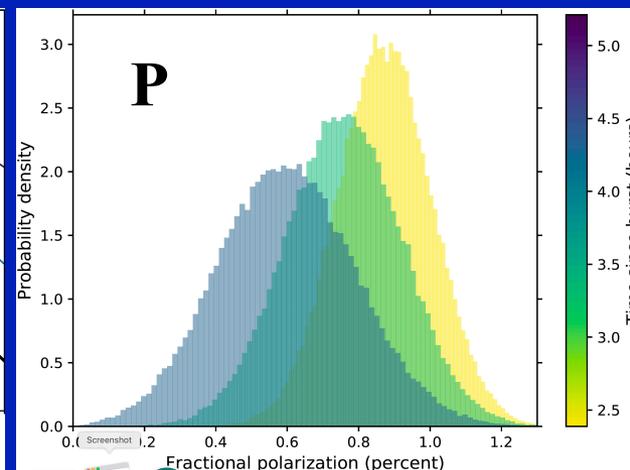
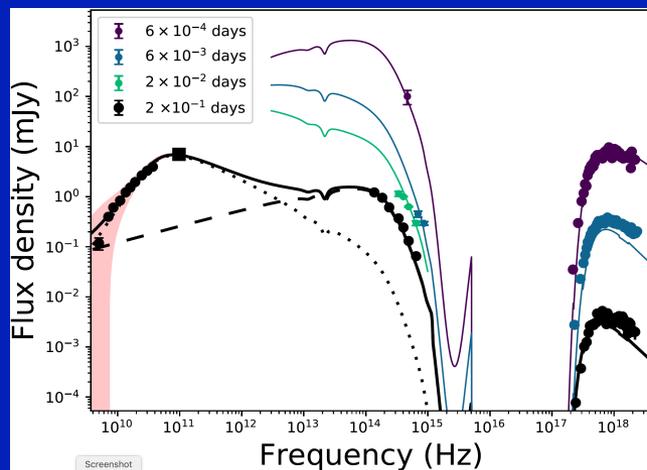
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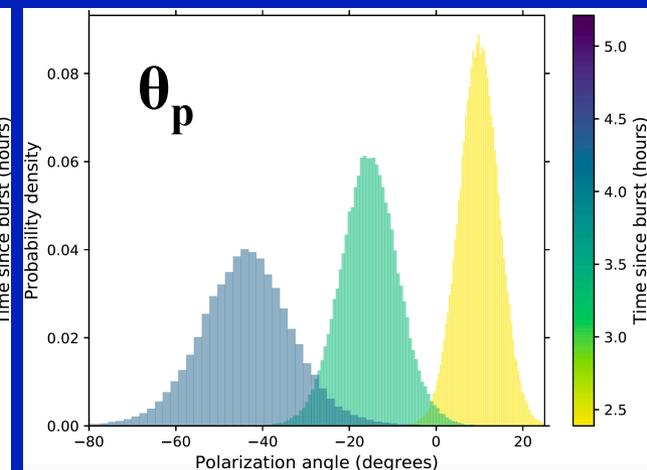
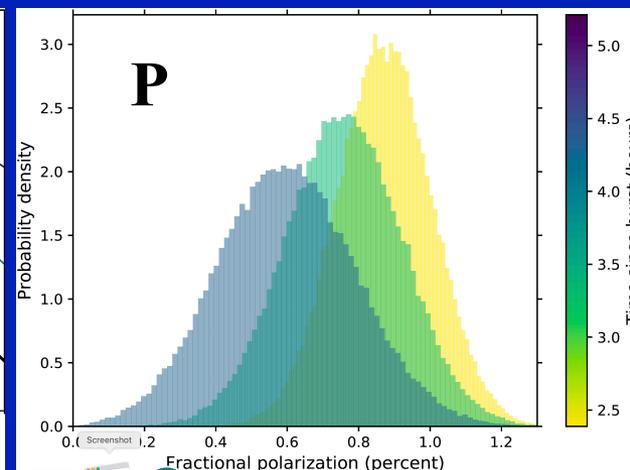
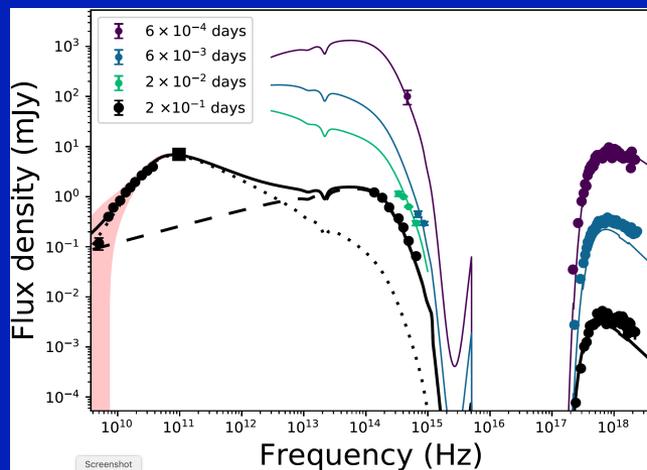
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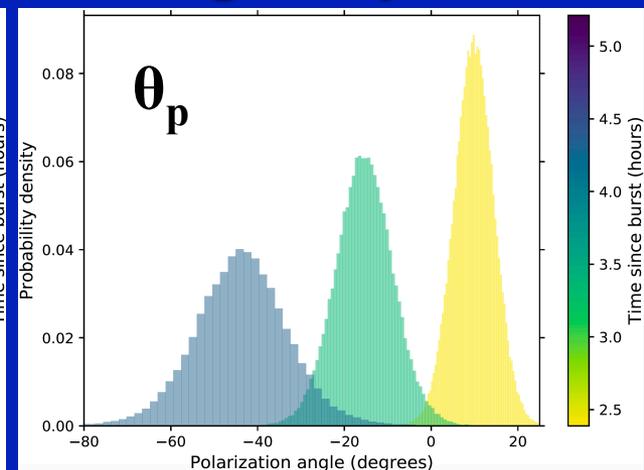
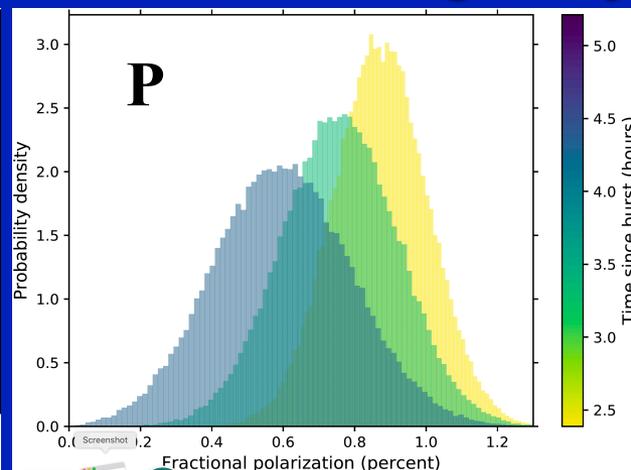
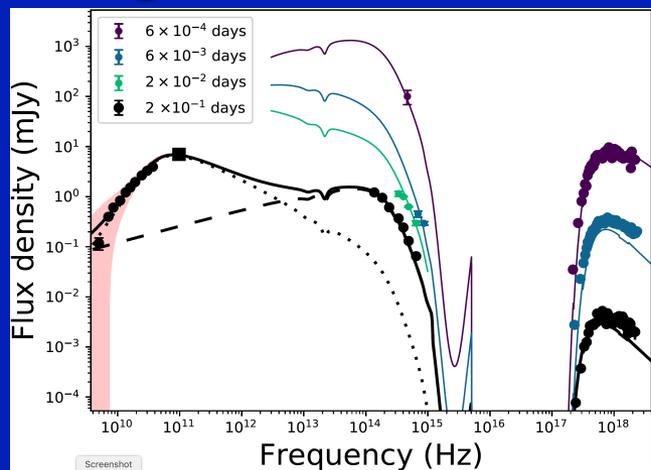
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- $N \sim (\Gamma_{\text{ej}} \theta_B)^{-2}$ incoherent patches: $\Gamma_{\text{ej}} \approx 15$, $P \sim P_{\text{max}}/N^{1/2} \Rightarrow \theta_B \sim P/P_{\text{max}} \Gamma_{\text{ej}} \sim 10^{-3}$ & $\Delta\theta_p \sim 1$ expected over $\Delta t \sim t$ ✓



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- $\Delta\theta_p \approx 54^\circ$ rules out an axi-symmetric configuration (e.g. a global toroidal B-field in the original jet; **A patchy shell?**)



Prompt γ -ray Polarization: hard to measure

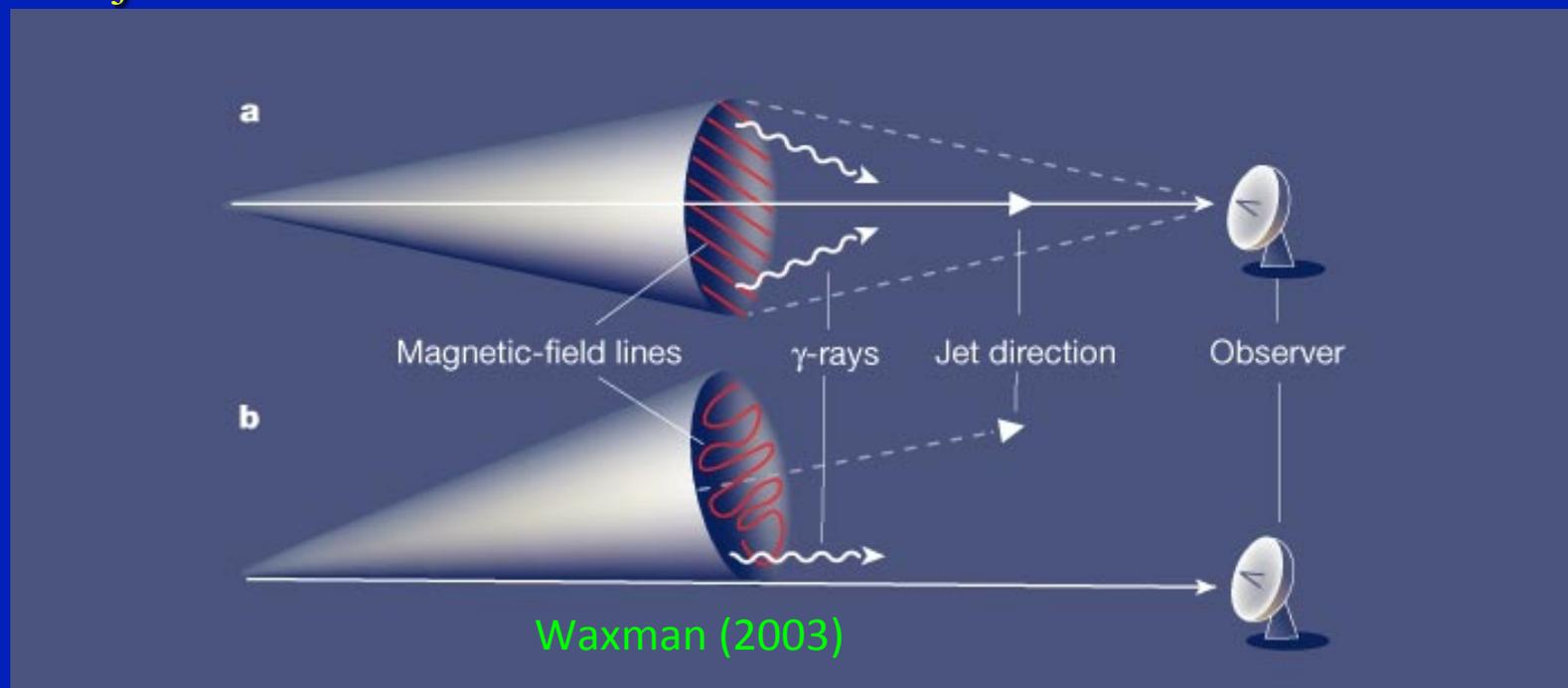
First consider synchrotron emission:

■ Shock produced B-field + $\theta_{\text{obs}} \lesssim \theta_j - 1/\Gamma \Rightarrow P \approx 0$

■ $P \sim P_{\text{max}}$ can be achieved in the following ways:

(1) **ordered magnetic field** in the ejecta,

(2) **special geometry**: $|\theta_{\text{obs}} - \theta_j| \lesssim 1/\Gamma \Rightarrow$ favors narrow jets: $\theta_j \lesssim 1/\Gamma$ (works with a shock produced B-field)



Narrow Jet + shock produced B-field

- High polarization + reasonable flux $\Rightarrow \theta_j < \theta_{\text{obs}} \lesssim \theta_j + 1/\Gamma$
- A reasonable probability for such $\theta_{\text{obs}} \Rightarrow \Gamma\theta_j \lesssim \text{a few}$
- Since $\Gamma \gtrsim 100$ & $\theta_j \gtrsim 0.05$, $\Gamma\theta_j \gtrsim 5$ and is typically larger

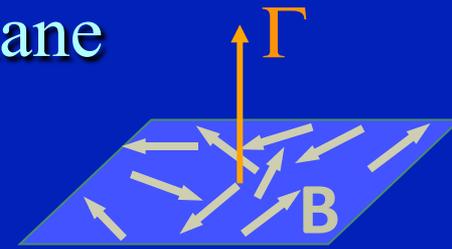
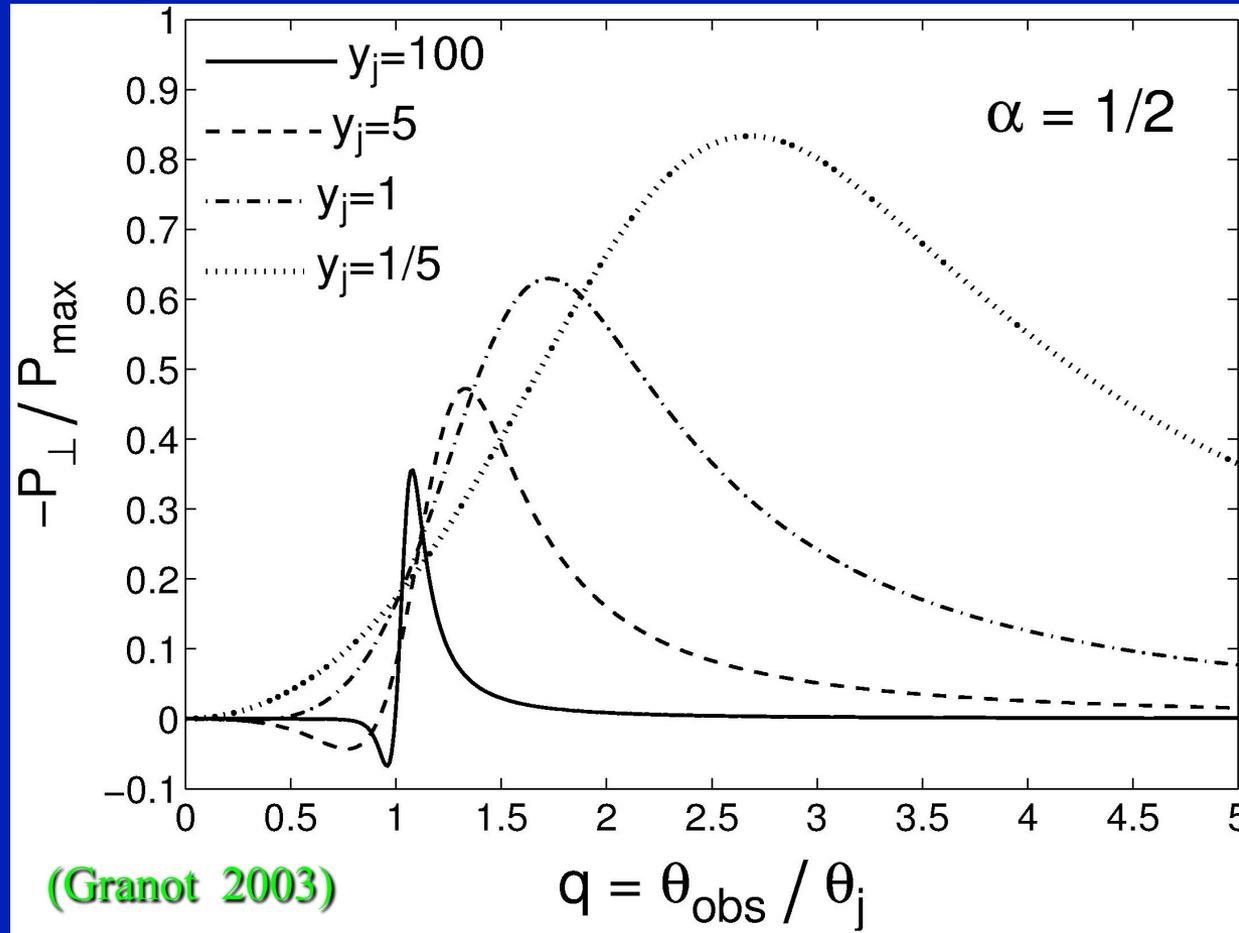
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- A reasonable probability for such $\theta_{\text{obs}} \Rightarrow \Gamma\theta_j \lesssim \text{a few}$
- Since $\Gamma \gtrsim 100$ & $\theta_j \gtrsim 0.05$, $\Gamma\theta_j \gtrsim 5$ and is typically larger
- The jet must have **sharp edges**: $\Delta\theta_j \lesssim 1/4\Gamma$ (Nakar et al. 03)
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- Afterglow obs. imply more random B_{rnd} : $0.48 < \xi_0 < 0.79$

Adding pulses: Random B-field in shock plane



$$y_j = (\Gamma \theta_j)^2$$

$$F_{\nu} \propto \nu^{-\alpha}$$

- $\Delta\Gamma \sim \Gamma$ between different shell collisions (different pulses in GRB light curve) reduces \mathbf{P} by a factor ~ 2

Prompt γ -ray Polarization: short summary

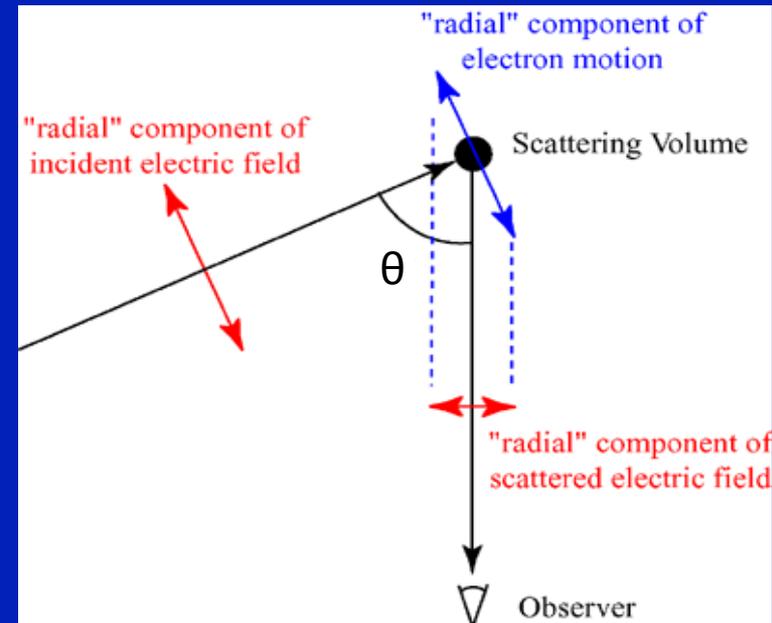
	Ordered Field	Sharp-edge Jet
$P \sim 80\%$	X	X
$P \sim 50\%$	✓	X
$P \sim 25\%$	with $B_{\text{rnd}} \lesssim B_{\text{ord}}$	✓
$P \lesssim 10\%$	$B_{\text{rnd}} > B_{\text{ord}}$	with $B_{\text{rnd}} \gtrsim B_{\text{ord}}$
statistics	High P in all GRBs	low P in most GRBs
Potential problems	Some B_{rnd} required for Fermi acceleration	$\Gamma\theta_j \lesssim$ a few, $\Delta\Gamma \sim \Gamma$, B_{rnd} ($0.48 < \xi_0 < 0.79$)

Alternative to Synchrotron: Compton Drag

(Bulk Inverse Compton Scattering of External photons)

(Lazzati et al. 2003; Dar & De Rujula 2003, Eichler & Levinson 2003)

- Requires special geometry/viewing angle, $\theta_j < \theta_{\text{obs}} \lesssim \theta_j + 1/\Gamma$
- Polarization properties similar to synchrotron + B_{rnd} with an advantage: **local polarization** $P = (1 - \cos^2\theta)/(1 + \cos^2\theta)$ can reach up to **100%** while $P_{\text{max}} \sim 70\%$ for synchrotron
- Shares drawbacks of shock produced field + narrow jet

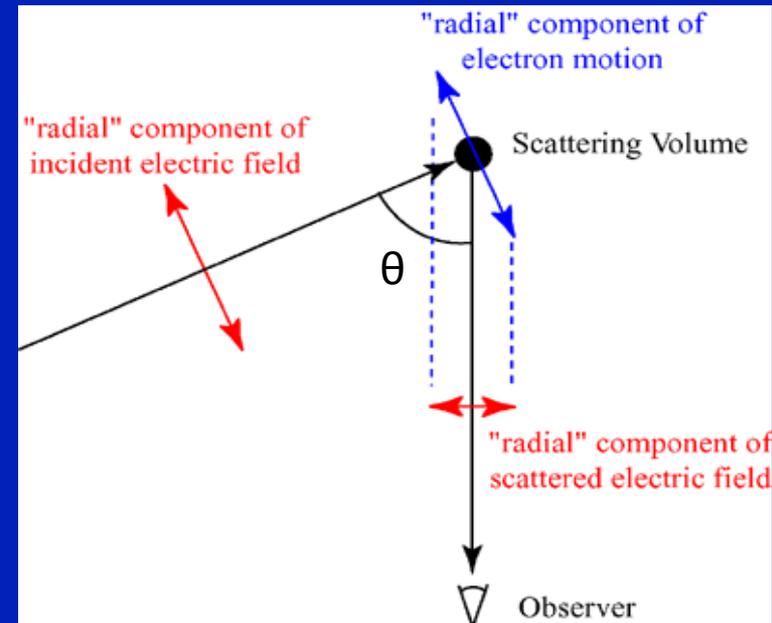


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- Shares drawbacks of shock produced field + narrow jet
- It has additional problems, unrelated to polarization:
 - ◆ Explaining prompt GRB spectrum
 - ◆ Supplying external photons for all the ejected shells
 - ◆ High photon density \Rightarrow small radii \Rightarrow high $\tau_{\gamma\gamma}$



Alternative to Synchrotron: Photospheric Emission

(Comptonized radiation advected from optically thick to thin region of the jet)

(Beloborodov 11; Thompson & Gill 14; Lundman+14; Vurm & Beloborodov 16; Lundman +16)

- Need to integrate radiation transfer equations for the Stokes parameters $I(r,\mu)$ & $Q(r,\mu)$ from $\tau_T \gg 1$ to $\tau_T \ll 1$.
- $P=0$ seed photons become anisotropic at $\tau_T \lesssim 10 \Rightarrow P \approx 0.45P_{\text{Compton-drag}}$

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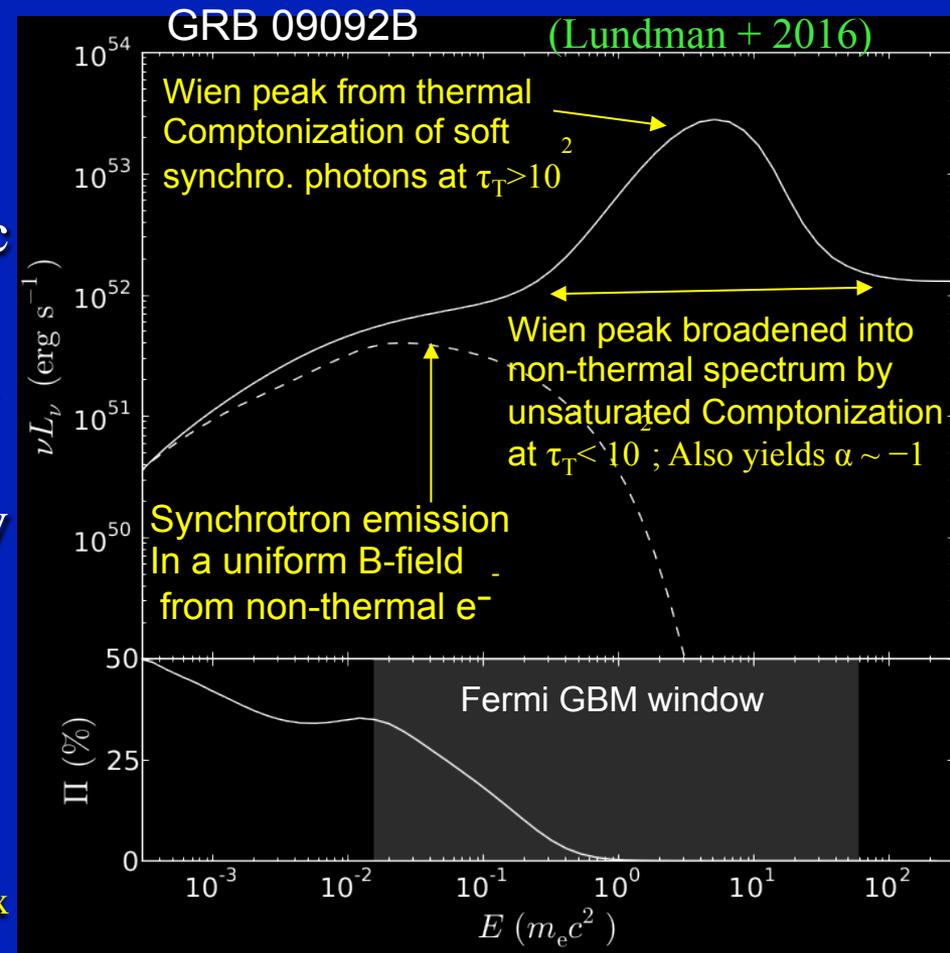
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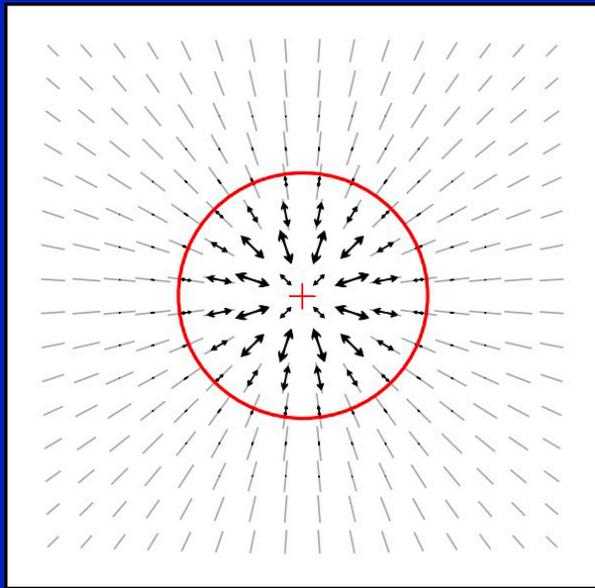
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- Synchrotron + B_{ord} (spherical flow): Unscattered syn. photons emitted at $\tau_T \sim 1$ dominate at $E \ll E_{\text{pk}} \Rightarrow P \sim P_{\text{syn,max}}$

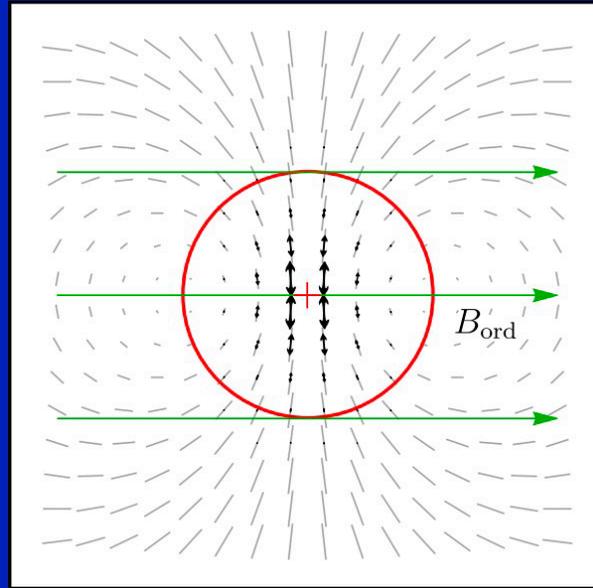


Prompt GRB Polarization (Gill, JG & Kumar 2018):

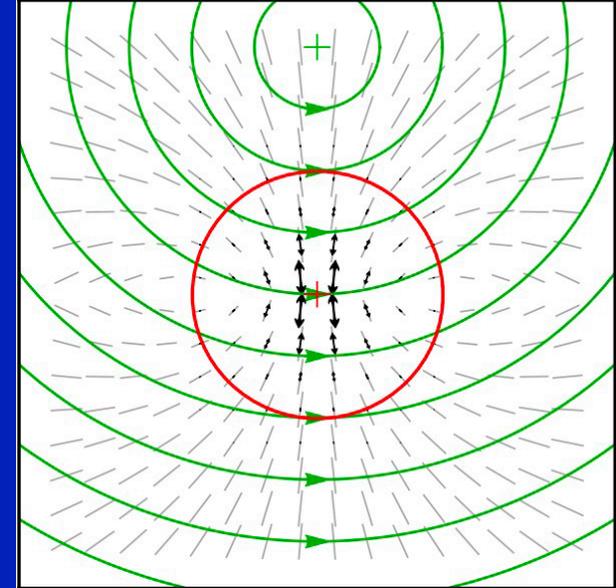
- Comprehensive study in view of γ -ray polarimetry missions
- **Jet structure:** top hat (sharp/smooth), Gaussian, core+power-law
- **Emission mechanism:** **synchrotron**, photospheric, Compton drag
- Time resolved, integrated over single or multiple pulses



Random B-field in 2D



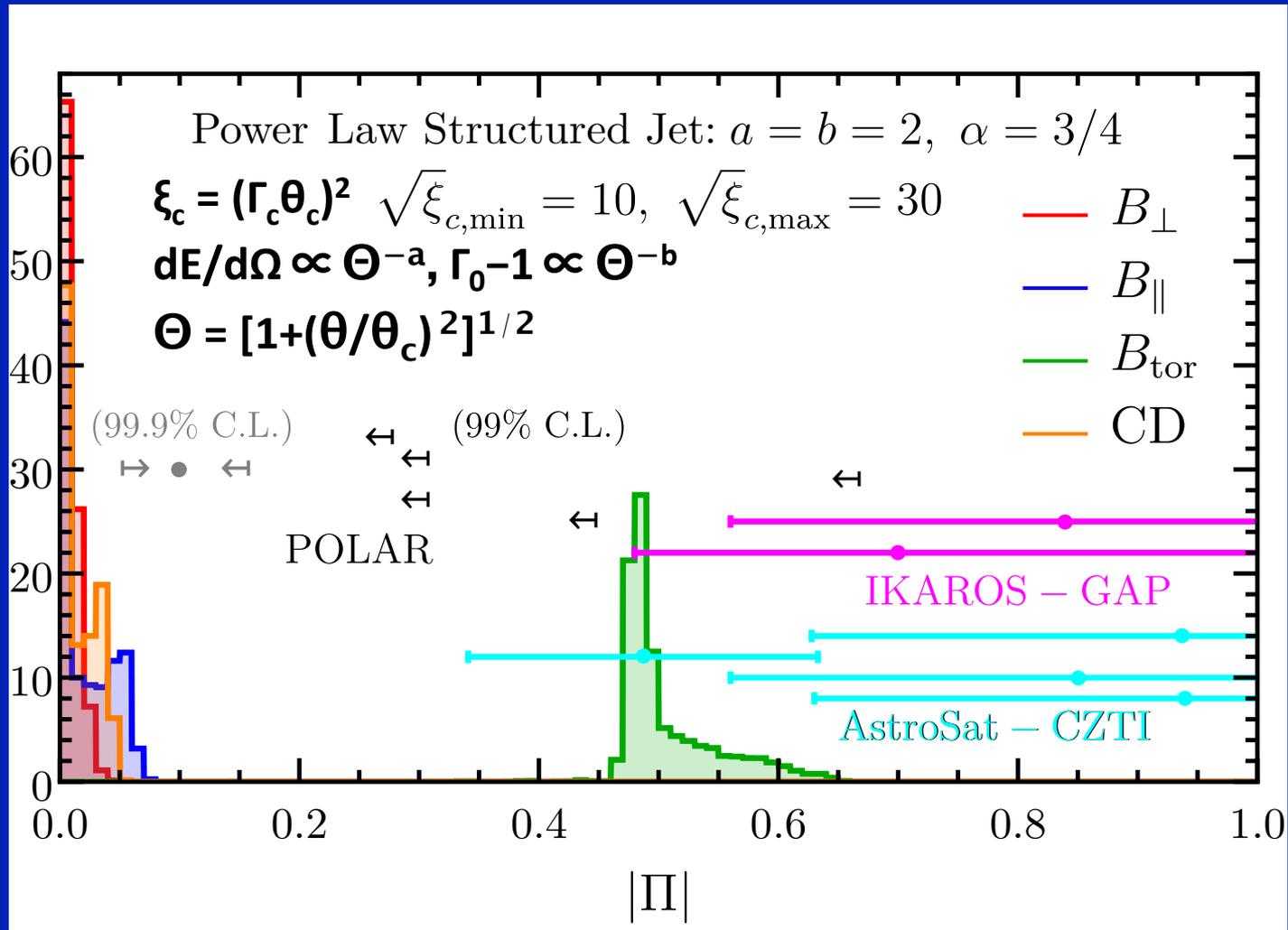
Ordered B-field



Toroidal B-field

Prompt GRB Polarization (Gill, JG & Kumar 2018):

Model comparison: structured jet, integrating 10 pulses



$B_{\text{tor}}/B_{\text{ord}}$ is favored if $P \sim 50-65\%$ in 1 ($\geq 20\%$ in most) GRBs

Conclusions:

- **Afterglow polarization** probes jet structure & dynamics + the B-field structure behind relativistic collisionless shocks
- ◆ \Rightarrow GW170817: $0.48 < \xi_0 < 0.79$ (B_{md}) + core-dominated jet

Conclusions:

- **Afterglow polarization** probes jet structure & dynamics + the B-field structure behind relativistic collisionless shocks
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- **Reverse shock polarization** probes B-field structure in ejecta
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 - ◆ Reverse & forward (afterglow) shock emission may overlap
 - ◆ GRB190114C: ~~\mathbf{B}_{ord} , axisymmetric ($\mathbf{B}_{\text{tor}}, \mathbf{B}_{\text{rnd}}$)~~, $\mathbf{B}_{\text{ord}} + \mathbf{B}_{\text{rnd}}$ X patchy shell?, incoherent patches: $\theta_B \sim 10^{-3}$ ✓

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- Prompt GRB pol. probes emission mechanism & jet structure
 - ◆ Observations are improving & new planned missions
 - ◆ Theory is improving to match the upcoming observations
 - ◆ $\mathbf{B}_{\text{ord}}/\mathbf{B}_{\text{tor}}$ favored if $P \sim 50\text{-}65\%$ in 1 ($\gtrsim 20\%$ in most) GRBs