# Astrophysical Searches for

# **Quantum Gravity Signals**

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on behalf of the Fermi LAT & GBM Collaborations

**Experimental search for quantum gravity: the hard facts Perimeter Institute (Colloquium), October 24, 2012** 

#### **Outline of the Talk:**

- Brief motivation & narrowing down the scope • Vacuum birefringence: helicity dependence of  $v_{ph}(E)$ Vacuum dispersion: energy dependence v<sub>ph</sub>(E) Using TeV flares from AGN Using GRBs: why, and how we set the limits ■ Limit from the bright long GRB 080916C at z~4.35 3 different types of limits from the short bright GRB 090510 at z = 0.903: detailed description & results Summary of Fermi GRB limits & future prospects
- Conclusions

#### **Quantum Gravity: a physics holy grail**

- Motivation: to unify in a self-consistent theory Einstein's general relativity that dominates on large scales & Quantum theory that dominates on small scales
- Quantum effects on space-time structure expected to become strong near the Planck scale:
- $l_{\text{Planck}} = (\hbar G/c^3)^{1/2} \approx 1.62 \times 10^{-33} \text{ cm}$  $E_{\text{Planck}} = M_{\text{Planck}}c^2 = (\hbar c^5/G)^{1/2}$ 
  - $\approx 1.22 \times 10^{19} \text{ GeV}$
- Many models / ideas out there: experimental constraints needed



#### Astrophysics as a test bed:

Advantage: large energies and distances available for free
 Disadvantage: uncontrolled experimental setup / conditions

• Vacuum birefringence: constrained by polarization Vacuum dispersion: by short timescale variability Pair production threshold: attenuation on the EBL Electron LIV: synchrotron radiation from the Crab nebula Space-time fuzziness: blur sources, broaden spectral lines ◆ UHECR / v LIV: energy spectrum / arrival time from GRBs Massive gravitons: supernovae cooling Cosmic string: gravitational lensing, gravity waves ◆ Early universe: CMB polarization, 21 cm HI line surveys...

Vacuum energy dispersion/birefringence Some quantum-gravity (QG) models (e.g. odd n SME) tie between vacuum dispersion & birefringence  $\Rightarrow$  makes life easier as birefringence is easier to constrain observationally Some models allow vacuum dispersion without birefringence • We directly constrain a simple form of LIV - dependence of the speed of light on the photon energy:  $v_{ph}(E_{ph}) \neq c$ This may be parameterized through a Taylor expansion of the LIV terms in the dispersion relation:  $c^2 p_{ph}^2 = E_{ph}^2 \left[ 1 + \sum_{k=1}^{\infty} S_k \left( \frac{E_{ph}}{M_{OG,k} c^2} \right)^k \right]$ , where  $M_{QG,k} \le M_{Planck}$  is naturally expected

s<sub>k</sub> = −1 ,0, 1 = model (helicity) dependent sign of the effect
 The most natural scale for LIV is the Planck scale
 l<sub>Planck</sub> ≈ 1.62×10<sup>-33</sup> cm ; E<sub>Planck</sub> = M<sub>Planck</sub>c<sup>2</sup> ≈ 1.22×10<sup>19</sup> GeV

Vacuum energy dispersion / birefringence
The photon propagation speed is given by the group velocity:

$$c^{2} p_{ph}^{2} = E_{ph}^{2} \left[ 1 + \sum_{k=1}^{\infty} S_{k} \left( \frac{E_{ph}}{M_{QG,k} c^{2}} \right)^{k} \right] , \quad v_{ph} = \frac{\partial E_{ph}}{\partial p_{ph}} \approx c \left[ 1 - S_{n} \frac{(1+n)}{2} \left( \frac{E_{ph}}{M_{QG,n} c^{2}} \right)^{n} \right]$$

■ Since  $E_{ph} \ll M_{QG,k}c^2 \leq E_{planck} \sim 10^{19}$  GeV the lowest order non-zero term, of order  $n = min\{k \mid s_k \neq 0\}$ , dominates

■ Usually n = 1 (linear) or 2 (quadratic) are considered

We focus here on n = 1, since only in this case are our limits of the order of the Planck scale

We try to constrain **both possible signs** of the effect:

◆ s<sub>n</sub> = 1, v<sub>ph</sub> < c: higher energy photons propagate slower</li>
 ◆ s<sub>n</sub> = -1, v<sub>ph</sub> > c: higher energy photons propagate faster
 ■ We stress: here c = v<sub>ph</sub>(E<sub>ph</sub>→0) is the low energy limit of v<sub>ph</sub>

Vacuum Birefringence: Polarization Helicity (left or right circular polarization) dependence of the photon propagation speed:  $c - v_{ph,L}(E) \approx v_{ph,R}(E) - c$ **\blacksquare** Rotates the position angle  $\theta$  of linearly polarized radiation:  $\Delta \phi_{R,L} = 2\Delta \theta = \omega \Delta t_{R,L} \approx \omega \Delta v_{R,L} D/c^2 \approx E^{n+1} D(1+n)/\hbar c (E_{OG^*,n})^n$  $\Delta E/E \gtrsim 0.2-1 \Longrightarrow \Delta \theta(E_2) \sim 2\Delta \theta(E_1)$  $\Delta \theta(E_1) \ge 1 \Longrightarrow$  depolarization  $\blacksquare \Rightarrow$  linear pol. constrains  $E_{QG^*,n} = \xi_{1^*}E_{planck}$ : ◆ Galaxy at D ~ 0.3 Gpc, optical:  $P \sim 10\% \implies \xi_{1*} > 5 \times 10^3$  (Gleizer & Nozameh 01)  $\diamond$  Crab nebula (Galactic SNR;  $D \approx 2$  kpc) X/ $\gamma$ -rays: P ~ 46% (INTEGRAL 150-300 keV)  $\Rightarrow \xi_{1*} > 1.1 \times 10^9$  (99% CL; Maccione et al. 2008)

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 $\Rightarrow \xi_{1*} > 10^{15}$ (Toma et al. 2012)

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Vacuum dispersion: time variability Relevant models without leading order vacuum birefringence Good candidate sources: **TeV flares** from AGN **AGN:** accreting super-massive black holes galaxy centers • Mass:  $M_{BH} \sim 10^6 - 10^9 M_{\odot}$  Jet Lorentz factor:  $\Gamma \sim 5 - 30^{\circ}$ Active for millions of years 10 kpc kpc Sometimes emit short bight flares AGN jet in M87 (VLBA 43 GHz)  $D \approx 16 \text{ Mpc}$ (X-rays; Chandra) composite: X-rays, **Centaurus A** ( $D \approx 3.6$  Mpc) optical, sub-mm) • 6 r.

• MAGIC (07,08): Mkn 501, D  $\approx$  140 Mpc; 0.17-10 TeV; t<sub>var</sub>  $\sim$  120 s  $\Delta t_{LIV} \approx t \Delta v_{LIV}/c \approx \frac{1}{2}(1+n) [\Delta E^n/(E_{QG,n})^n] D/c$ claimed a possible detection:  $\xi_1 \sim 0.03$ ,  $E_{QG,2} \sim 6 \times 10^{10}$  GeV or alternatively lower limits:  $\xi_1 > 0.02$ ,  $E_{OG,2} > 4 \times 10^{10}$  GeV



#### Vacuum dispersion: time variability

#### ◆ HESS (08,11): PKS 2155-304, D ≈ 480 Mpc; 0.2-5 TeV









#### **Probing Vacuum dispersion Using GRBs** (first suggested by Amelino-Camelia et al. 1998)

Why GRBs?Very bright & short<br/>transient events, at cosmological<br/>distances, emit high-energy γ-rays(D. Pile, Nature Photonics, 2010)

vannin

#### **GRB** Theoretical Framework:

#### Progenitors:

Long: massive stars
Short: binary merger?
Acceleration: fireball or magnetic?
Prompt γ-rays: internal shocks? emission mechanism?



Deceleration: the outflow decelerates (by a reverse shock for σ ≤ 1) as it sweeps-up the external medium
 Afterglow: from the long lived forward shock going into the external medium; as the shock decelerates the typical frequency decreases: X-ray → optical → radio

Fermi Gamma-ray **Space Telescope** (launched on June 11, 2008)



LAT

 $10^{3}$ 

 $10^{4}$ 

10<sup>5</sup>

Fermi GRB Monitor (GBM): 8 keV – 40 MeV  $(12 \times \text{NaI 8} - 10^3 \text{ keV}, 2 \times \text{BGO 0.15} - 40 \text{ MeV})$ , full sky Comparable sensitivity + larger energy range than its predecessor - BATSE ■ Large Area Telescope (LAT): 20 MeV – >300 GeV FoV ~ 2.4 sr; up to  $40 \times EGRET$  sensitivity,  $\ll$  deadtime



# The Fermi Observatory

Large Area Telescope (LAT) -• Large Field of

- View (>2.4 sr)
- views entire sky every 3 hrs
- 20 MeV -300 GeV



Gamma-ray Burst Monitor (GBM) • Views entire unocculted sky • Nal: 8 keV - 1 MeV • BGO: 0.15 - 30 MeV

## The Large Area Telescope

#### **Pair-conversion** γ-ray detector

- Energy range: 20MeV >300GeV
- GeV photons useful for LIV studies with GRBs
- Wide field of view (~±70°); large effective area
- Helps with detecting many GRBs with ample photon statistics per detection
- Good angular (~0.2° at 1GeV) and energy resolution (~10% over 1GeV), low bkg rate (<1Hz in ROI over 20MeV)</li>
- Provides high-quality data for LIV studies
- In first 3 years:
- Detected 10 GRBs with a measured redshift
- 21 GRBs with emission over 1 GeV

Range of redshifts extends from 0.74 to 4.35



#### **Constraining LIV Using GRBs**

A high-energy photon E<sub>h</sub> would arrive after (in the sub-luminal case: v<sub>ph</sub> < c, s<sub>n</sub> = 1), or possibly before (in the super-luminal case, v<sub>ph</sub> > c, s<sub>n</sub> = -1) a low-energy photon E<sub>l</sub> emitted together

The time delay in the arrival of the high-energy photon is:

$$\Delta t_{\rm LIV} = S_n \frac{(1+n)}{2H_0} \frac{E_h^n - E_l^n}{\left(M_{\rm QG,n}c^2\right)^n} \int_0^z \frac{(1+z')^n}{\sqrt{\Omega_m (1+z')^3 + \Omega_\Lambda}} dz'$$
(Jacob & Piran 2008)

The photons E<sub>h</sub> & E<sub>l</sub> do not have to be emitted at exactly the same time & place in the source, but we must be able to limit the difference in their effective emission times, i.e. in their arrival times to an observer near the GRB along our L.O.S



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• Our limits apply to any source of energy dispersion on the way from the source to us, and may constrain some (even more) exotic physics ( $\Delta t_{LIV} \rightarrow \Delta t_{LIV} + \Delta t_{exotic}$ )

#### Method 1

■ Limits only  $s_n = 1$  - the sub-luminal case:  $v_{ph} < c$ , & positive time delay,  $\Delta t_{LIV} = t_h - t_{em} > 0$  (here  $t_h$  is the actual measured arrival time, while  $t_{em}$  would be the arrival time if  $v_{ph} = c$ )

• We consider a single high-energy photon of energy  $E_h$  and assume that it was emitted after the onset time  $(t_{start})$  of the relevant low-energy  $(E_l)$  emission episode:  $t_{em} > t_{start}$ 

 $\blacksquare \Longrightarrow \Delta t_{\rm LIV} = t_{\rm h} - t_{\rm em} < t_{\rm h} - t_{\rm start}$ 

A conservative assumption: t<sub>start</sub> = the onset of any observed emission from the GRB

#### Limits on LIV: GRB080916C ( $z \approx 4.35$ )

**GRB080916C:** highest energy photon (13 GeV) arrived 16.5 s after lowenergy photons started arriving (=the GRB trigger)  $\Rightarrow$  conservative lower limit:  $M_{OG,1} > 1.3 \times 10^{18} \text{ GeV/c}^2$  $\approx 0.11 M_{Planck}$ 

This improved upon the previous limits of this type, reaching 11% of M<sub>Planck</sub>

AGN

(Biller 98)

GRB

Pulsar

1015

(Kaaret 99) (Ellis 06)

 $1.8 \times 10^{15} 0.9 \times 10^{16} 10^{16}$ 



GRB090510: L.I.V ■ A short GRB (duration ~1 s) Redshift:  $z = 0.903 \pm 0.003$ ■ A ~31 GeV photon arrived at  $t_{\rm h} = 0.829$  s after the trigger We carefully verified it is a photon; from the GRB at  $>5\sigma$ We use the  $1-\sigma$  lower bounds on the measured values of  $E_{h}$ (28 GeV) and z (0.900)Intrinsic spectral lags known on timescale of individual pulses: weak effect expected



GRB090510: L.I.V

Method 1: different choices of t<sub>start</sub> from the most conservative to the least conservative

 $t_{start} = -0.03$  s precursor onset  $\Rightarrow \xi_1 = M_{OG.1}/M_{Planck} > 1.19$  $t_{start} = 0.53 s$  onset of main emission episode  $\Rightarrow \xi_1 > 3.42$ For any reasonable emission spectrum a ~31 GeV photon is accompanied by many  $\gamma$ 's above<sup>§</sup> 0.1 or 1 GeV that "mark" its t<sub>em</sub>  $t_{start} = 0.63 \text{ s}, 0.73 \text{ s} \text{ onset of}$ emission above 0.1, 1 GeV  $\Rightarrow \xi_1 > 5.12, \xi_1 > 10.0$ 



## GRB090510: L.I.V

Troja et al. 2010: detection of low level emission ~13 s before the GRB090510 trigger Highly unlikely that no other LAT photons were emitted together with the ~31 GeV photon (none were observed) • Fine tuning is required for the ~31 GeV photon to arrive on top of brightest emission episode (+on a narrow spike)





### GRB090510: L.I.V

Method 2: least conservative Associating a high energy photon with a sharp spike in the low energy lightcurve, which it falls on top of Limits both signs:  $s_n = \pm 1$ Non-negligible chance probability (~5-10%), but still provides useful information ■ For the 31 GeV photon (*shaded* vertical region)  $\Rightarrow \Delta t < 10 \, ms$ and  $\xi_1 = M_{OG,1}/M_{Planck} > 102$ For a 0.75 GeV photon during precursor:  $|\Delta t| < 19 \text{ ms}, \xi_1 > 1.33$ 



#### Method 3: DisCan (Scargle et al. 2008)

- Based on lack of smearing of the fine time structure (sharp narrow spikes in the lightcurve) due to energy dispersion
- Constrains both possible signs of the effect:  $s_n = \pm 1$
- Uses all LAT photons during the brightest emission episode (obs. range 35 MeV – 31 GeV); no binning in time or energy
- Shifts the arrival time of photons according to a trail energy dispersion (linear in our case), finding the coefficient that maximizes a measure of the resulting lightcurve variability
- We found a symmetric upper limit on a linear dispersion:  $|\Delta t/\Delta E| < 30 \text{ ms/GeV} (99\% \text{ CL}) \Rightarrow M_{QG,1} > 1.22M_{Planck}$
- Remains unchanged when using only photons < 1 or 3 GeV (a very robust limit)

	$t_{\rm start}$ limit on		Reason for choice of	$E_l$	valid	lower limit on	limit on $M_{\rm QG,2}$
	(ms)	$ \Delta t $ (ms)	$t_{ m start}$ or limit on $\Delta t$	(MeV)	for $s_n$	$M_{\rm QG,1}/M_{\rm Planck}$	in $10^{10} \text{ GeV}/c^2$
<b>GRB090510</b>	-30	< 859	start of any observed emission	0.1	1	> 1.19	> 2.99
b	530	< 299	start of main $< 1 \mathrm{MeV}$ emission	0.1	1	> 3.42	> 5.06
line to an I IV	630	< 199	start of $> 100$ MeV emission	100	1	> 5.12	> 6.20
limits on LIV	730	< 99	start of $> 1 \text{ GeV}$ emission	1000	1	> 10.0	> 8.79
e		< 10	association with $< 1 \mathrm{MeV}$ spike	0.1	±1	> 102	> 27.7
Summary:	—	< 19	if $0.75{ m GeV}~\gamma$ is from $1^{ m st}$ spike	0.1	-1	> 1.33	> 0.54
Summary.	$ \Delta t/\Delta $	$E  < 30 \mathrm{ms/GeV}$	lag analysis of all LAT events	—	±1	> 1.22	—

**Our results disfavor QG models with linear (n = 1)**  $v_{ph}(E)$ 

**a**-e based on 31 GeV γ-ray a-d method 1:  $t_{em} \ge t_{strat}$ **e**,**f**: method 2: association with a low-energy spike **g**: method 3: DisCan sharpness of HE spikes All of our lower limits on M<sub>OG.1</sub> are above M<sub>Planck</sub>



# Limits on LIV from Fermi GRBs

GRB	duration or class	# of events > 0.1 GeV	# of events > 1 GeV	method	Lower Limit on M <sub>QG,1</sub> /M <sub>Planck</sub>	Valid for S <sub>n</sub> =	Highest photon Energy	redshift
080916C	long	145	14	1	0.11	+1	~ 13 GeV	~ 4.35
				1	<b>1.2</b> , 3.4, <b>5.1</b> , <b>10</b>	+1		
090510	short	> 150	> 20	2	102	±1	~ 31 GeV	0.903
				3	1.2	±1		
090902B	long	> 200	> 30	1	0.068	+1	~ 33 GeV	1.822
090926	long	> 150	> 50	1,3	0.066, 0.082	+1	~ 20 GeV	2.1062

Method 1: assuming a high-energy photon is not emitted before the onset of the relevant low-energy emission episode
 Method 2: associating a high-energy photon with a spike in the low-energy light-curve that it coincides with
 Method 3: DisCan (dispersion cancelation; very robust) – lack of smearing of narrow spikes in high-energy light-curve

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				1	1.2, 3.4, 5.1, 10	+1		
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New (work in progress): using 3 different analysis methods that are complimentary & somewhat more sensitive on the same 4 brightest Fermi/LAT GRBs with known redshifts:
 We improve the limits (the new numbers are still not public)

### Limits on LIV from Fermi GRBs

 A. PairView: calculates spectral lags *l*<sub>i,j</sub> between all pairs of photons in a dataset and identifies the most prominent value of *l*<sub>i,j</sub> as the best estimate of the LIV parameter *l*<sub>i,j</sub> = t<sub>i</sub> - t<sub>j</sub>/E<sup>n</sup><sub>i>j</sub> = t<sub>i</sub> - t<sub>j</sub>/E<sup>n</sup><sub>i>j</sub>
 The distribution of *l*<sub>i,i</sub> will

have a peak approximately centered at the true value  $\tau_n$ 

- -If data has no lag, there will still be a peak but at zero
- -Peak width/height depend on statistical strength of the dataset:



many GeV photons in a bright pulse will give the strongest signal
B. Sharpness Maximization: based on idea similar to DisCan
C. Likelihood analysis: used before on AGN – low-energy lightcurve + spectrum template used to claculate unbinned likelihood for high-energy data as a function of τ<sub>n</sub> = Δt/ΔE<sup>n</sup>

# The Future - CTA Energy range: ~20 GeV to ~500 TeV an order of magnitude more sensitive than current instruments around 1 TeV (~150M€ price tag), better angular/energy resolution >1000 members in 27 countries Preparatory Phase 2010-2013, construction 2013-2018 2 sites (southern + northern hemispheres) Hundreds of telescopes of 3 different sizes

#### A bigger difference for transient sources



**Prospects for LIV studies with CTA GRBs** Method 1: it may be difficult to do much better • Our current limit  $|\Delta t/\Delta E| < 30 \text{ ms/GeV}$  would require  $E_h > 1$  TeV for a response time of 30 s  $\diamond$  at > 1 TeV intrinsically fewer photons + EBL Method 3: might work best 7RB090926A Sharp bright spikes up to high energies exist also 1000 500 well within long GRBs  $\bullet t_{var} \sim 0.1 \text{ s \& } E_h \sim 0.1 \text{ TeV}$ Counts/Bin could do ~30 times better A short GRB in CTA FoV ounts/Bin (survey mode) would be great **10 ms, 1 TeV: >10<sup>3</sup> times better** 

# **Conclusions:**

Astrophysical tests of QG can help – look for them **GRBs** are very useful for constraining LIV Bright short GRBs are more useful than long ones A very robust and conservative limit on a linear energy dispersion of either sign:  $M_{QG,1} > 1.2M_{Planck}$ Still conservative but somewhat less robust limits:  $M_{QG,1}/M_{Planck} > 5.1, 10$  (onset of emission >0.1, 1 GeV) "Intuition builder" liberal limit: M<sub>OG,1</sub> / M<sub>planck</sub> > 102 Quantum-Gravity Models with linear (n = 1)photon energy dispersion are disfavored