Prompt GRB Emission: High-Energy Spectra & Polarization

> **Jonathan Granot** Open University of Israel

Collaborators: Fermi-LAT collaboration, Arieh Königl, Greg Taylor, Ramandeep Gill, Pawan Kumar

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Outline of the Talk:

Polarization in GRBs:

 Motivation, synchrotron emission, B-field structure Relativistic beaming, jet + viewing angle effects • Alternative models: photospheric, Compton drag Model comparison: different pulses vs. whole GRB Fermi: new perspectives on prompt GRB spectra High-energy & photospheric spectral components \square Constraints on Γ : the compactness problem **Time-dependent intrinsic** $\gamma \gamma \rightarrow e^+e^-$ opacity model

Why is GRB Polarization Interesting

- It teaches us about the magnetic field structure in the GRB ejecta & provides clues as to whether most of energy is in Poynting flux or kinetic energy:
 - $E_{EM} \gg E_{kin} \Rightarrow$ ordered magnetic field is expected
 - ◆ $E_{kin} \ge E_{EM}$ ⇒ ordered & random fields are possible
- Provides a strong test for the jet structure in GRBs, both in the prompt GRB & in the afterglow
- **Constrains the prompt GRB emission mechanism**
- Probes magnetic field structure behind afterglow shock
 Helps pin down cause of time variability in afterglows

Polarization of Synchrotron Emission



linear polarization perpendicular to the projection of
 B on the plane of the sky

The maximal polarization is for the local emission from an ordered **B**-field: $P_{max} = (\alpha+1)/(\alpha+5/3)$ where $F_v \propto v^{-\alpha}, -1/3 \leq \alpha \leq 1.5 \Rightarrow 50\% \leq P_{max} \leq 80\%$ (Rybicki & Lightman 1979; Granot 2003)

Shock Produced Magnetic Field:

A magnetic field that is produced at a relativistic collisionless shock, due to the two-stream instability, is expected to be tangled within the plane of the shock (Medvedev & Loeb 1999)





Polarization in the observer frame

Random field in shock plane



Ordered field in shock plane



Sari 99; Ghisellni & Lazzati 99



Granot & Königl 03

 $P \sim P_{max}$

Polarization of Prompt γ-ray emission:

Observations: very hard to measure in γ-ray or hard X-rays
GRB021206 P=80±20% RHESSI (Coburn & Boggs 2003)

Soon refuted (Rutledge & Fox 2003; Wigger et al. 2004)

- GRB930131, 960924 P>35, 50% CGRO/BATSE (Willis+05)
- GRB041219a 98±33%, 98±33%, 43±25% Integral/SPI,IBIS (Kalemci et al. 2007; McGlynn et al. 2007; Götz et al. 2009)
- GRB061122 P>60% (1σ),33% (90%) Integral/IBIS (Götz+13)
- GRB140206a P>48% (1σ), 28% (90%) Integral/IBIS (Götz+14)
- GRB100826a 27±11%, 110301a 70±22%, 110721a 80±22%
 Ikaros/GAP (Yonetoku et al. 2011, 2012)
- Astrosat/CZTI ???

Polarization of Prompt γ-ray emission:
Theory: first consider synchrotron emission
Shock produced B-field + θ_{obs} ≤ θ_j-1/Γ ⇒ P ≈ 0
P ~ P_{max} can be achieved in the following ways:

ordered magnetic field in the ejecta,
special geometry: |θ_{obs} - θ_j| ≤1/Γ ⇒ favors narrow jets: θ_j ≤ 1/Γ (works with a shock produced B-field)



Ordered Magnetic Field in the Ejecta:

0.9 0.8

max

Total emission from jet

Afterglow: instantaneous emission

Prompt GRB: time integrated emission



0.70.60.5 ord afterglow 0.4 prompt GRB Ω 0.3 0.70.6 0.5 0.4 0.3afterglow 0.2 prompt GRB 0.1 0 _0.25 0.25 0.5 0 0.751.25 α (Granot 2003) **F** $\propto v^{-\alpha}$, **P** increases with α **P** from an ordered **B**-field is slightly larger in afterglow

1.5

Narrow Jet + shock produced B-field

- High polarization + reasonable flux $\Rightarrow \theta_i < \theta_{obs} \leq \theta_i + 1/\Gamma$
- A reasonable probability for such $\theta_{obs} \Rightarrow \Gamma \theta_i \leq a$ few
- Since $\Gamma \ge 100 \& \theta_i \ge 0.05$, $\Gamma \theta_i \ge 5$ and is typically larger
- However GRB 021206 was very bright, suggesting a very narrow jet: $f = 1.6 \times 10^{-4}$ erg which for z~1 implies $E_{iso} \sim 10^{54}$ erg & $\theta_j \sim (10^{51} \text{erg/E}_{iso})^{1/2} \sim 0.03$ (Frail et al. 01)

 $\blacksquare \Rightarrow \Gamma \theta_{j} \sim 3(\Gamma/100) \Rightarrow \Gamma \theta_{j} \leq a \text{ few is possible (Waxman 03)}$

- The jet must have sharp edges: $\Delta \theta_i \leq 1/4\Gamma$ (Nakar et al. 03)
- a 'structured jet' produces low polarization (several %)
- Most GRBs are viewed from $\theta_{obs} < \theta_j$ and are expected to have a very low polarization in this scenario



ΔΓ ~ Γ between different shell collisions (different pulses in GRB light curve) reduces P by a factor ~ 2

	Ordered Field	Narrow Jet
P ~ 80%	X	X
P ~ 50%	\checkmark	X
P~25%	with B _{rnd} ≤ B _{ord}	\checkmark
P ≤ 10%	with $B_{rnd} > B_{ord}$	with B _{rnd} ≥ B _{ord}
statistics	High P in all GRBs	low P in most GRBs
Optical flash	High P - similar to the prompt GRB	Similar to prompt GRB (low P in most GRBs)
Potential problems	Some B _{rnd} required for Fermi acceleration	$ \Gamma θ_j ≤ a \text{ few, } ΔΓ ~ Γ, Brnd (afterglow obs.) $

Alternative to Synchrotron: Compton Drag (Bulk Inverse Compton Scattering of External photons) (Lazzati et al. 2003; Dar & De Rujula 2003, Eichler & Levinson 2003)
Requires special geometry/viewing angle, θ_j < θ_{obs} ≤ θ_j+1/Γ
Polarization properties similar to synchrotron + B_{rnd} with an advantage: local polarization P=(1-cos²θ)/(1+cos²θ) can reach up to 100% while P_{max}~70% for synchrotron

Shares drawbacks of shock produced field + narrow jet

- It has additional problems, unrelated to polarization:
 - Explaining prompt GRB spectrum
 - Supplying external photons for all the ejected shells
 - High photon density \Rightarrow small radii \Rightarrow high $\tau_{\gamma\gamma}$



Alternative to Synchrotron: Photospheric Emission (Comptonized radiation advected from optically thick to thin region of the jet) (Beloborodov 2011; Thompson & Gill 2014; Lundman + 2014; Vurm & Beloborodov 2016; Lundman + 2016)

- Need to integrate radiation transfer equations for the Stokes parameters $I(r,\mu) \& Q(r,\mu)$ from $\tau_T \gg 1$ to $\tau_T \ll 1$.
- P=0 seed photons become anisotropic at $\tau_T \leq 10 \implies P \approx 0.45P_{Compton-drag}$
- This requires symmetry breaking e.g.
 special viewing angle: |θ_{obs}-θ_j| ≤ 1/Γ
 θ-dependent bulk-Γ and/or luminosity (in structured jets P ≤ 40%)
- Synchrotron + B_{ord} (spherical flow): Unscattered syn. photons emitted at $\tau_T \sim 1$ dominate at $E \ll E_{pk} \Rightarrow P \sim P_{syn,max}$



Different Pulses in the same GRB:

■ Harder to measure (requires very bright GRB), more rewarding

• $\theta_i = \text{const}^*$, $\Gamma \text{ varies} \Rightarrow q = \theta_{obs}/\theta_i = \text{const}^*$, $y_i = ((\Gamma \theta_i))^2$ varies

- Synchrotron + B_{ord} : $\theta_p = const$, P ~ $P_{max} \sim 50\%$ also when integrating over all the GRB pulses
- Models requiring special viewing angle (Synchrotron+ \mathbf{B}_{rnd} , Compton drag, photospheric emission) $\theta_i < \theta_{obs} \le \theta_i + 1/\Gamma$:
- ◆ For a sharp-edged jet the special condition on θ_{obs} occurs only in some pulses ⇒ P varies between pulses (θ_p may flip by 90°) ⇒ smaller integrated P over all the GRB
- For a smooth-edged jet, e.g. Gaussian or core+power-law wings, P is low for any θ_{obs} even in a single pulse (unless Γθ_{i/e}<1)



Fermi Gamma-ray Space Telescope (launched on June 11, 2008)



 Fermi GRB Monitor (GBM): 8 keV – 40 MeV (12×NaI 8 – 10³ keV, 2×BGO 0.15 – 40 MeV), full sky
 Comparable sensitivity + larger energy range than its predecessor - BATSE
 Large Area Telescope (LAT): 20 MeV – >300 GeV FoV ~ 2.4 sr; up to 40× EGRET sensitivity, « deadtime





Delayed onset of High-Energy Emission GRB080916C GRB090510



(Abdo et al. 2009, Science, 323, 1688)

The 1st LAT peak coincides with the 2nd GBM peak Delay in HE onset: ~ 4-5 s

The first few GBM peaks are missing in LAT but later peaks coincide; the delay is 0.1-0.2s

Distinct High-Energy Spectral Component







Clearly (>5o) exists in several LAT GRBs, but very common in the brightest LAT GRBs
Suggests that it is common but good photon statistics is needed for clear evidence





Late onset/HE spectral component: Possible Origin Leptonic: inverse-Compton (or synchrotron self-Compton)? • Hard to produce a delayed onset longer than spike widths (the seed photon field builds-up on the dynamical time) \diamond A gradual increase in the HE photon index β (determined by the electron energy dist.) is not naturally expected • Hard to account for the different photon index values of the HE component & the Band spectrum at low energies ◆ Hard to produce a low-energy power-law (GRB090902B)



Late onset/HE spectral component: Possible Origin **Hadronic**: (pair cascades, proton synchrotron)? ◆ Late onset: time to accelerate protons+develop cascades? \diamond Does not naturally account the gradual increase in β • Hard to produce the observed sharp spikes that coincide with those at low energies (+ a longer delay in the onset) • GRB090510: large energy needed: $E_{total}/E_{v,iso} \sim 10^2 - 10^3$ ◆ GRB090902B: synchrotron emission from secondary e[±] pairs can naturally explain the power-law at low energies



GRB: High Energy Emission Processes Leptonic: Inverse-Compton or Synchrotron-Self Compton: $E_{p,SSC}/E_{p,syn} \sim \gamma_e^2$, $L_{SSC}/L_{syn} = Y$, $Y(1+Y) \sim \epsilon_{rad}\epsilon_e/\epsilon_B$



■ Hadronic processes: photopair production $(p+\gamma \rightarrow p+e^+e^-)$, proton synchrotron, pion production via $p-\gamma$ (photopion) interaction or p-p collisions

GRB: High Energy Emission Processes Leptonic: Inverse-Compton or Synchrotron-Self Compton: $E_{p,SSC}/E_{p,syn} \sim \gamma_e^2$, $L_{SSC}/L_{syn} = Y$, $Y(1+Y) \sim \epsilon_{rad}\epsilon_e/\epsilon_B$ **GRBs 090510, 090926A:** Y varies, and sometimes Y > 1 **GRB080916C**: single dominant emission mechanism \Rightarrow if synchrotron, SSC is expected, and can avoid detection if $E_{\text{peak,SSC}} \gg 10 \text{ GeV} (\gamma_e \gg 100), \text{ or if } Y \approx \epsilon_e / \epsilon_B \leq 0.1$



GRB: High Energy Emission Processes **Leptonic**: Inverse-Compton or Synchrotron-Self Compton: $E_{p,SSC}/E_{p,syn} \sim \gamma_e^2$, $L_{SSC}/L_{syn} = Y$, $Y(1+Y) \sim \epsilon_{rad}\epsilon_e/\epsilon_B$ **GRBs 090510, 090926A:** Y varies, and sometimes Y > 1 **GRB080916C**: single dominant emission mechanism \Rightarrow if synchrotron, SSC is expected, and can avoid detection if $E_{\text{peak},\text{SSC}} \gg 10 \text{ GeV} (\gamma_e \gg 100), \text{ or if } Y \approx \epsilon_e / \epsilon_B \leq 0.1$ **Parameter space study** (Benyamini & Piran 2013): $0.1 < \epsilon_{\rm e}/\epsilon_{\rm B} < 10^4 \ (0.1 \leq Y \leq 100), 300 \leq \Gamma \leq 3000,$ $3 \times 10^3 \leq \gamma_e \leq 10^5$, 10^{15} cm $\leq \mathbf{R} \leq 10^{17}$ cm $(E_{\text{peak,SSC}} \sim E_{\text{KN}} \sim \Gamma \gamma_{e} m_{e} c^{2} \sim 1.6(1+z)^{-1} \Gamma_{2.5} \gamma_{e.4} \text{ TeV} \Longrightarrow \text{CTA?})$

Thermal components in prompt spectrum? ■ Usually sub-dominant ⇒ degeneracy with the assumed (usually phenomenological Band) dominant component









Photospheric components

- Suggested in some cases by low energy data ($kT \leq 0.1 \text{ MeV}$)
- Usually sub-dominant energetically (+not unique interpretation)
- In the Fireball Model: a remnant of the thermal acceleration
 - $E_{ph}/E = T_{ph}/T_0 = 0.05E_{52}^{-2/3}R_{0,6}^{2/3}t_1^{2/3}\Gamma_{2.5}^{8/3} \text{ (Nakar et al. 2005)}$ $kT_0 = 3(1+z)^{-1}E_{52}^{1/4}R_{0,6}^{-1/2}t_1^{-1/4} \text{ MeV} \qquad t = T_{GRB}/(1+z)$

 $kT_{ph} = 300(1+z)^{-1}E_{th,51}E_{52}^{1/4}R_{0,6}^{-1/2}t_1^{-1/4} keV$

For magnetic acceleration:
 Dissipation below the photosphere can give such a spectral component
 can arise from gradual reconnection or multiple passages of weak shocks



Constraints on \Gamma for Fermi LAT GRBs F_{min}: no high-energy cutoff due to intrinsic pair production \Rightarrow lower limit on the Lorentz factor of the emitting region **Fermi**: more robust limits – don't assume photons >E_{obs max} For bright LAT GRBs (long/short): $\Gamma \ge 10^3$ for simple model (steady-state, uniform, isotropic) but $\Gamma \gtrsim 10^{2.5}$ for more realistic time-dependent self-consistent thin shell model (JG et al. 2008) **GRB 090926A:** high-energy cutoff – if due to intrinsic pair production then $\Gamma \sim 300 - 700$



Chine Compactness Problem: (Schmidt 1978; Fenimore et al. 1993; Woods & Loeb 1995;...)

The large γ-ray flux implies huge luminosities for cosmological GRBs, L_{iso} ~ 10⁵⁰ - 10⁵³ erg/s

■ For sources at rest: short variability time $\Delta t \Rightarrow$ small source R < c Δt & $\epsilon = E_{ph}/m_ec^2 \sim 1 \Rightarrow$ large fraction of γ 's can pair produce ($\gamma\gamma \rightarrow e^+e^-$)

τ_{γγ}(ε) ~ σ_Tn_{ph}(1/ε)R, n_{ph}(1/ε) ~ L_{1/ε}/4πR²m_ec³ ⇒ τ_{γγ}(ε) ~ σ_TL_{1/ε}/4πm_ec³R ≥ 10¹⁴ L_{1/ε,51}(Δt / 1 ms)⁻¹
 Such a huge τ_{γγ} would produce a thermal spectrum ⇒ inconsistent with the observed high energy tail

Solution: Relativistic Motion Γ »1

The effects of special relativity help us out here:

The (relativistic) Doppler effect:
 E_{ph}/E'_{ph} = 1/Γ(1-βcosθ) ~ Γ ^y
 Aberration of light or

relativistic beaming:

• For $\theta' = 90^\circ$, $\cos\theta = \beta$ $(\theta \approx 1/\Gamma \text{ for } \Gamma \gg 1)$

The direction of photons in the lab frame for a point source emitting isotropically in its own rest frame and moving to the right at different proper speeds Γβ:





Time Dependence of Intrinsic $\gamma\gamma \rightarrow e^+e^-$ Opacity Semi-Analytic Model (JG, Cohen-Tanugi & do Couto e Silva 08) ■ Ultra-relativistic ($\Gamma \gg 1$) spherical thin ($\Delta \ll R/\Gamma^2$) shell emits over a finite range of radii $R_0 \leq R \leq R_0 + \Delta R$ Emission: uniform over the shell & isotropic in its local (co-moving) rest frame; $L'_{\nu\nu} \propto (\nu')^{1-\alpha} R^{b}$, $\Gamma^{2} \propto R^{-m}$ **F**, calculated by integral over **equal arrival time surface The photon field is fully calculated at all space & time** $\mathbf{I}_{\gamma\gamma}$ calculated by **integral along trajectory** of each photon expanding shell Corresponds to a single spike/flare in the light curve $\Delta R \leq R_0$ impulsive $\gamma\gamma \rightarrow e^+e^-$ 0.9 $\Delta R >> R_0$ quasi-steady 0.8 normalized flux 0.7 0.6 GRB 0.5 emission "turns "turns observer 0.4 0.3 region on" Off at 0.2 infinity 0.1 R_0 0.5 3.5 4.5 5 $R_0 + \Delta R$ (T/T_0)

Results: Light Curves & Instantaneous Spectra
 There is no initial photon field - opacity gradually builds-up: ε > ε₁(steady state) photons can escape early on while ε₁(t) > ε
 ⇒ a distinct temporal & spectral signature



Time Integrated Spectrum: power-law high energy tail



Unique signature: High-energy γ -rays above break in time integrated spectrum escape mainly near the onset of a flare/ spike in the light curve



The opacity builds-up & saturates on a dynamical time scale Spectral Cutoffs at ~100 MeV (Vianello, Gill, Granot, Omodei, Cohen-Tanugi & Longo 2017)
 Time-resolved spectra of 2 GRBs with cutoffs E_c ≤ 100 MeV
 Best fit phenomenological model + 2 theoretical models fits
 Gill & Thompson 2014: photospheric (high-σ breakout from star)



Spectral Cutoffs at ~100 MeV (Vianello, Gill, Granot, Omodei, Cohen-Tanugi & Longo 2017) Granot et al. 2008: model parameters inferred from the fits: $\epsilon'(\Gamma_{max}) = 1 \Longrightarrow$ model is self-consistent for $\Gamma_0 < \Gamma_{max}$ ($\tau_{T\pm} < 1$) GRB100724B E_c~ 20 – 60 MeV GRB160509A z = 1.17 E_c ~ 80 - 150 MeV -1.4800 -1.6700 500 -1.6 -1.8600 () 500 () 400 (keV) -1.8 400 ∞_ −2.0 ∞_-2.0 400 , 400 ° 300 Ш -2.2 300 ເມີ -2.2 -2.4200 200-2.4-2.610² ΔR / R₀ s പ്പ 10¹ erg erg 9 ΔR 10^{52} 3 3 2 10⁰ 10 4×1 3×10^{2} **∔**∳ 4×10^{2} 10⁰ Γ₀/Γ_{max} 2×10^{2} 3×10^{2} 10^{0} டீ 2×10^{2} 10² 10^{-1} 6×10 10^{2} $+ R_0 + R_f$ + R₀ + R_f 10² R (10¹³ cm) R (10¹³ cm) ک 10¹ 10¹) bo 10⁰ 10⁰ 100 15 20 25 50 100 150 0 50 150 10 10 15 20 25 Time since trigger (s) Time since trigger (s) Time since trigger (s) Time since trigger (s)

